

# Short range correlation physics in low resolution pictures

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UNIVERSITY

Scott Bogner  
Facility for Rare Isotope Beams  
Michigan State University

**Sushant More**



FRIB

# Low and High Resolution Scale Pictures

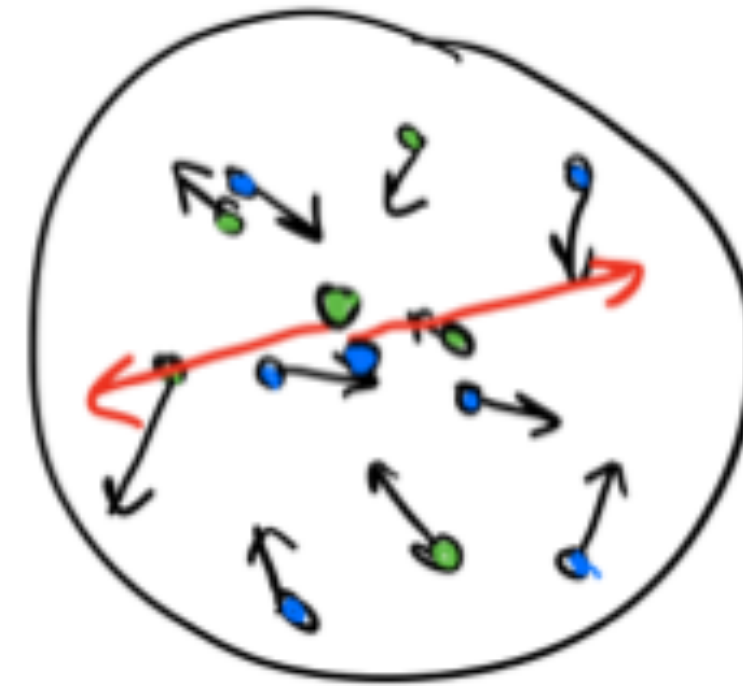


(RG) Resolution Scale  $H = H(\Lambda)$   $\rightarrow$  max. momenta in low-energy wf's  $\sim \Lambda$

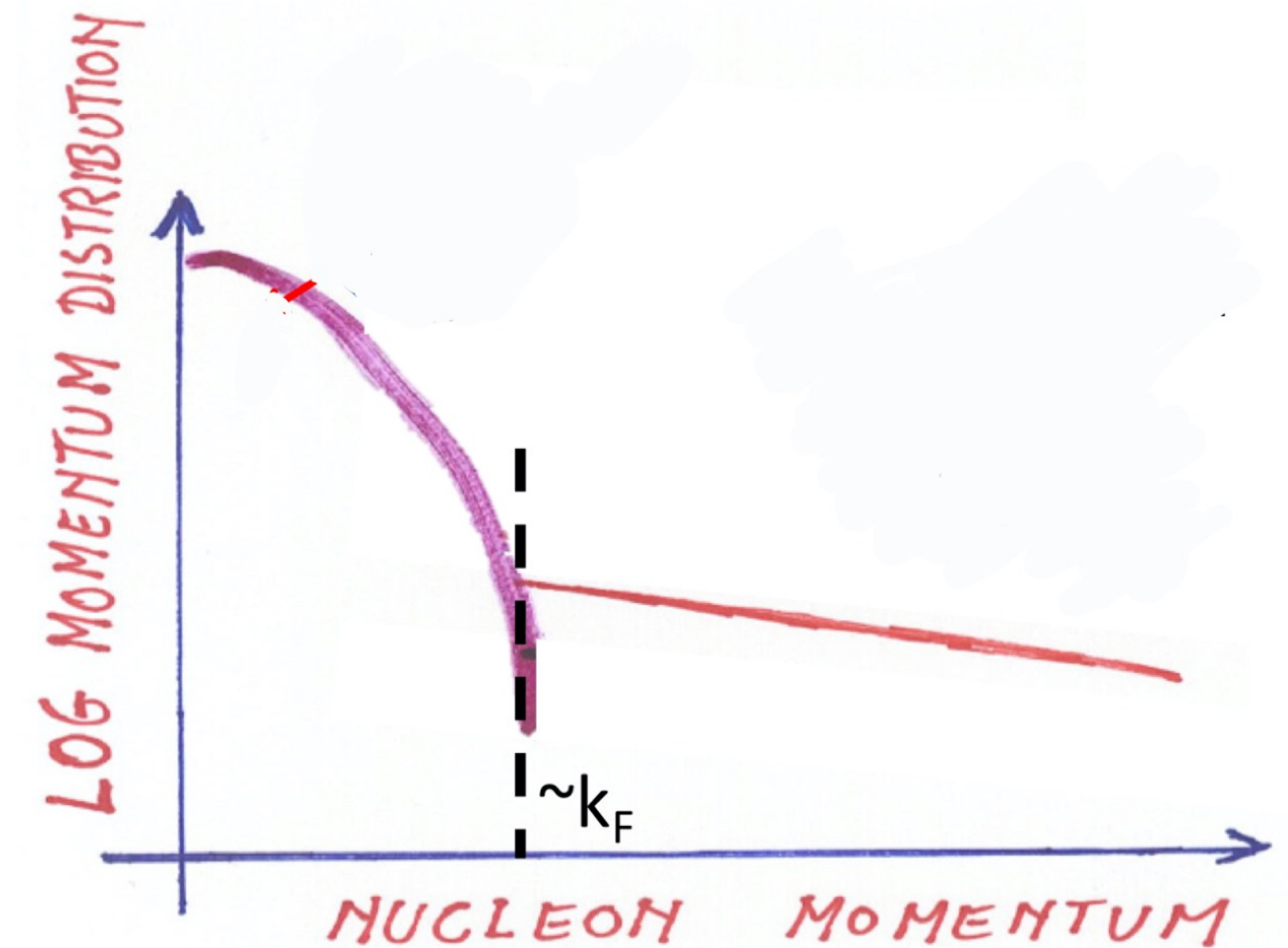
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**High resolution** picture:



correlated pairs w/large  $\sim$  back-to-back momenta

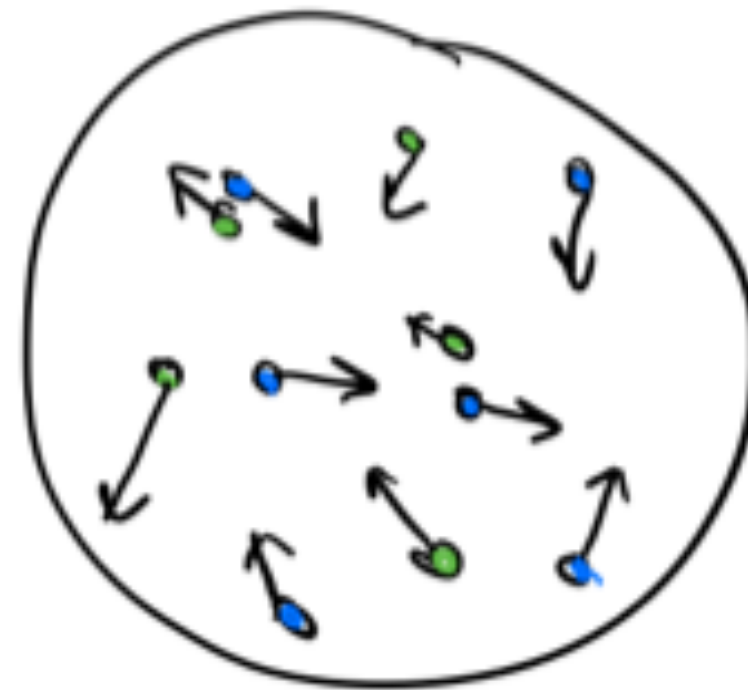


high-k tails ( $k \gg k_F$ ) present

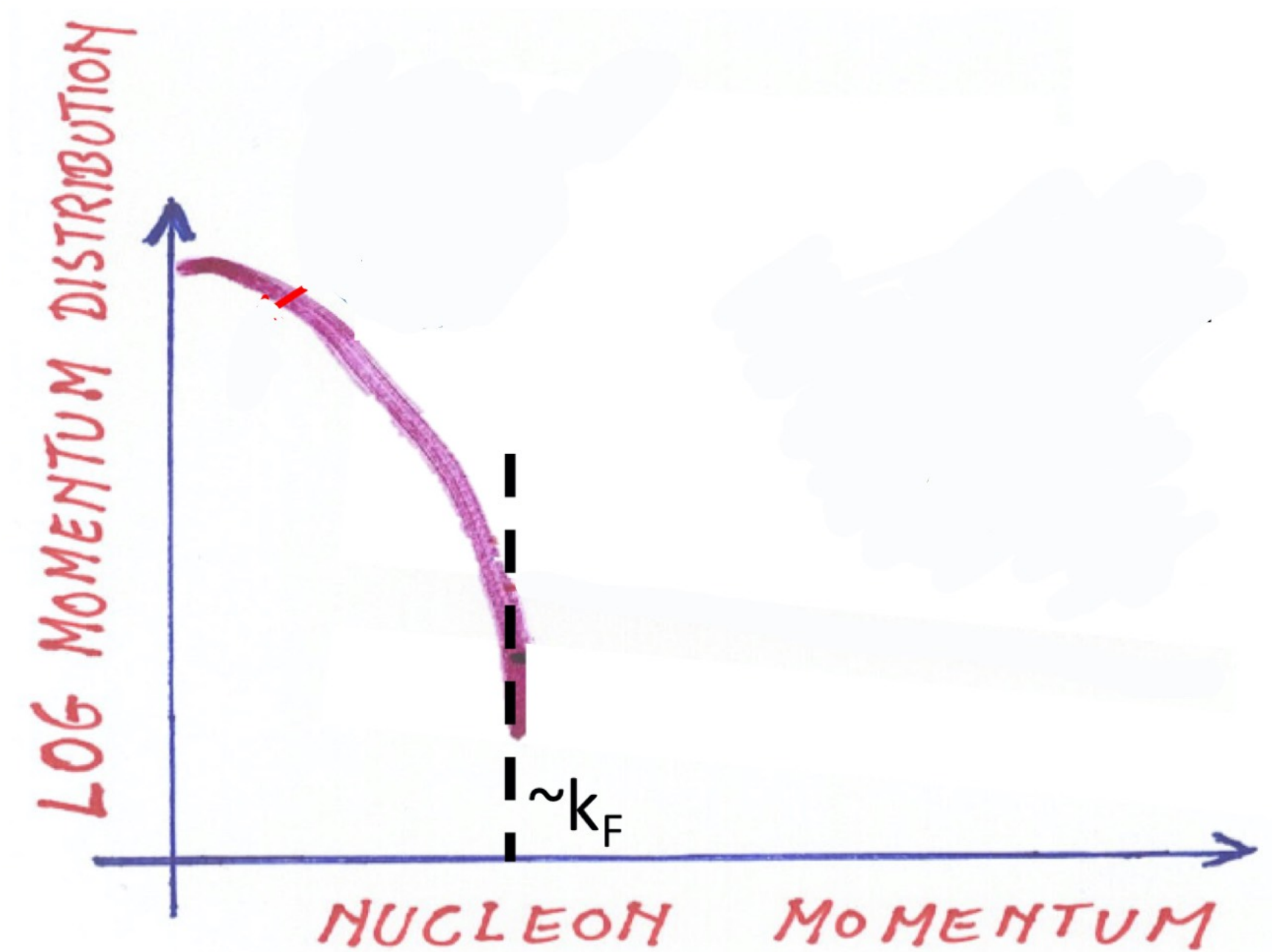
# Low and High Resolution Scale Pictures

(RG) Resolution Scale  $H = H(\Lambda)$   $\rightarrow$  max. momenta in low-energy wf's  $\sim \Lambda$

**Low resolution** picture:



resembles “mean field” picture



no high-k tails ( $k \gg k_F$ )

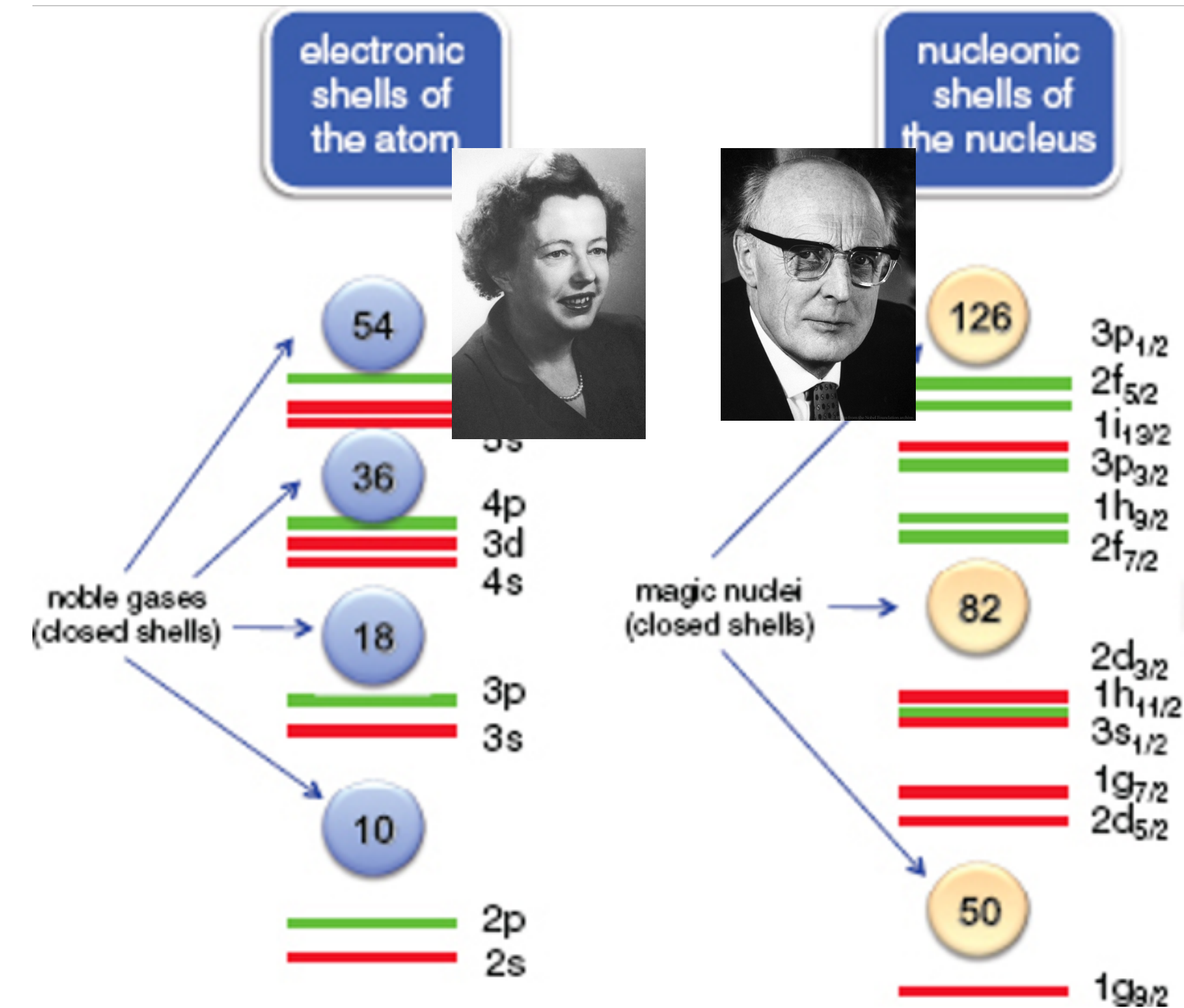
# Low and High Resolution Scale Pictures



Examples of **Low resolution** pictures

# Low and High Resolution Scale Pictures

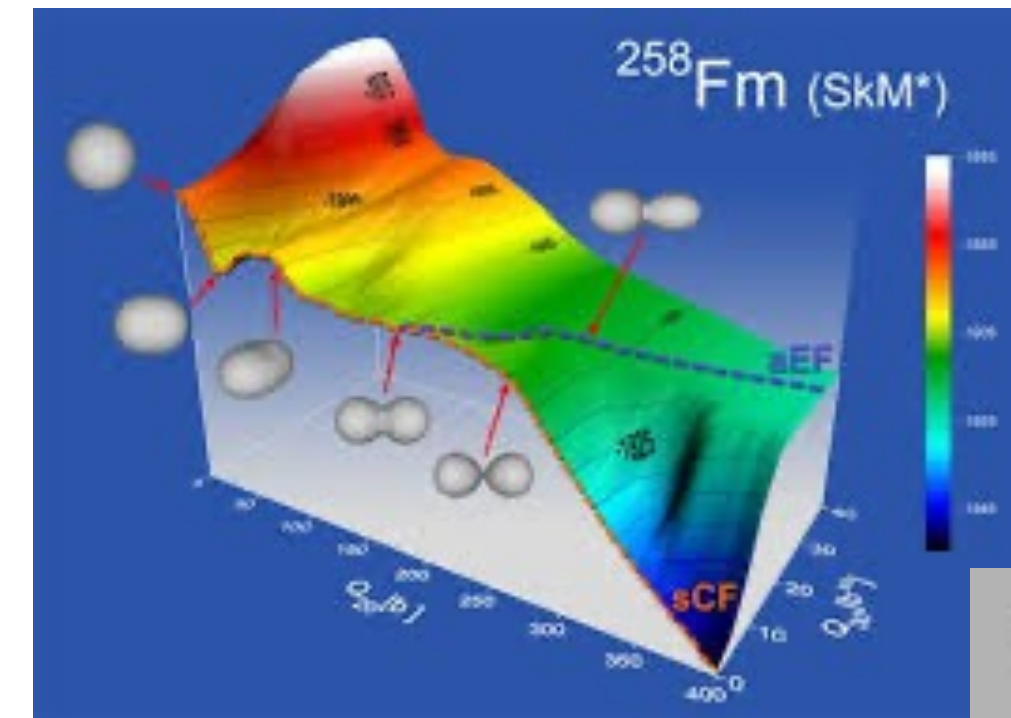
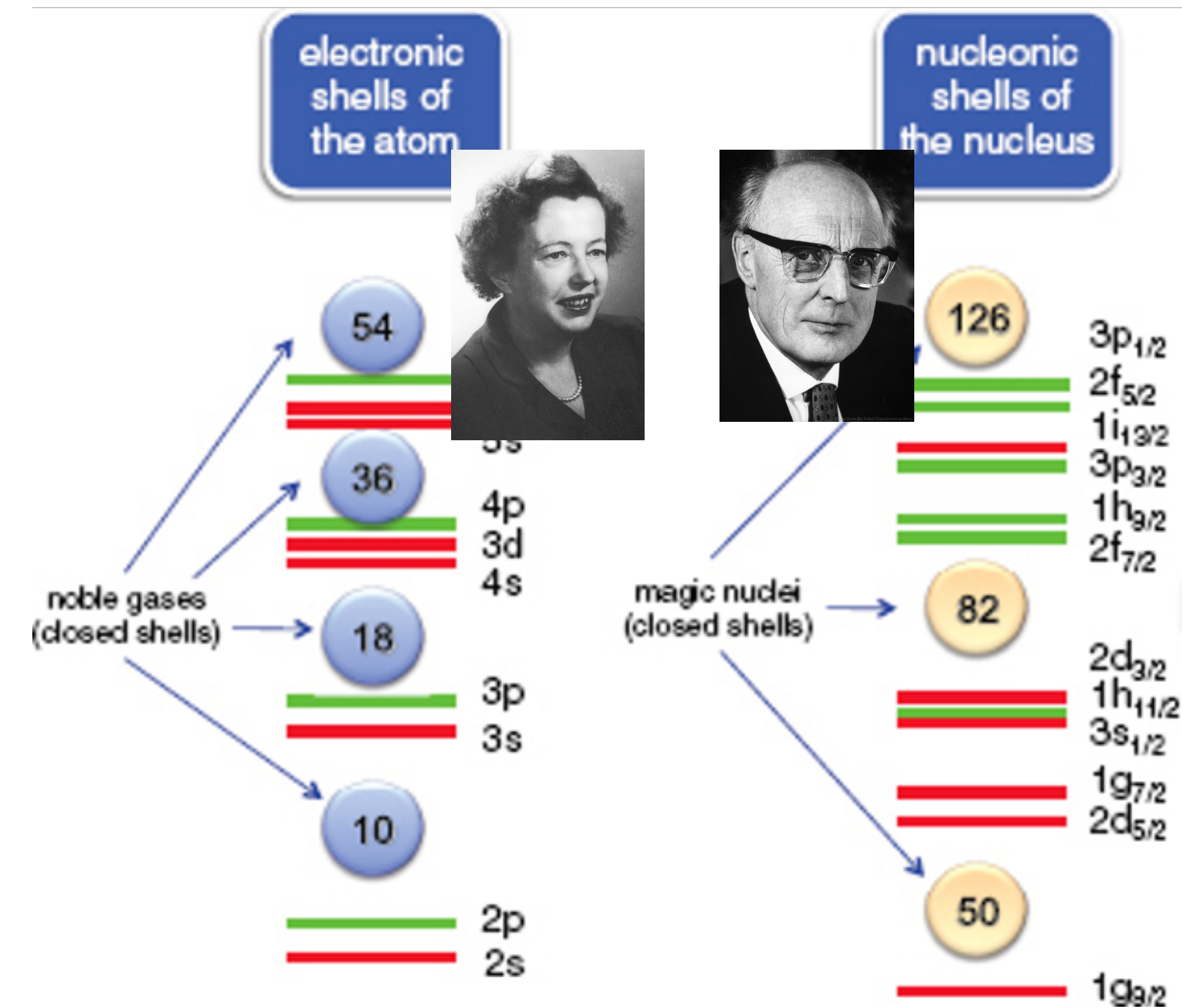
## Examples of **Low resolution** pictures



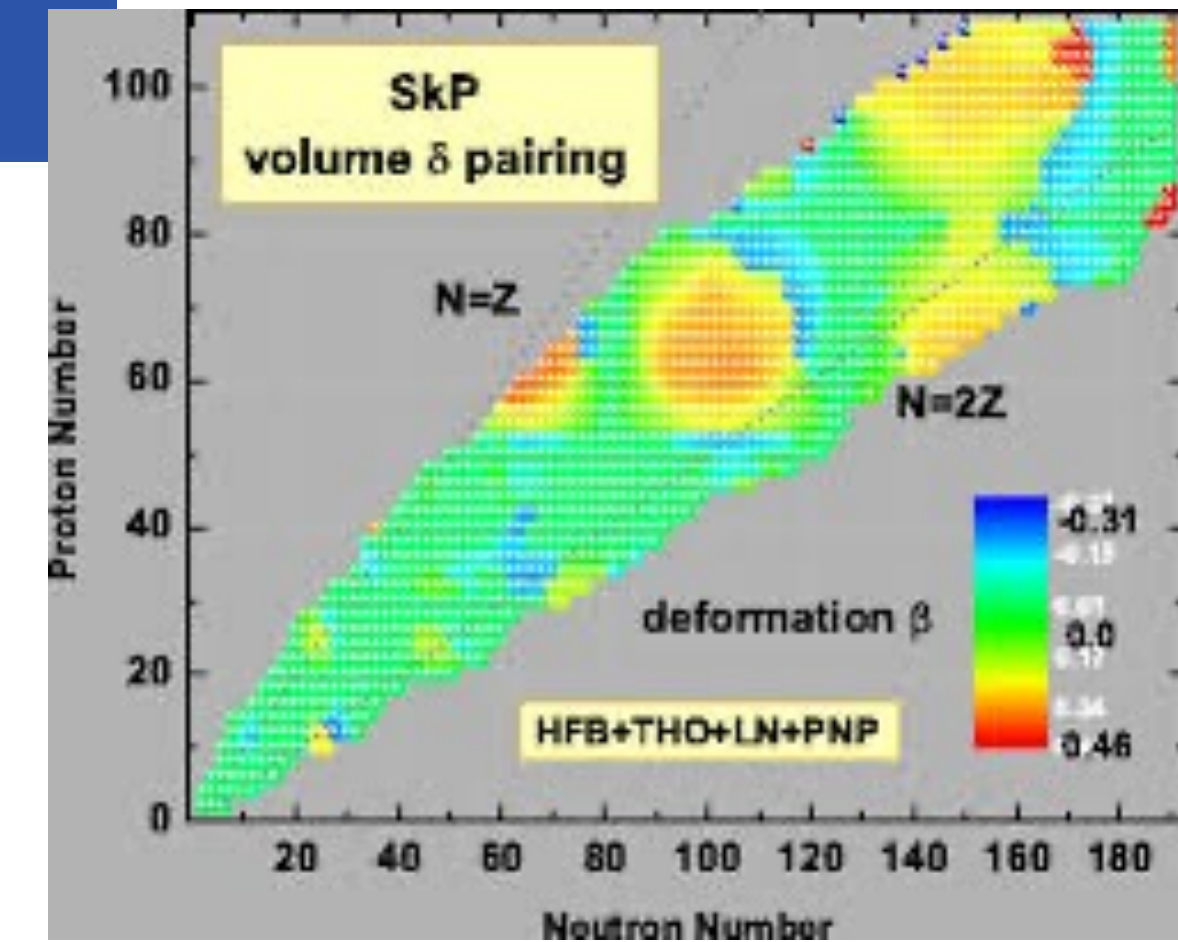
Nuclear Shell model

# Low and High Resolution Scale Pictures

## Examples of **Low resolution** pictures



$$\begin{aligned}
 V_{\text{Skyrme}} = & t_0(1 + x_0 \hat{P}_\sigma) \delta(\vec{r}) \\
 & + \frac{1}{2} t_1(1 + x_1 \hat{P}_\sigma) [\delta(\vec{r}) \vec{k}^2 + \vec{k}'^2 \delta(\vec{r})] \\
 & + t_2(1 + x_2 \hat{P}_\sigma) \vec{k}' \cdot \delta(\vec{r}) \vec{k} \\
 & + \frac{1}{6} t_3(1 + x_3 \hat{P}_\sigma) \rho^\alpha(\vec{R}) \delta(\vec{r}) \\
 & + i t_4 \vec{k}' \cdot \delta(\vec{r}) (\vec{\sigma}_i + \vec{\sigma}_j) \vec{k},
 \end{aligned}$$



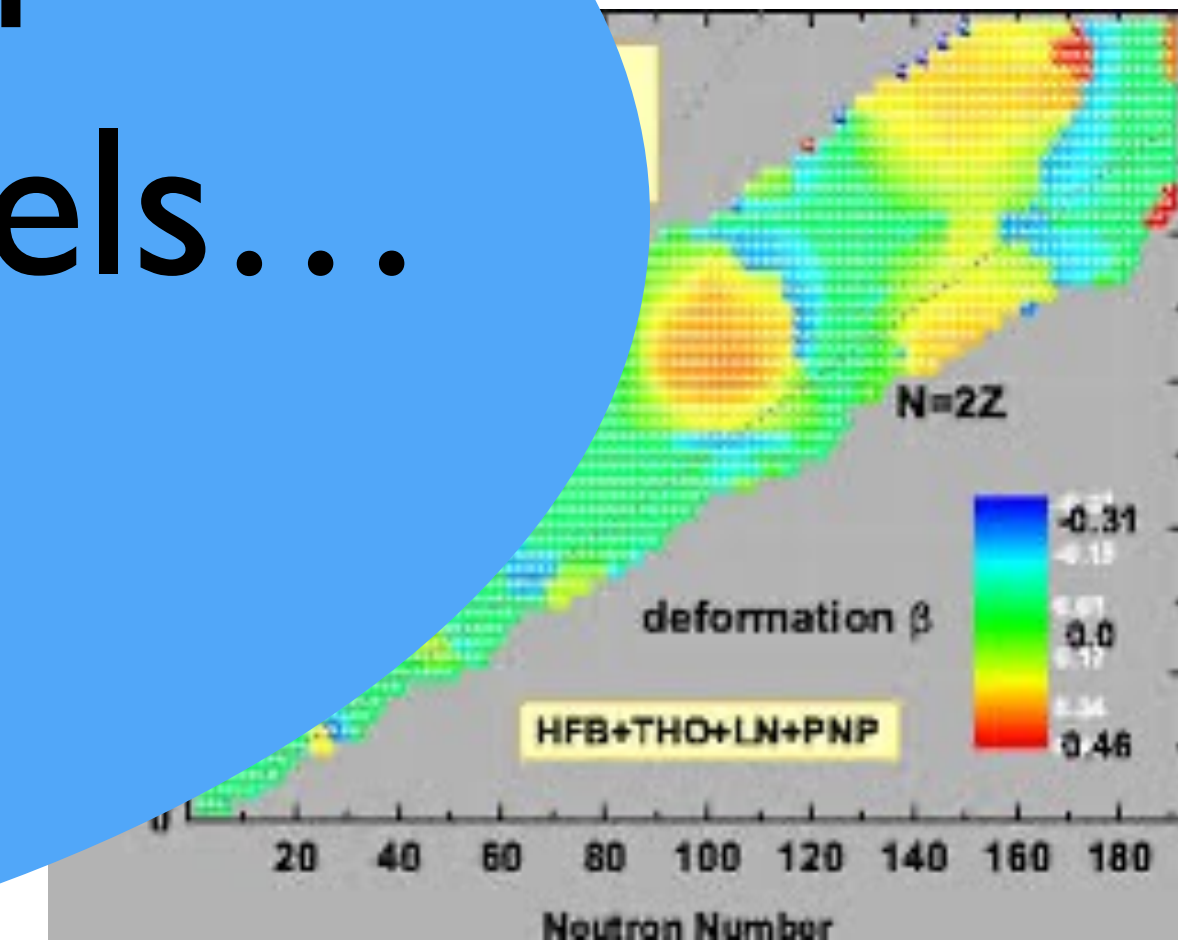
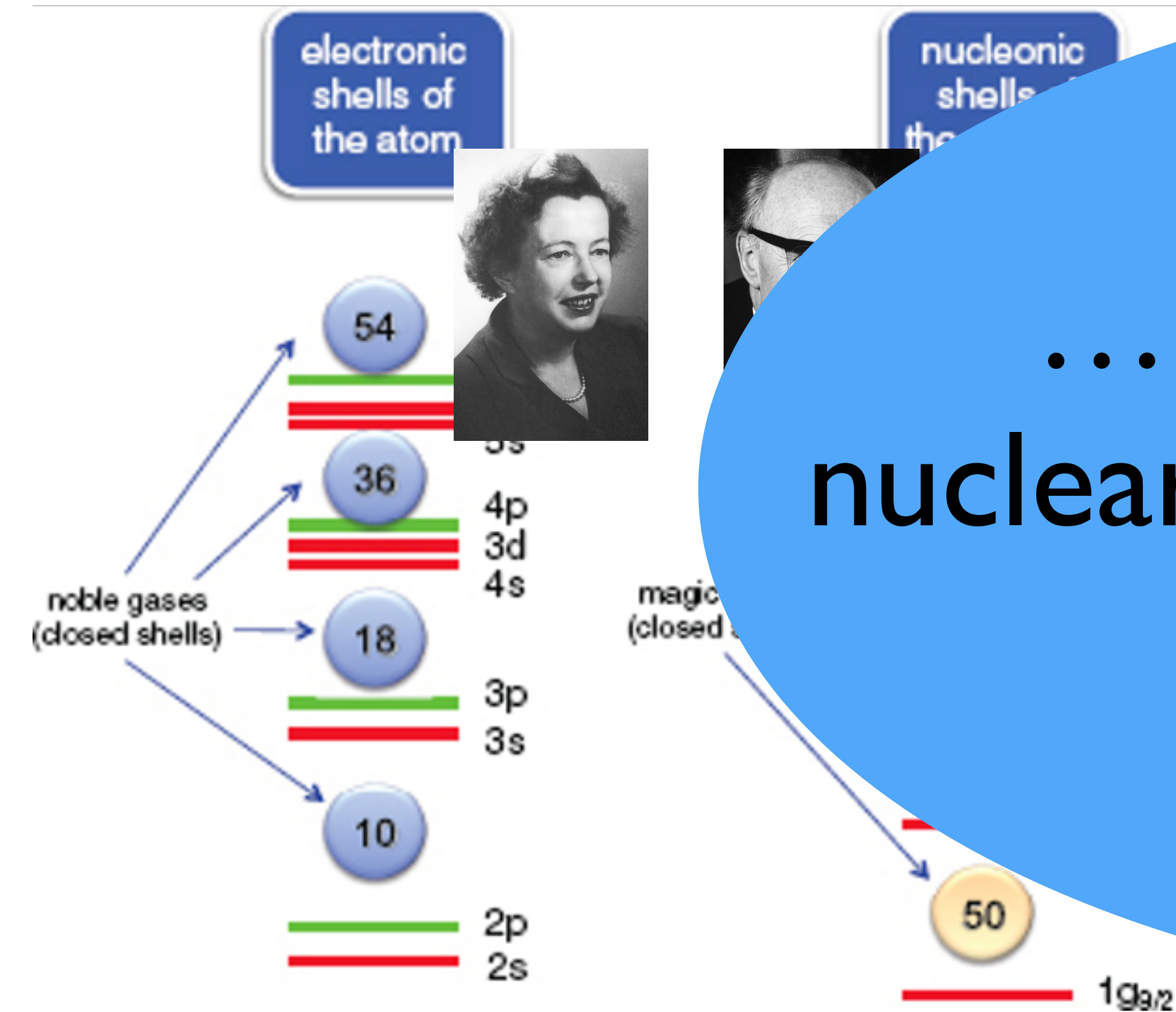
Nuclear Shell model

Nuclear density functional

# Low and High Resolution Scale Pictures

Examples of **Low resolution** pictures

...Pretty much all our nuclear **structure** models...



Nuclear Shell model

Nuclear density functional



# Short History: High Resolution Scale Picture



All nuclear structure models  $\Leftrightarrow$  **Low resolution** pictures

How did the high resolution picture arise?

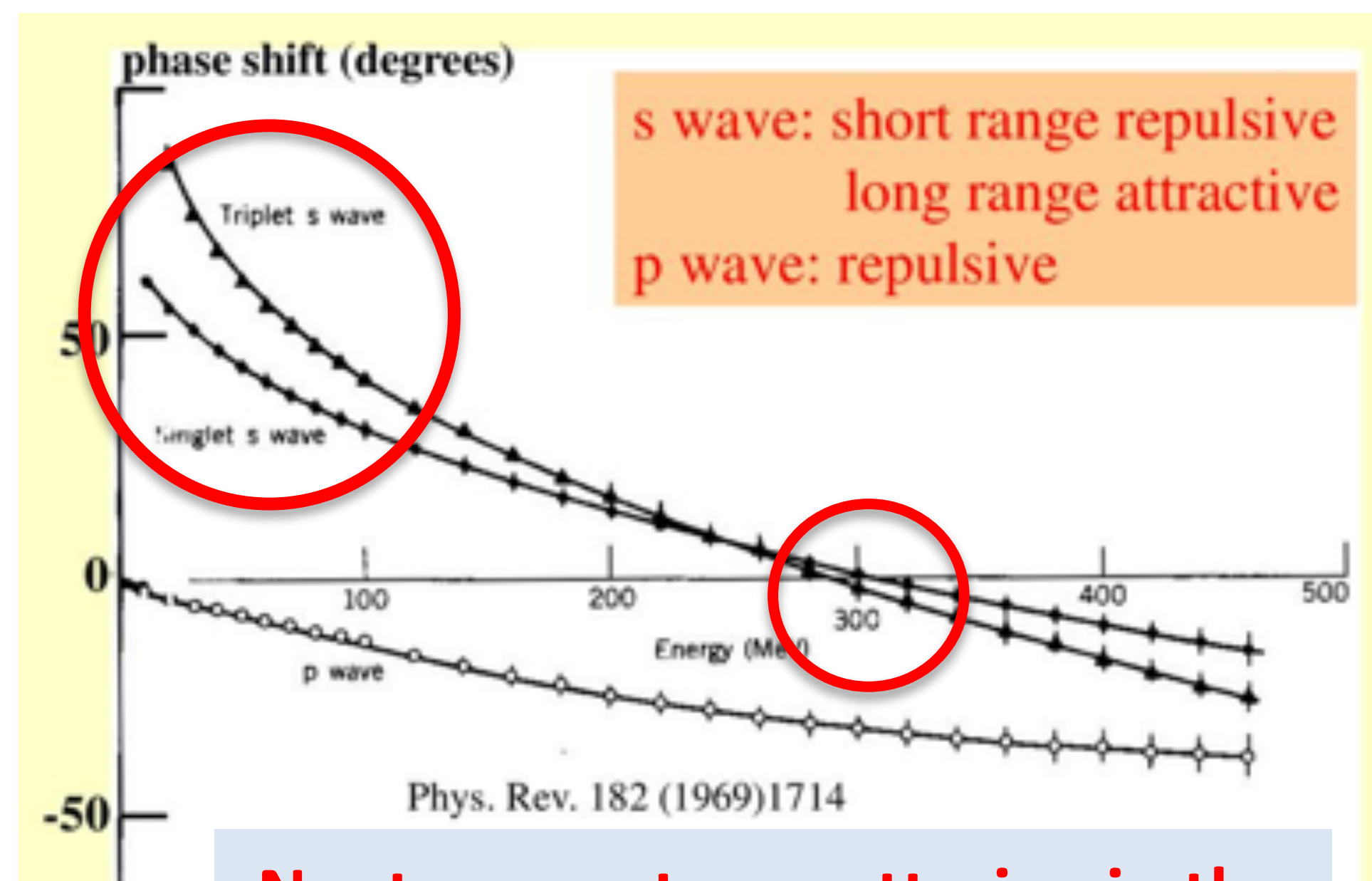
exhibit A: NN scattering (1950s-60s)

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**Neutron-proton scattering in the S-waves changes sign**

**IF** you insist on a local  $V(r)$ :

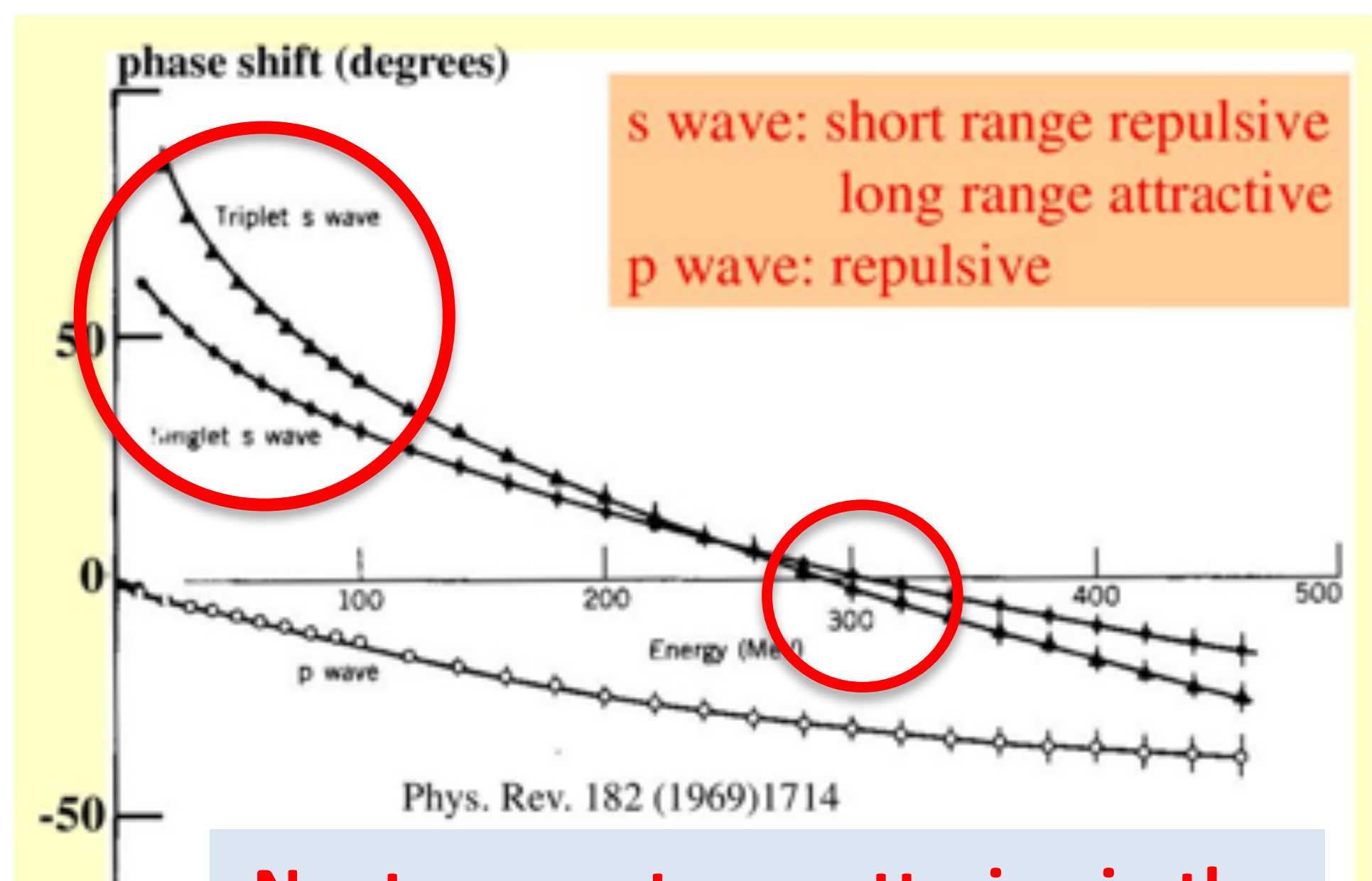
strong short-range repulsive core needed to get s-wave sign change

# Short History: High Resolution Scale Picture

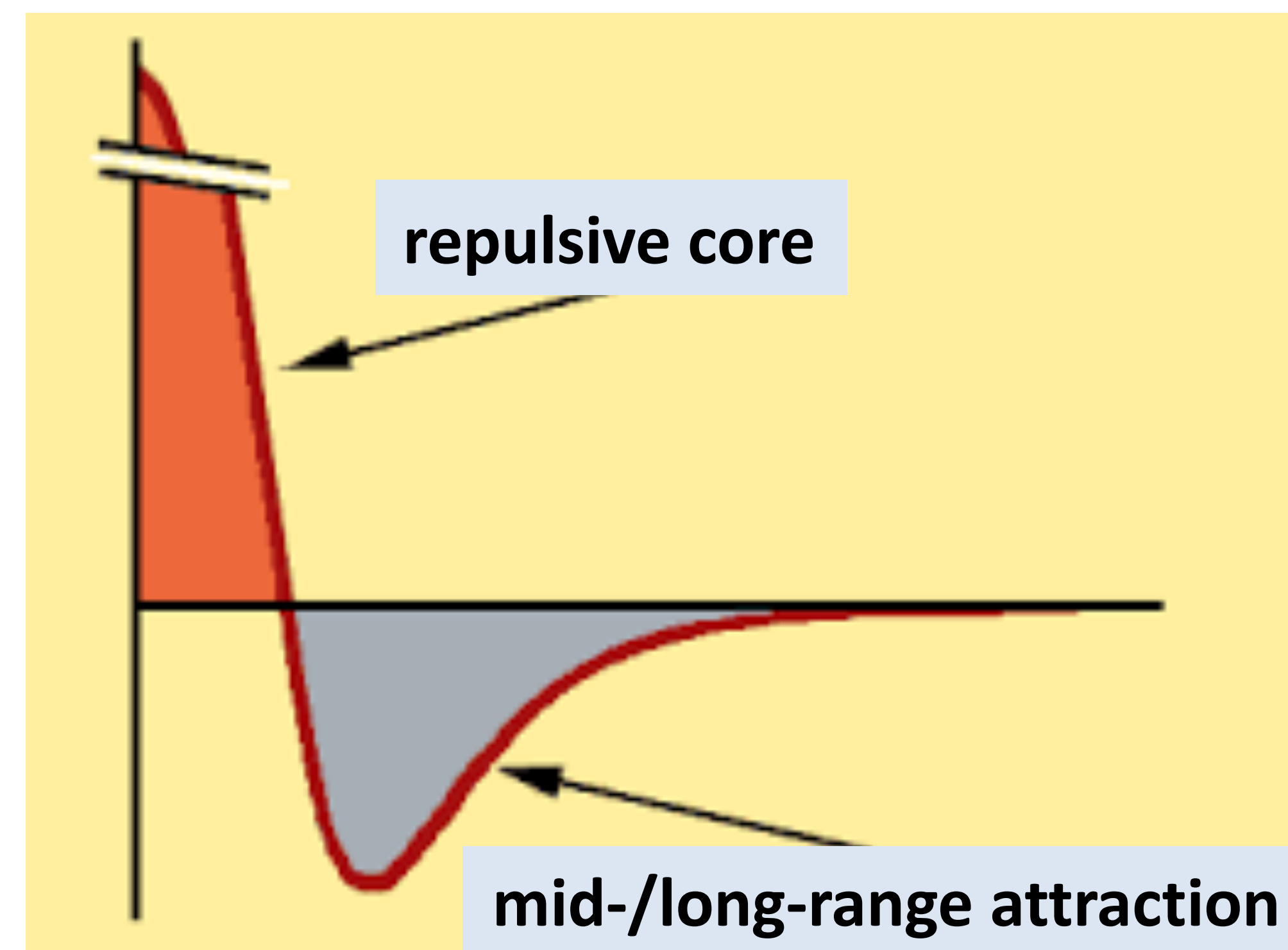
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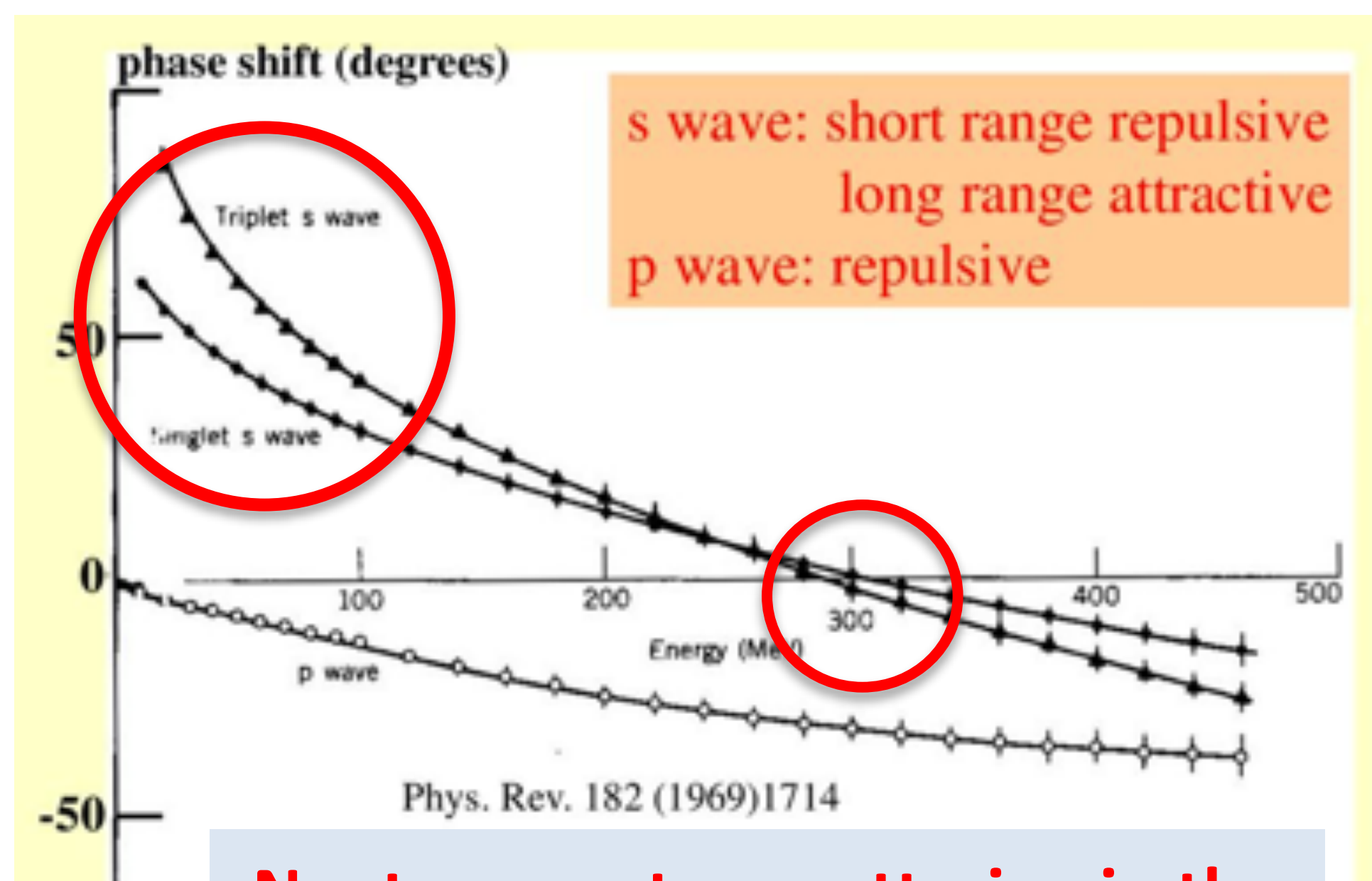


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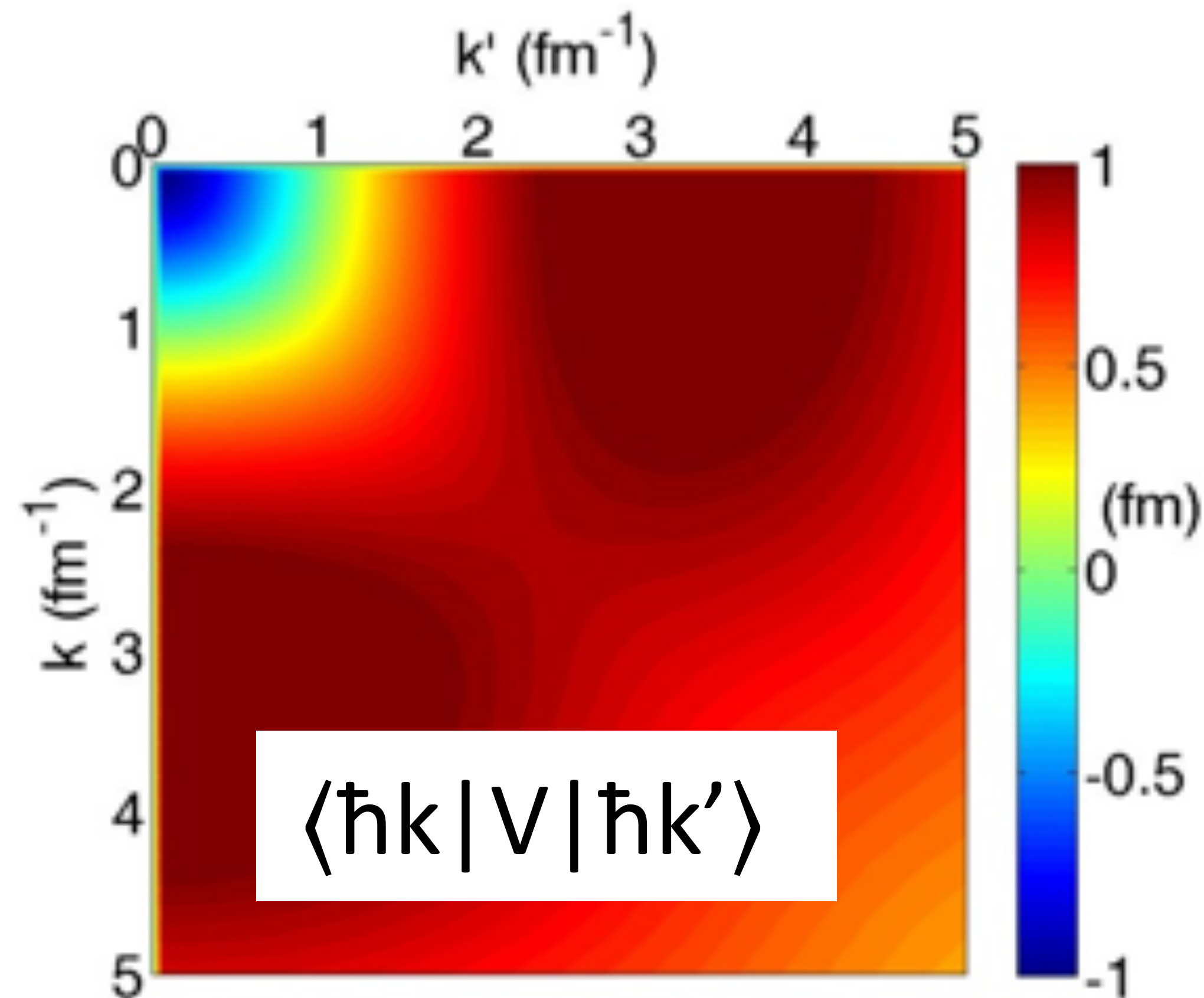
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## exhibit B: Brueckner 1955 paper

PHYSICAL REVIEW

VOLUME 98, NUMBER 5

JUNE 1, 1955

### High-Energy Reactions and the Evidence for Correlations in the Nuclear Ground-State Wave Function\*

K. A. BRUECKNER, R. J. EDEN,<sup>†</sup> AND N. C. FRANCIS

*Indiana University, Bloomington, Indiana*

(Received January 13, 1955)

High-energy nuclear reactions which depend strongly on nucleon position correlations in the nuclear ground state are analyzed and shown to give evidence for the existence of marked correlation effects. The following high-energy experiments are considered: nuclear photoeffect, meson absorption in nuclei, deuteron pickup, proton-proton scattering in a nucleus, and meson production in proton-nucleus collisions. The corresponding cross sections depend on a nucleon momentum distribution which can be represented at high energies by a single function giving reasonable agreement with all the experiments considered. This momentum distribution differs substantially from that for the shell model of the nucleus and thus provides strong evidence for correlation in the nuclear ground-state wave function.

The transformation methods developed in previous papers are used to provide a unified theory of the above five processes. The momentum distribution predicted by this theory is estimated by two methods each of which gives close agreement with the experimentally determined function in the relevant energy ranges.

# Short History: High Resolution Scale Picture

## exhibit B: Brueckner 1955 paper

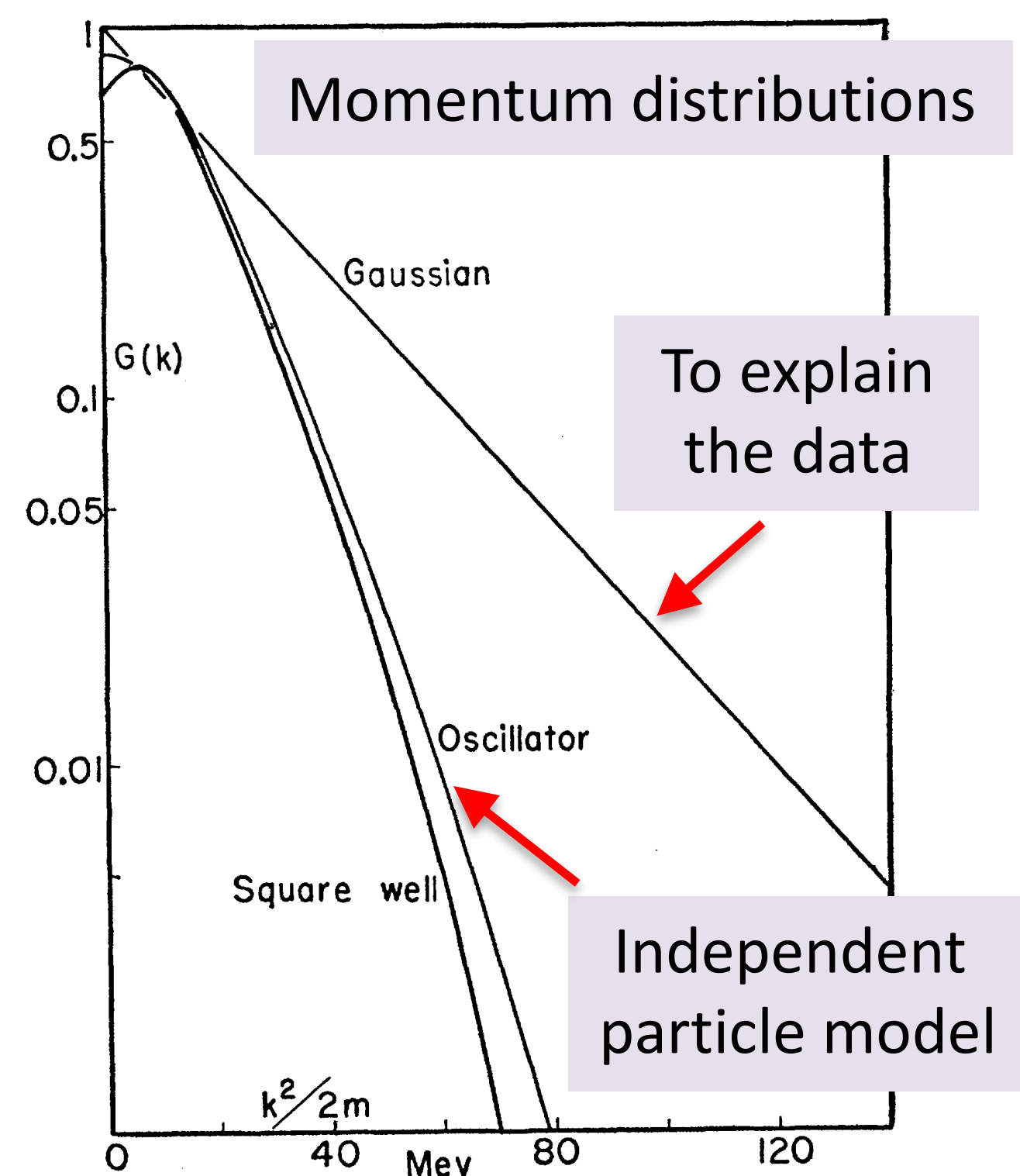


FIG. 1. Momentum distribution  $G(k)$  of 8 neutrons and 8 protons in the independent-particle states of a square well with infinite walls and of a harmonic oscillator well. For comparison the Gaussian distribution of Eq. (3) is also given.

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The transformation methods developed in previous papers are used to provide a unified theory of the above five processes. The momentum distribution predicted by this theory is estimated by two methods each of which gives close agreement with the experimentally determined function in the relevant energy ranges.

“Consequently it follows that the usual assumptions of the shell-model theory of the nucleus, that the particles move independently in a uniform potential, cannot be other than very approximately correct.”

# Short History: High Resolution Scale Picture

## exhibit B: Brueckner 1955 paper

### Key configurations in Brueckner's analysis:

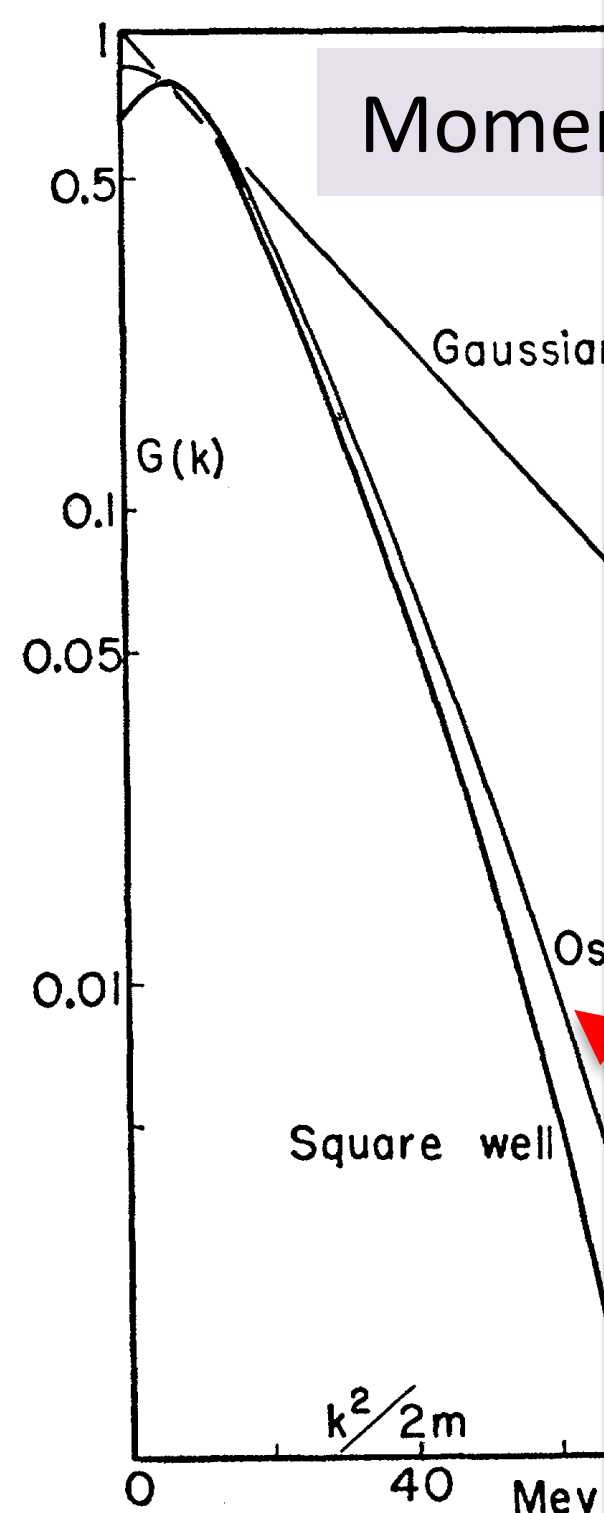
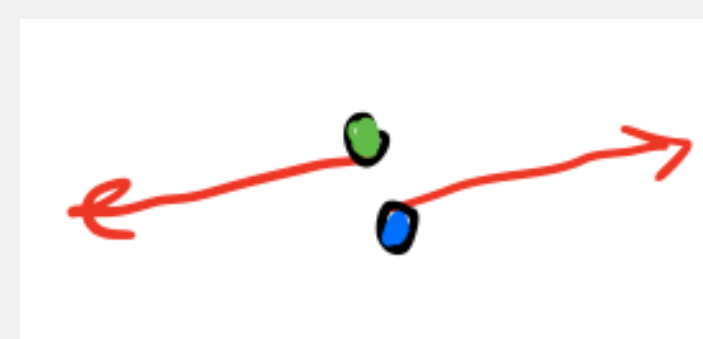
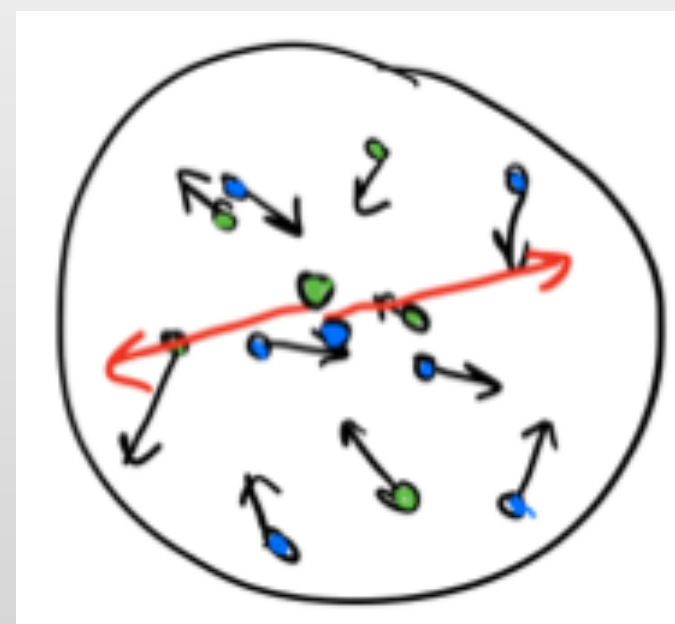


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This is a *short-range correlation* or SRC

$$\Psi(k_1, k_2, k_3, \dots, k_A) \sim \phi(k_1 - k_2) \chi_A(K_{12}, k_3, \dots, k_A)$$

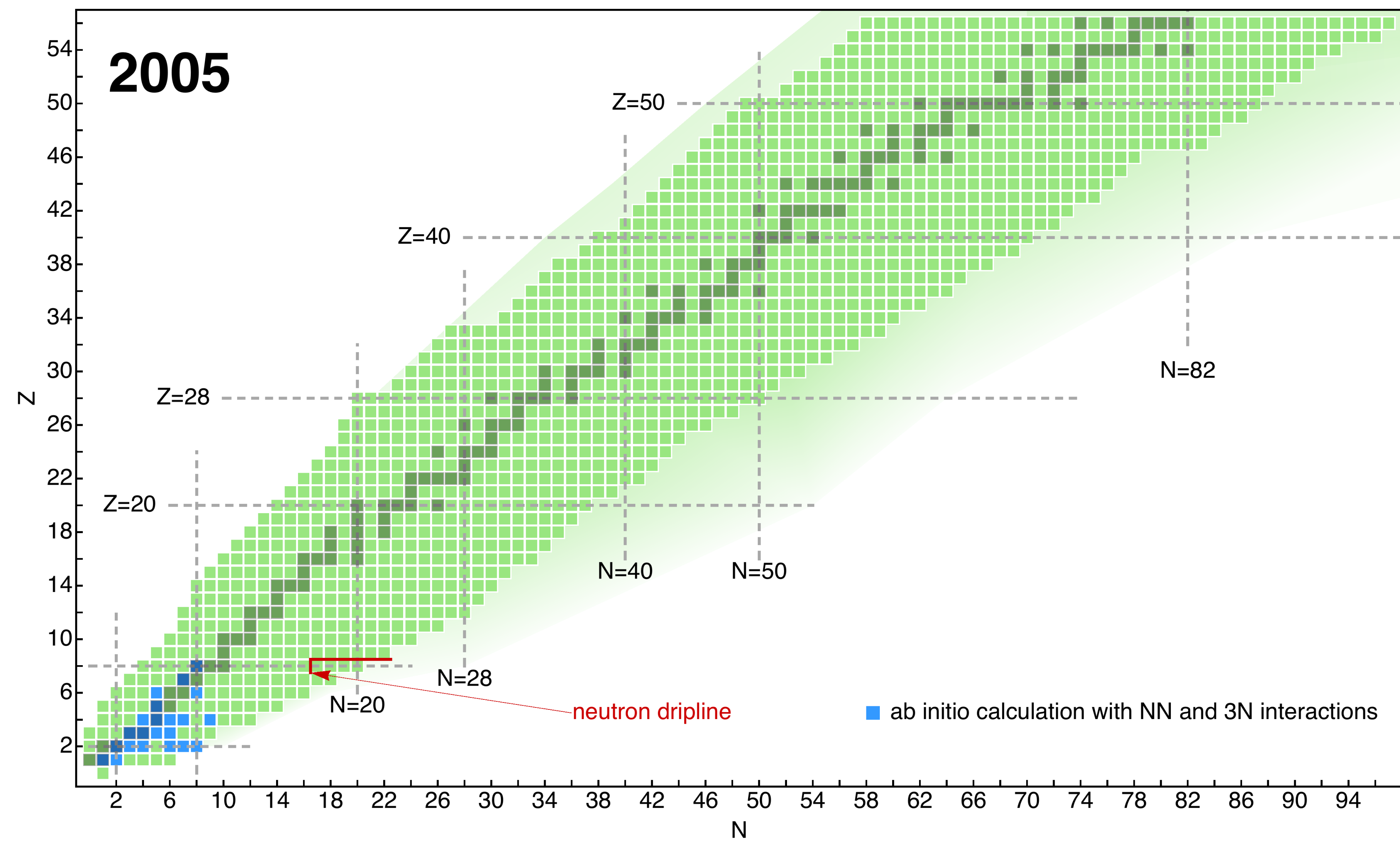
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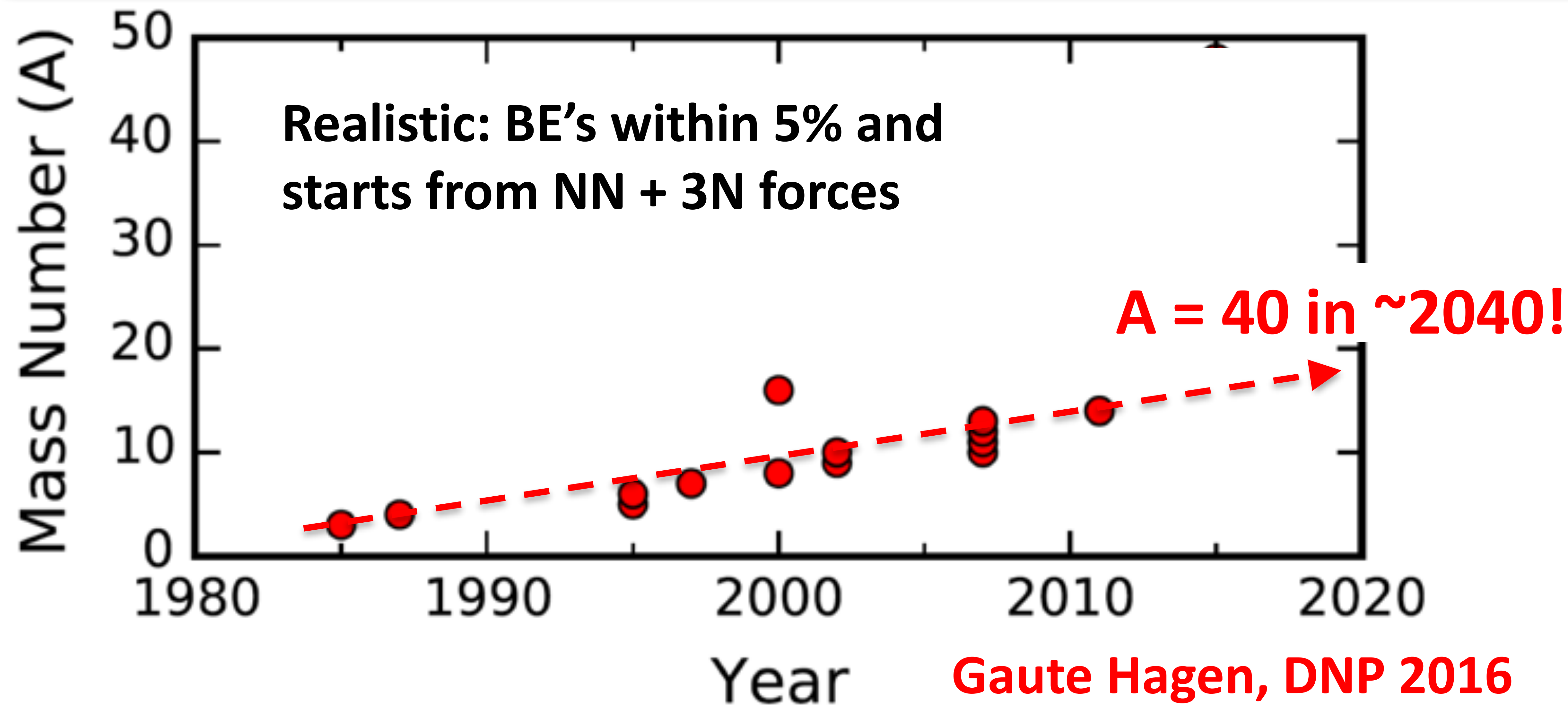
-model theory  
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# Ab Initio with high-resolution NN + NNN

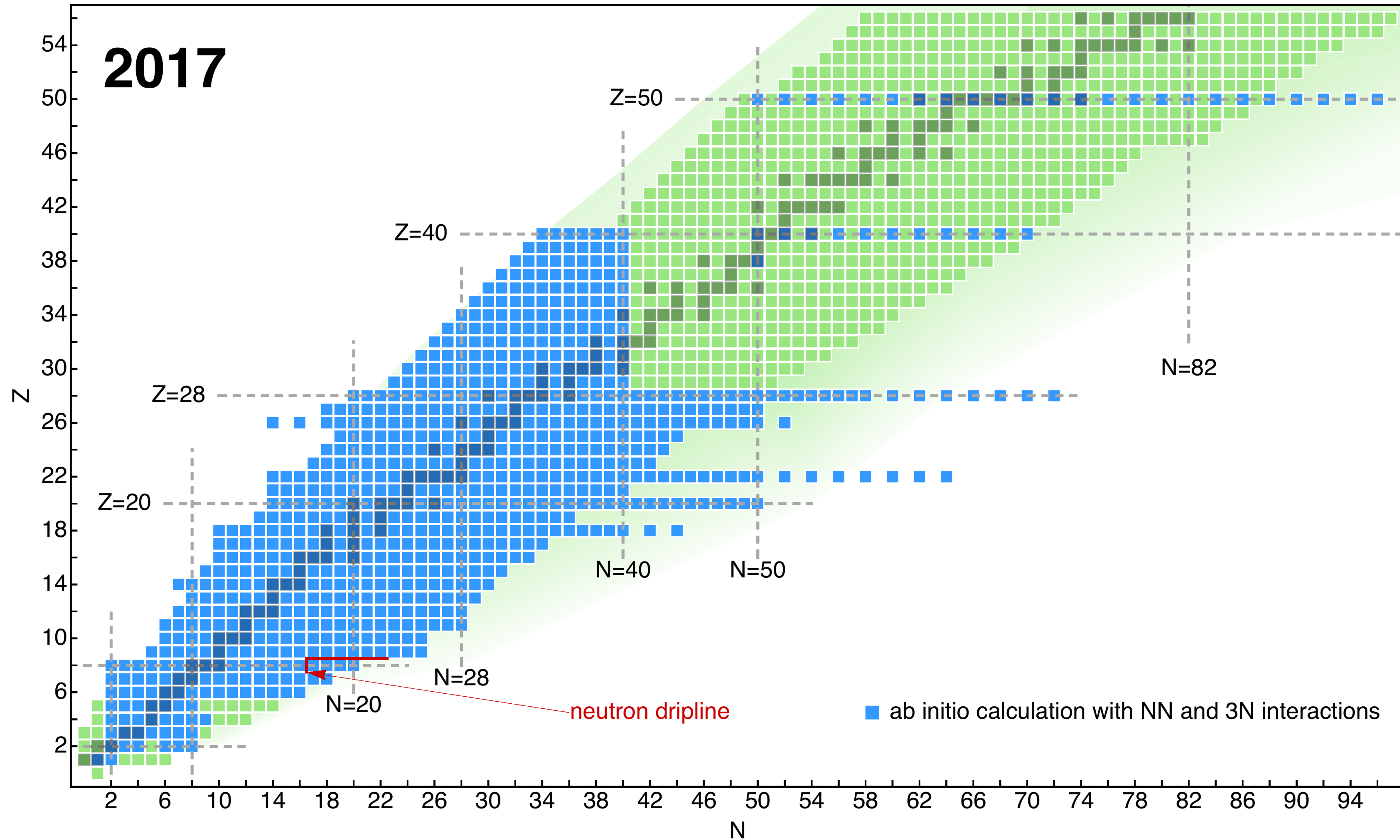




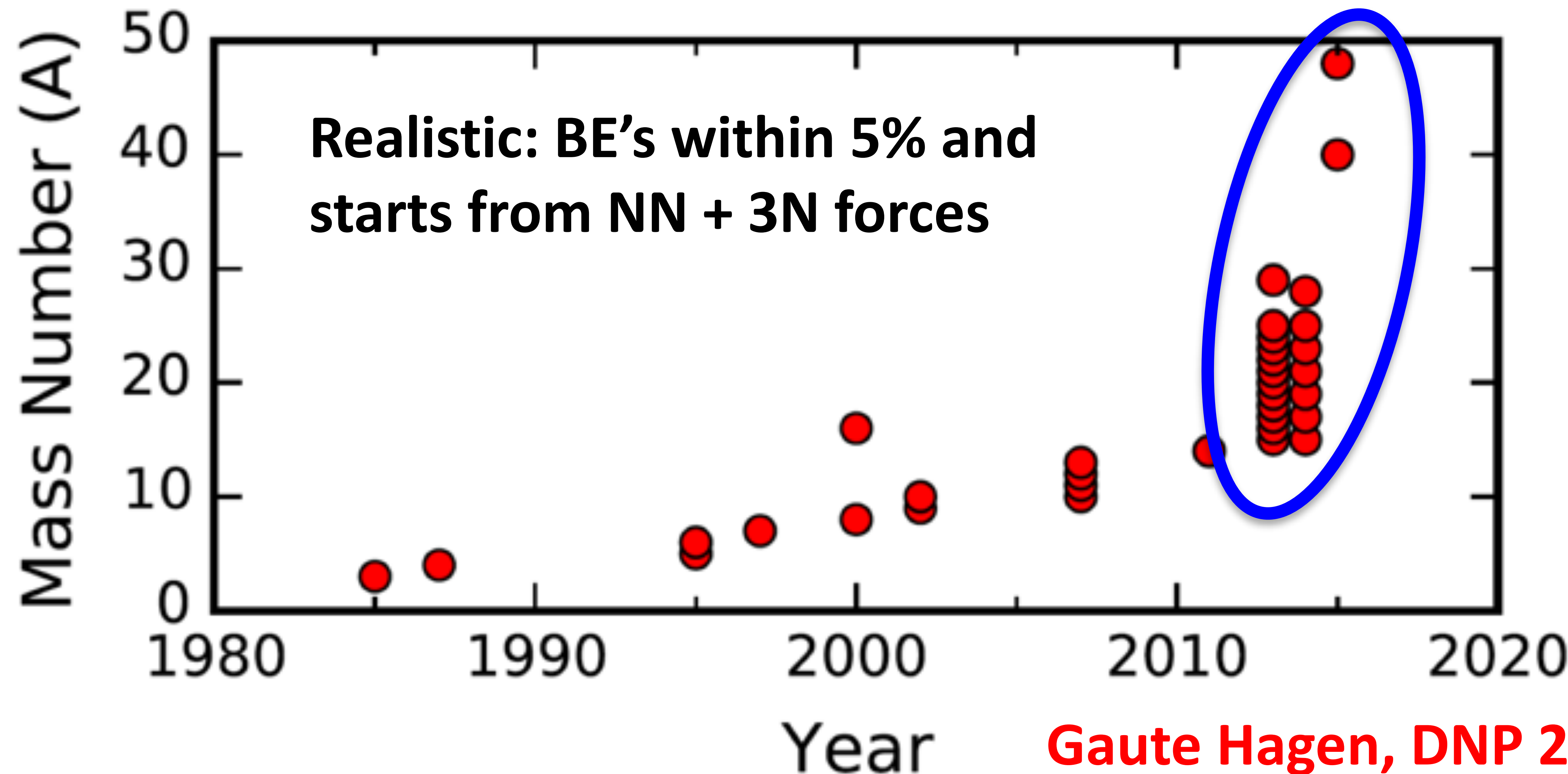
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# Ab Initio with low-resolution NN + NNN

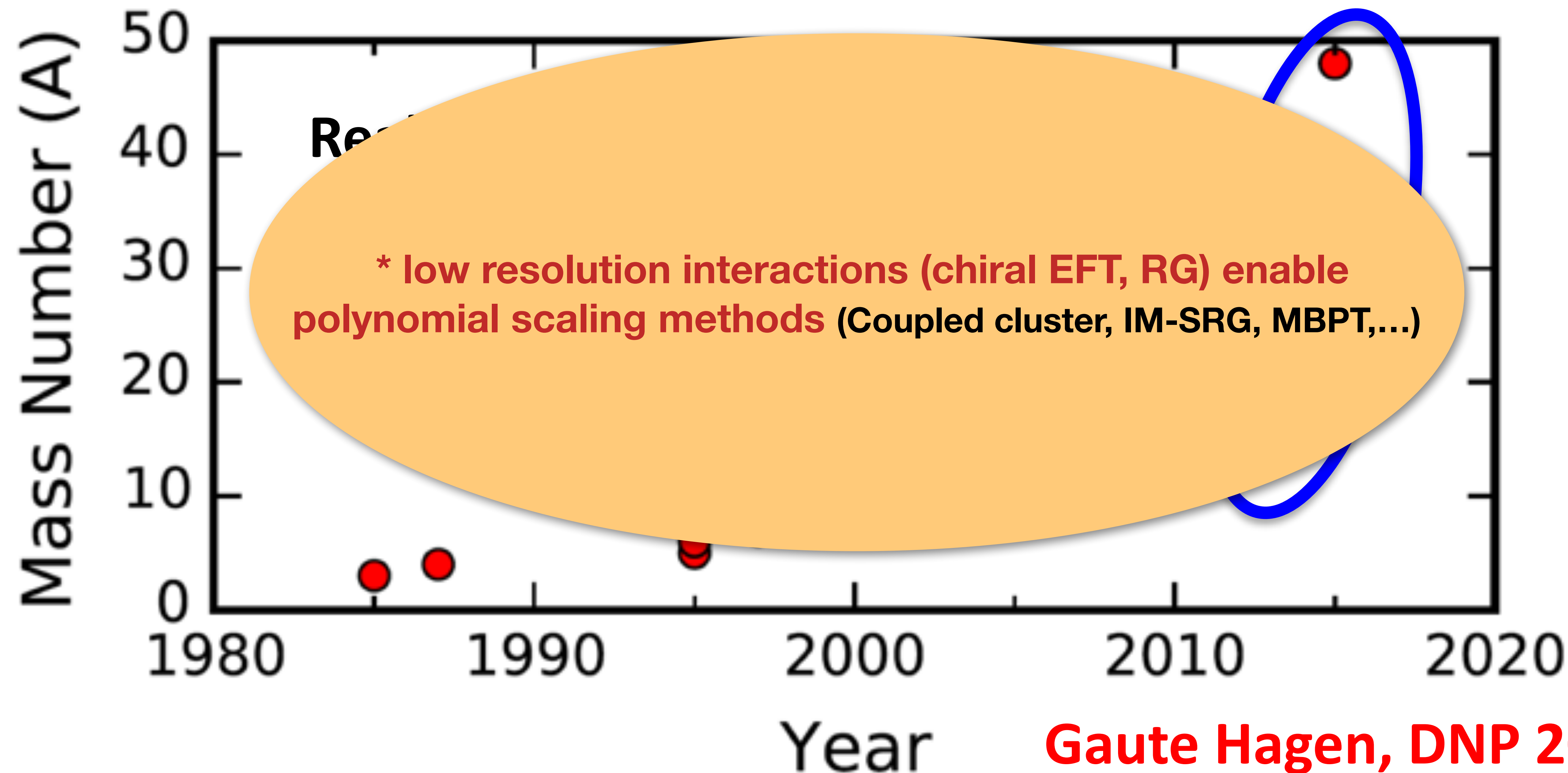


# Ab Initio with low-resolution NN + NNN



Gaute Hagen, DNP 2016

# Ab Initio with low-resolution NN + NNN



# Modern SRC Phenomenology (2000's - present)



Experiments at BNL and JLab to detect knocked-out nucleons from an SRC *pair*

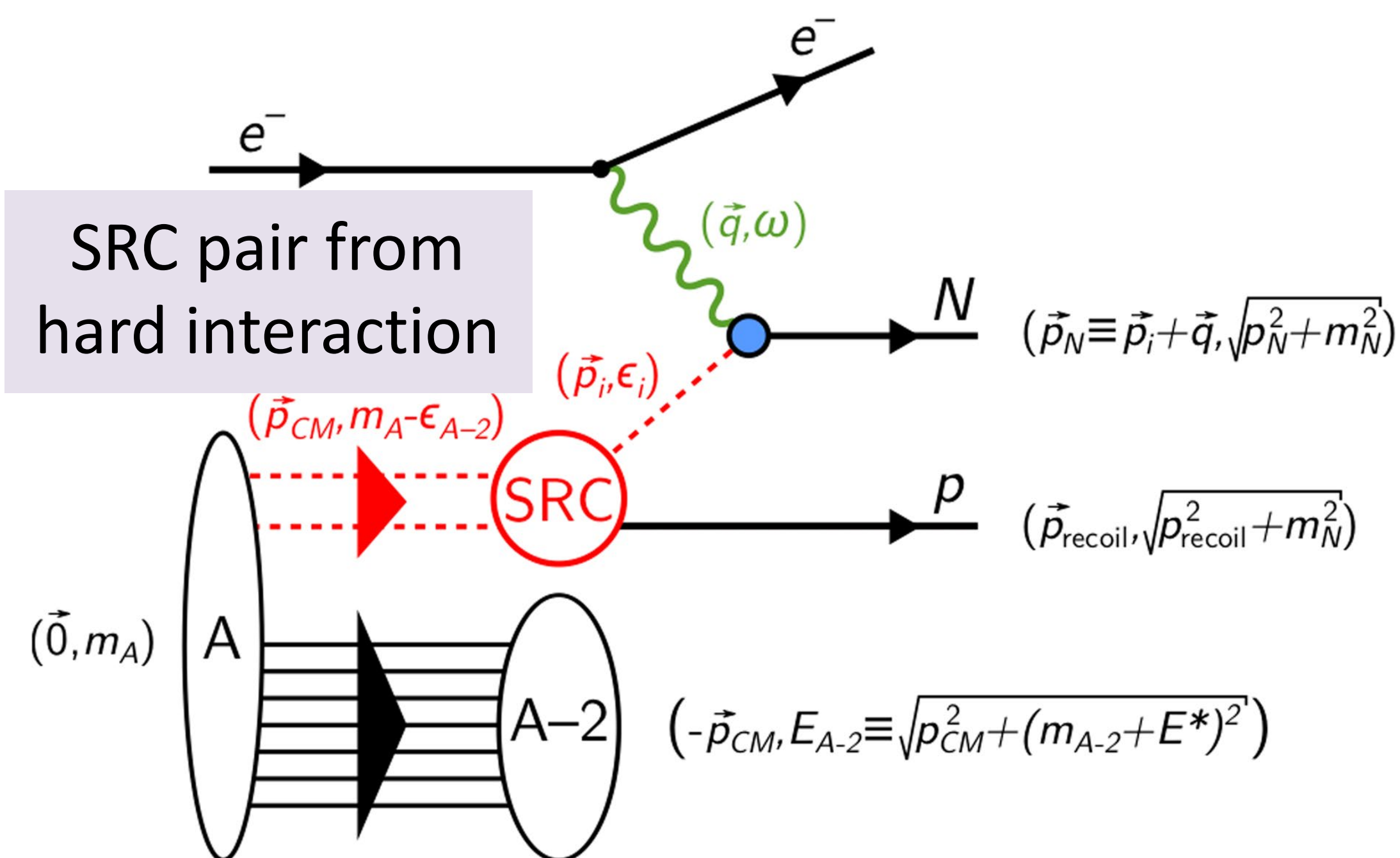
Breakup the pair =>  
Detect **both** nucleons =>  
Reconstruct 'initial' state

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Kinematics chosen to minimize ambiguities from MECs, FSI, etc.



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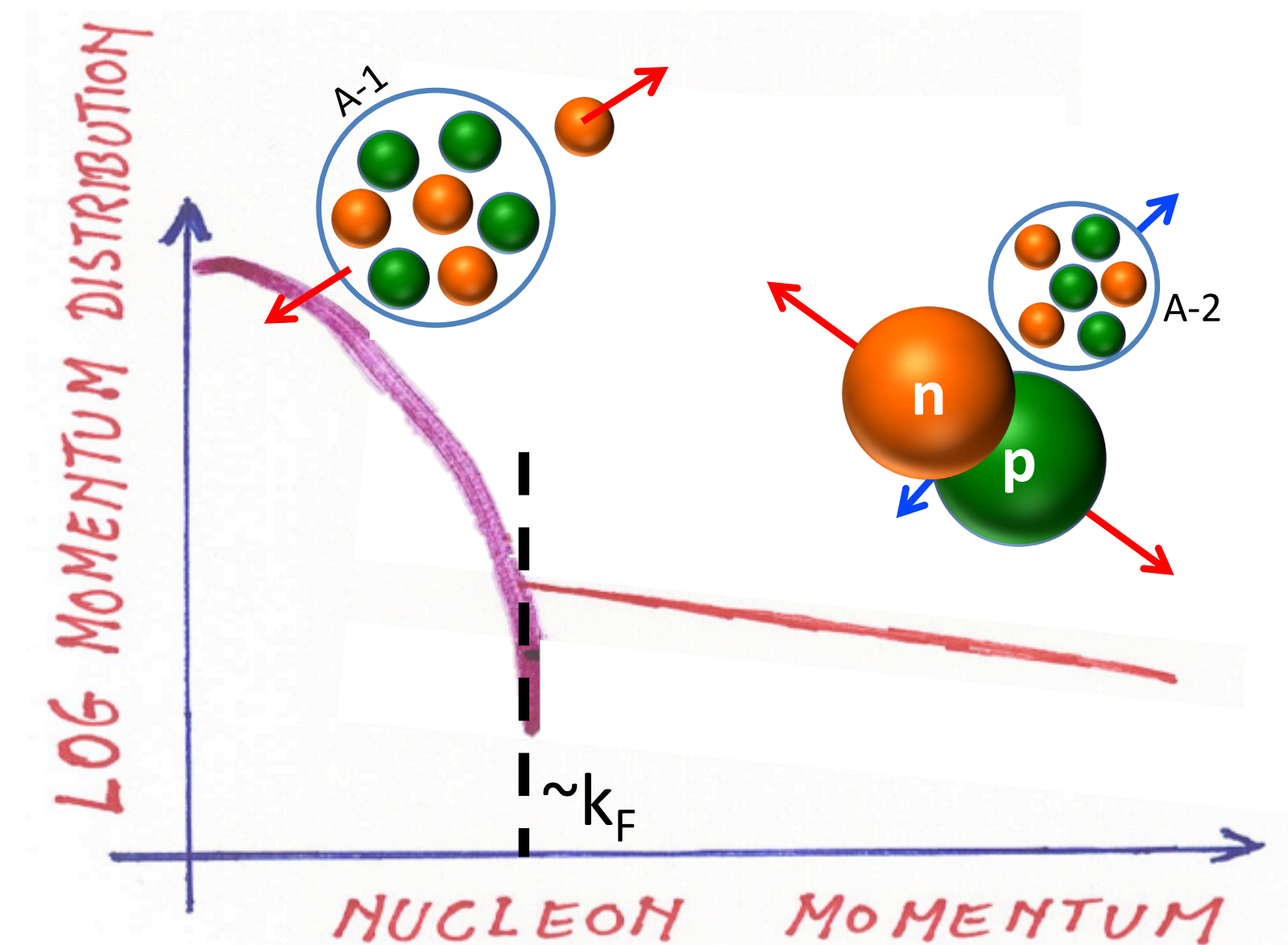
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Interpretation (high resolution picture)

2 regions of momenta in nuclei

~ 20% of nucleons in SRC pairs

~ 70% of KE from SRC pairs



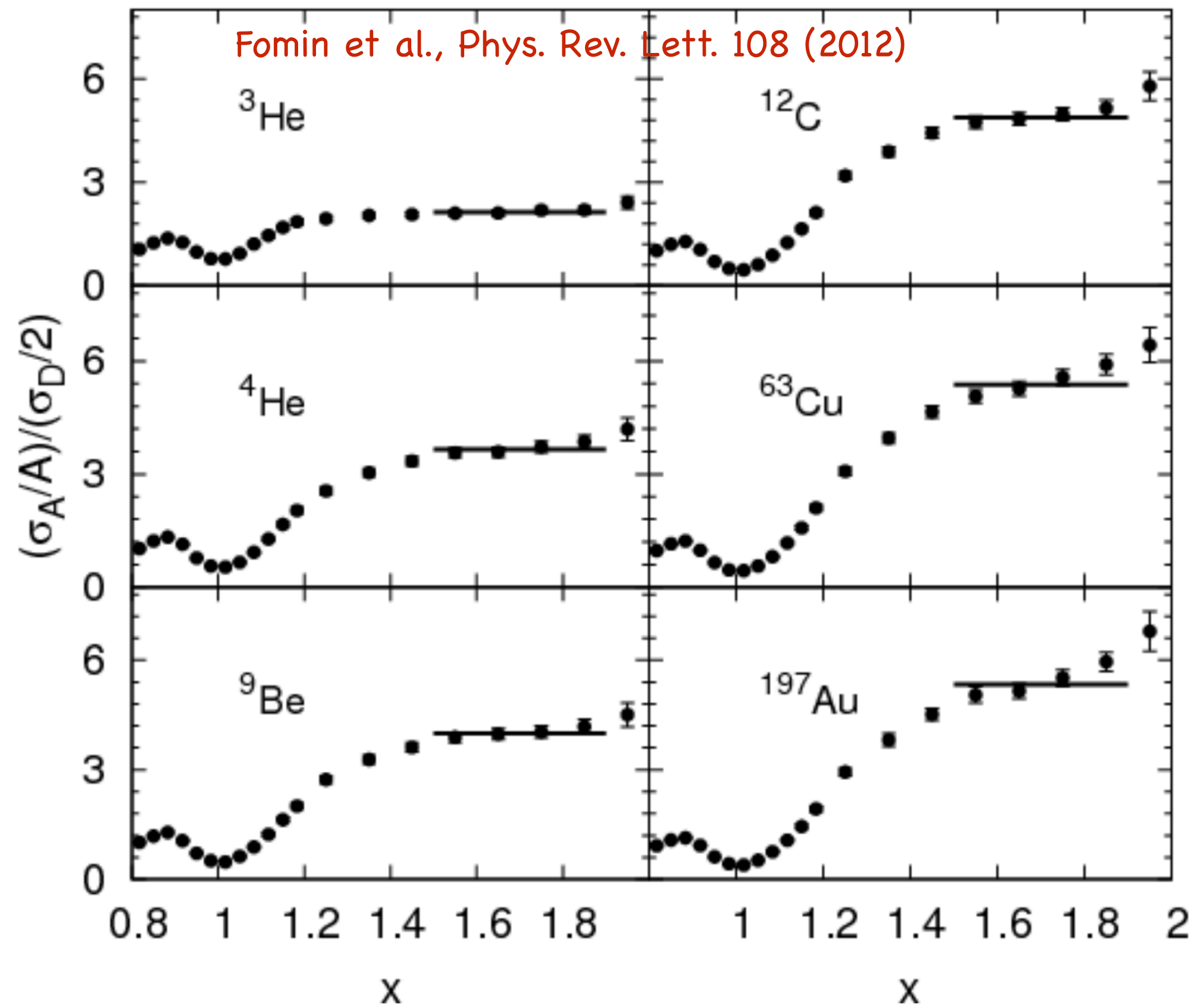
# Modern SRC Phenomenology (2000's - present)



## I) Universal high-momentum tails

inclusive ratios

$$a_2(A) = \frac{2}{A} \frac{\sigma_A(x_B, Q^2)}{\sigma_d(x_B, Q^2)}$$

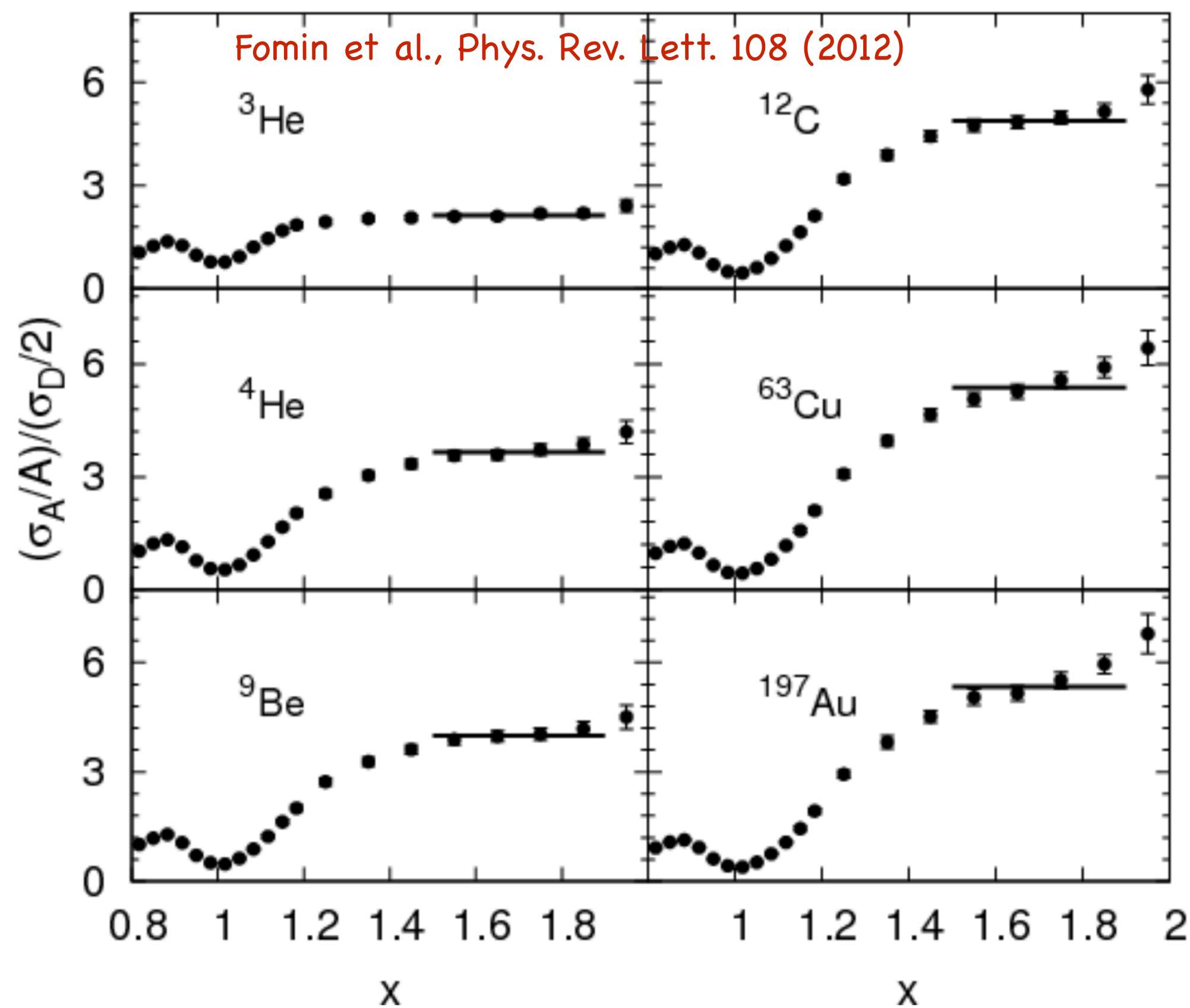




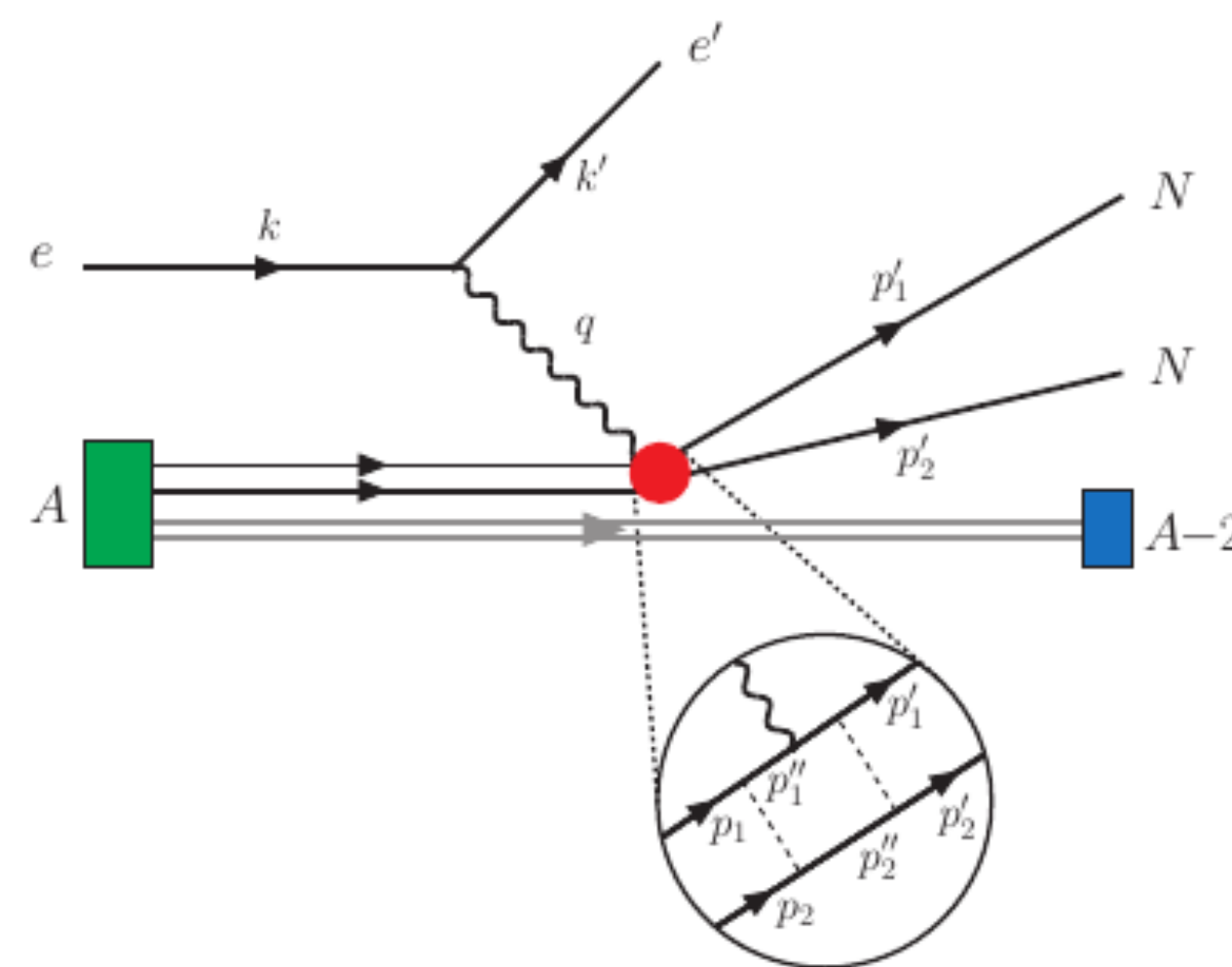
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**SRC interpretation:**

NN interaction scatters pair  $p_1, p_2 < k_F$  to intermediate-state momenta  $\gg k_F$  which are then knocked out by photon

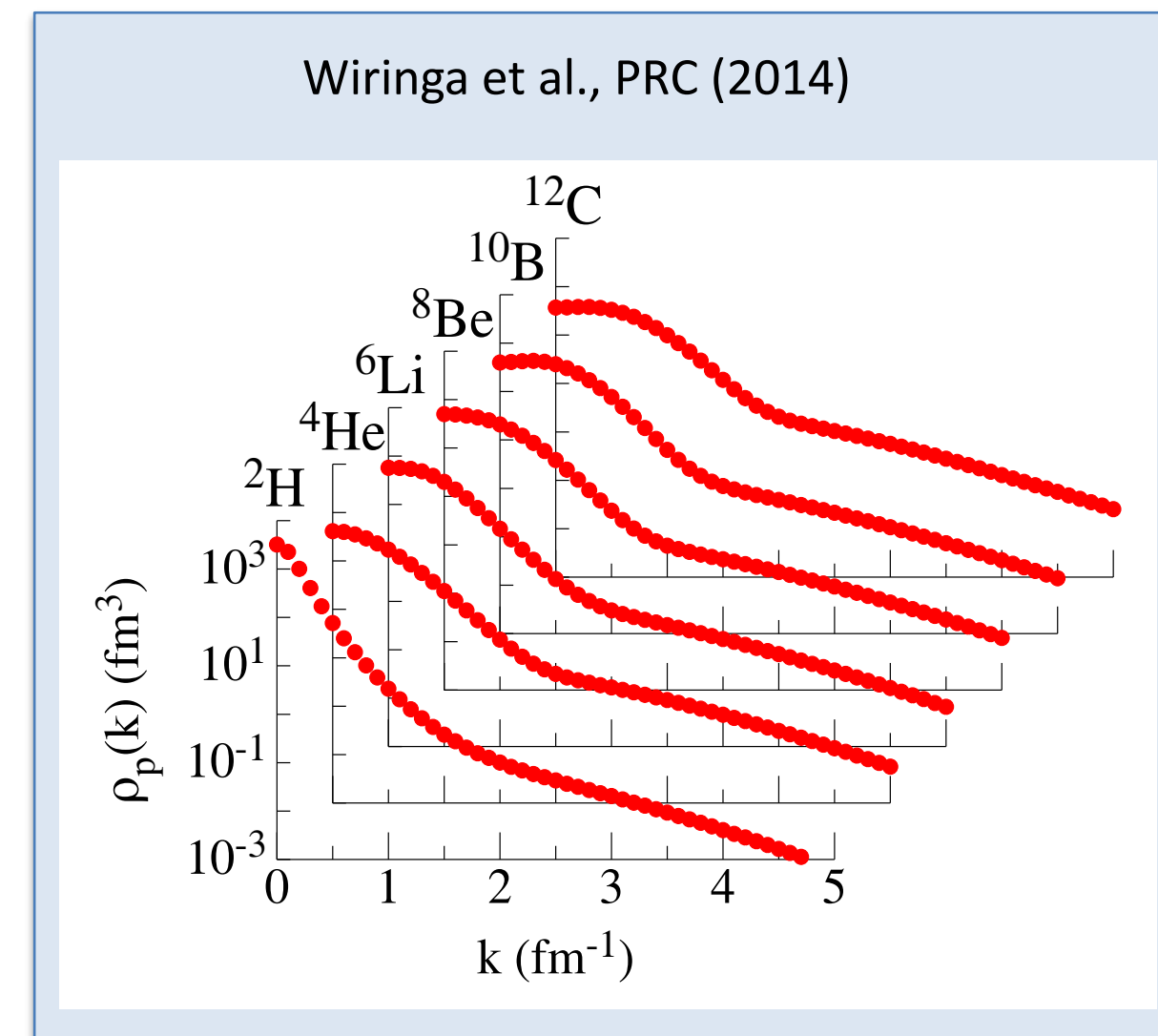
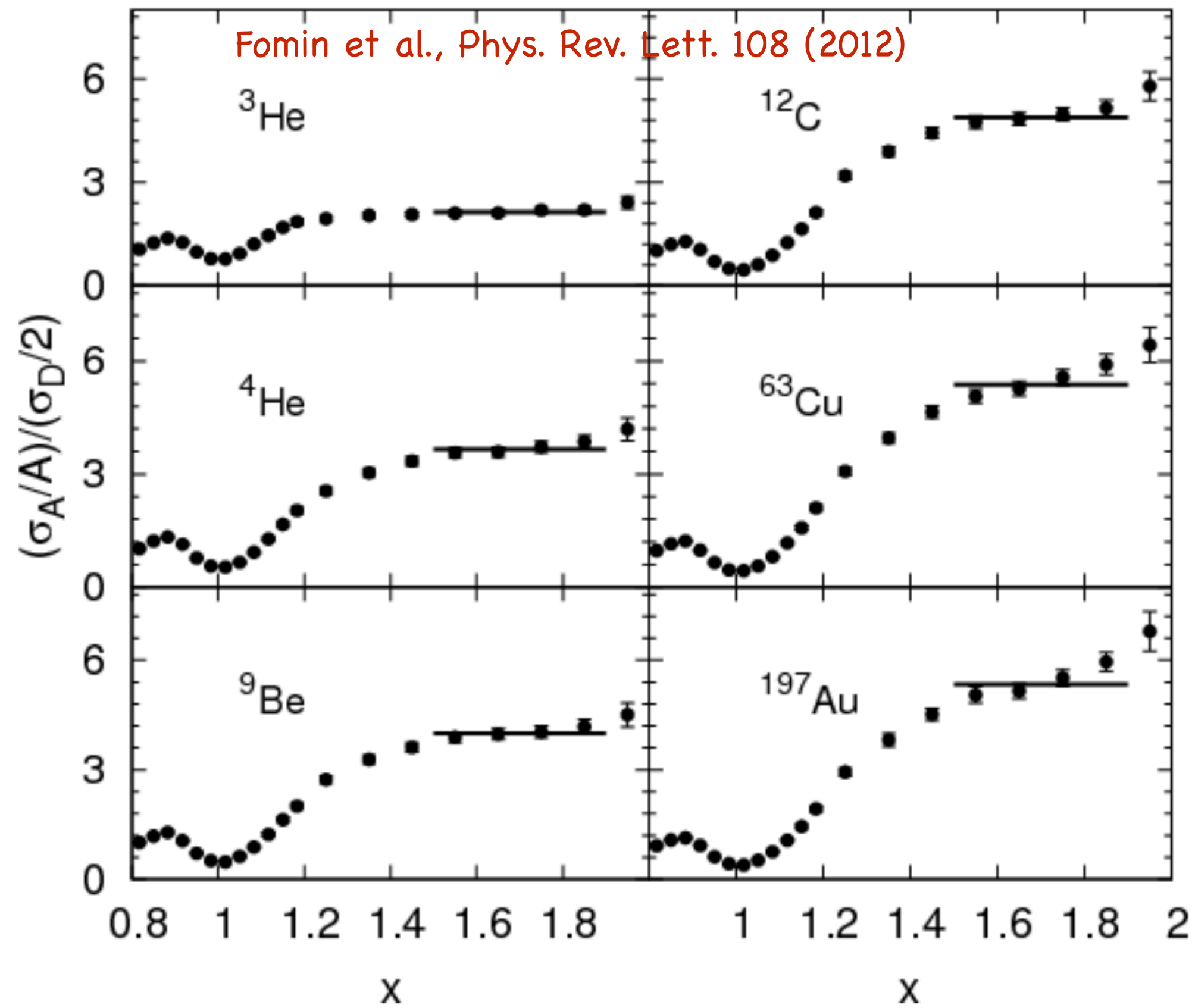
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plateaus in  $x \Rightarrow$  universal (all nuclei) high- $q$  momentum distributions



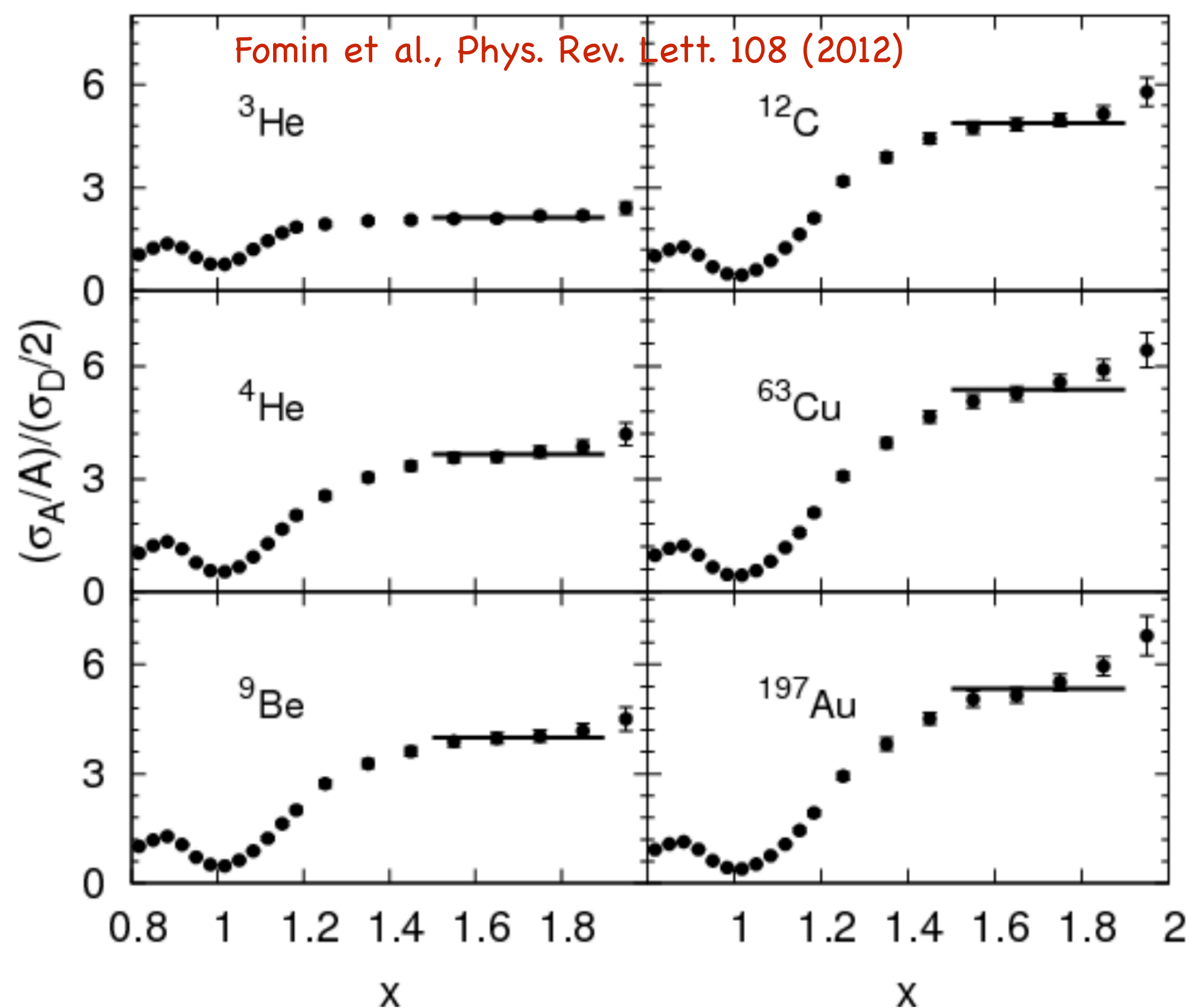
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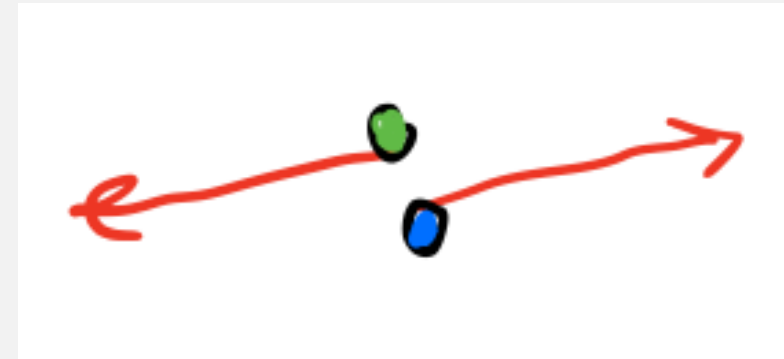
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relative plateau height  $\Rightarrow$  relative prob. of finding  $2N$  SRC

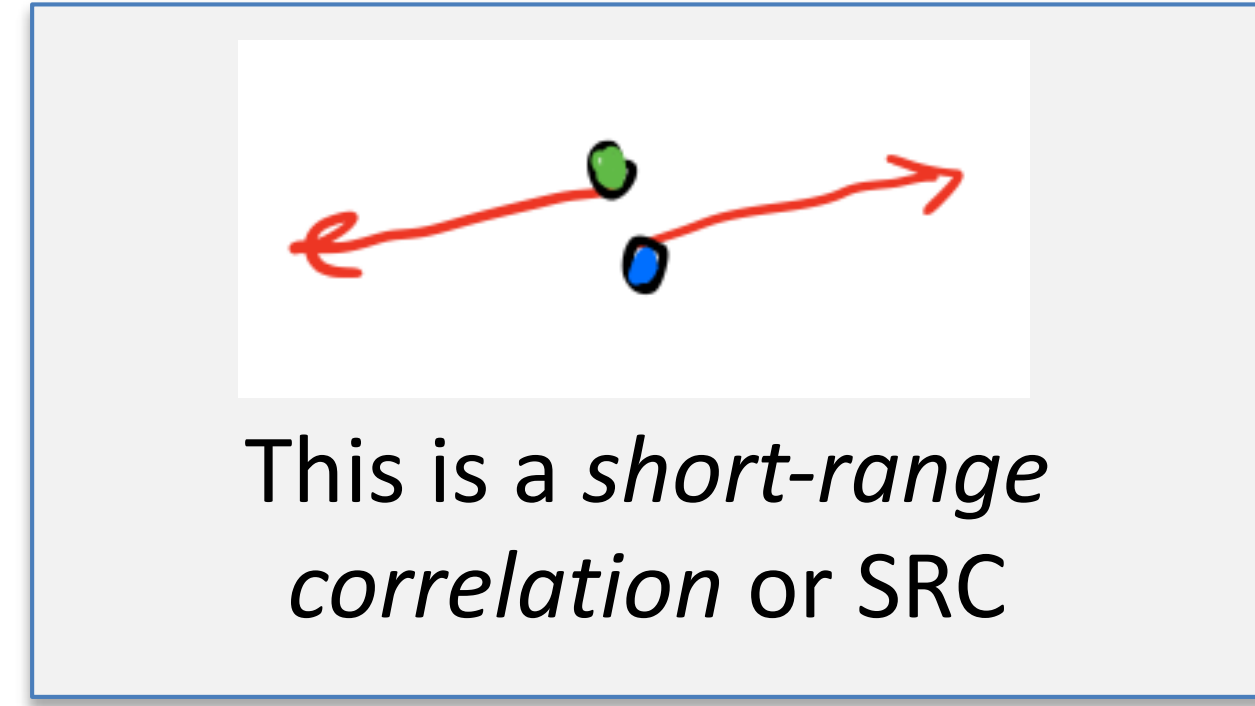
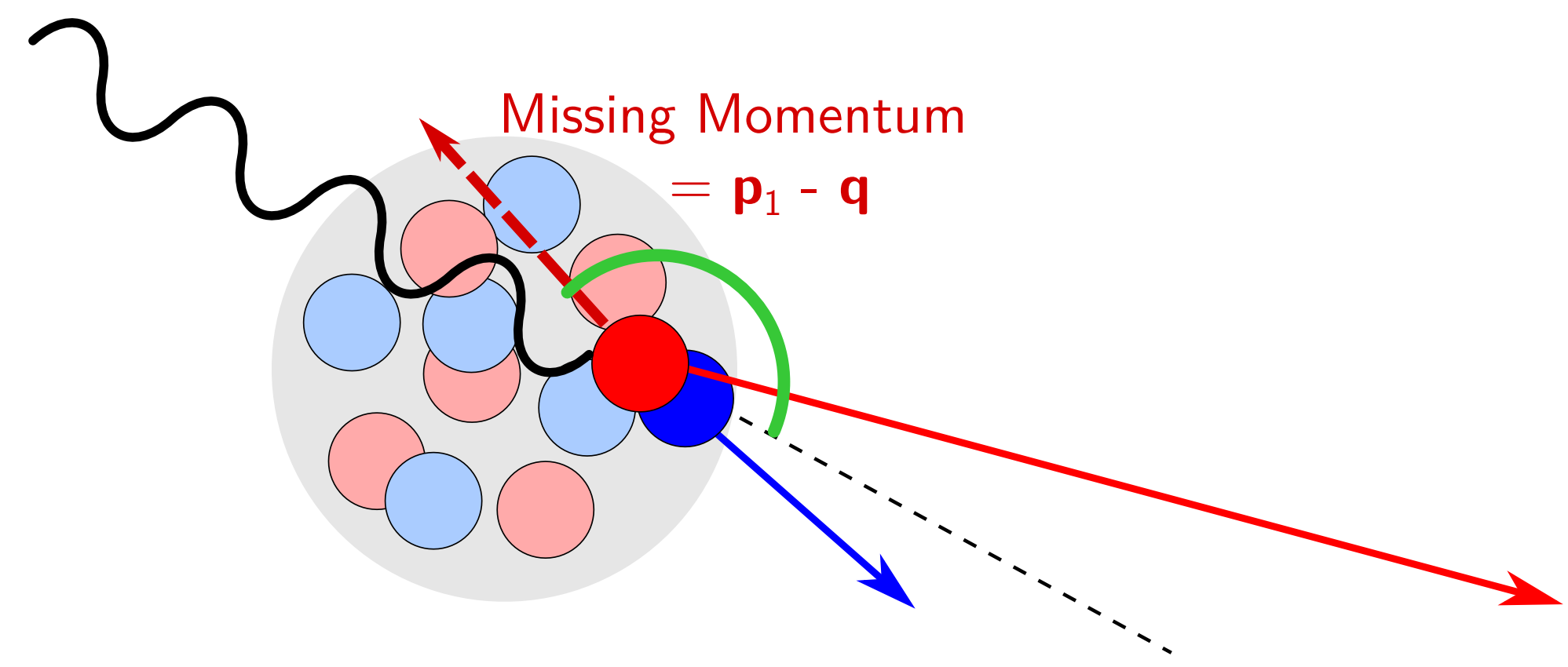


## 2) Kinematics of knocked-out nucleons



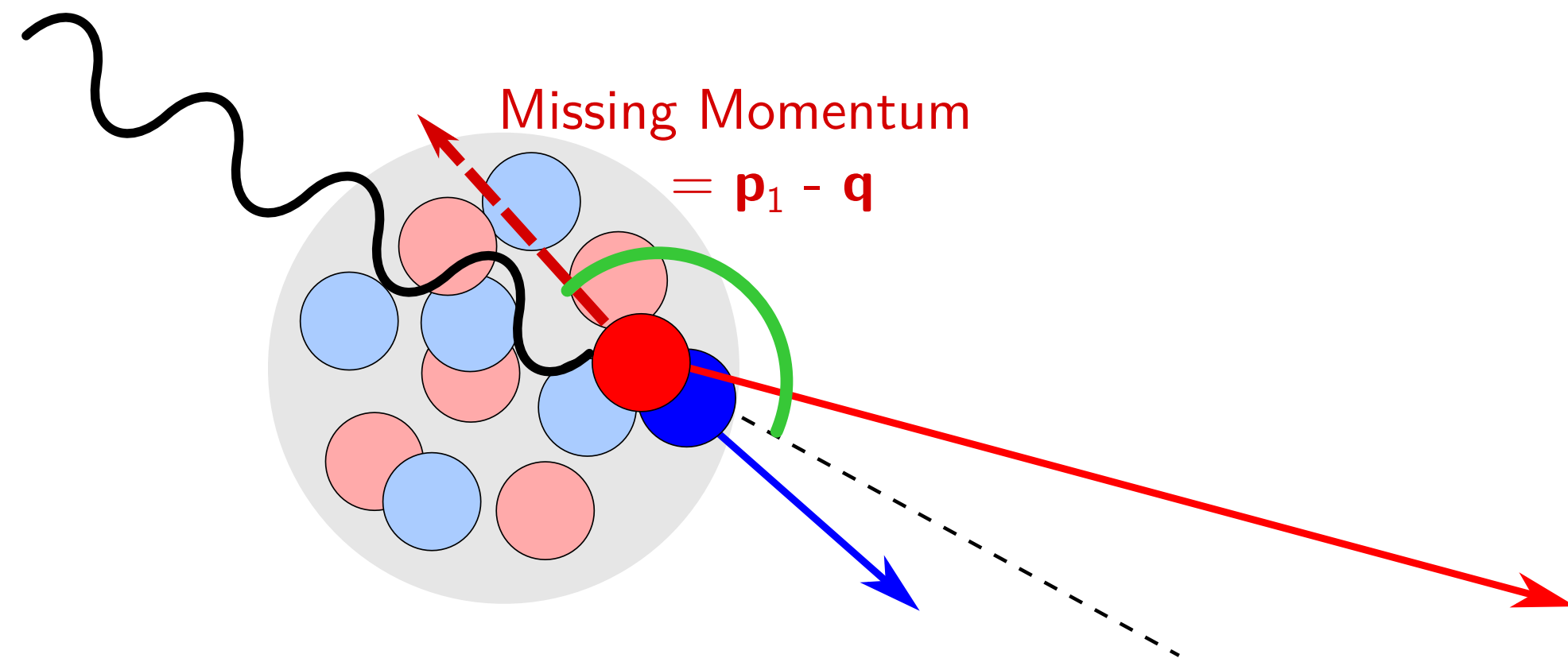
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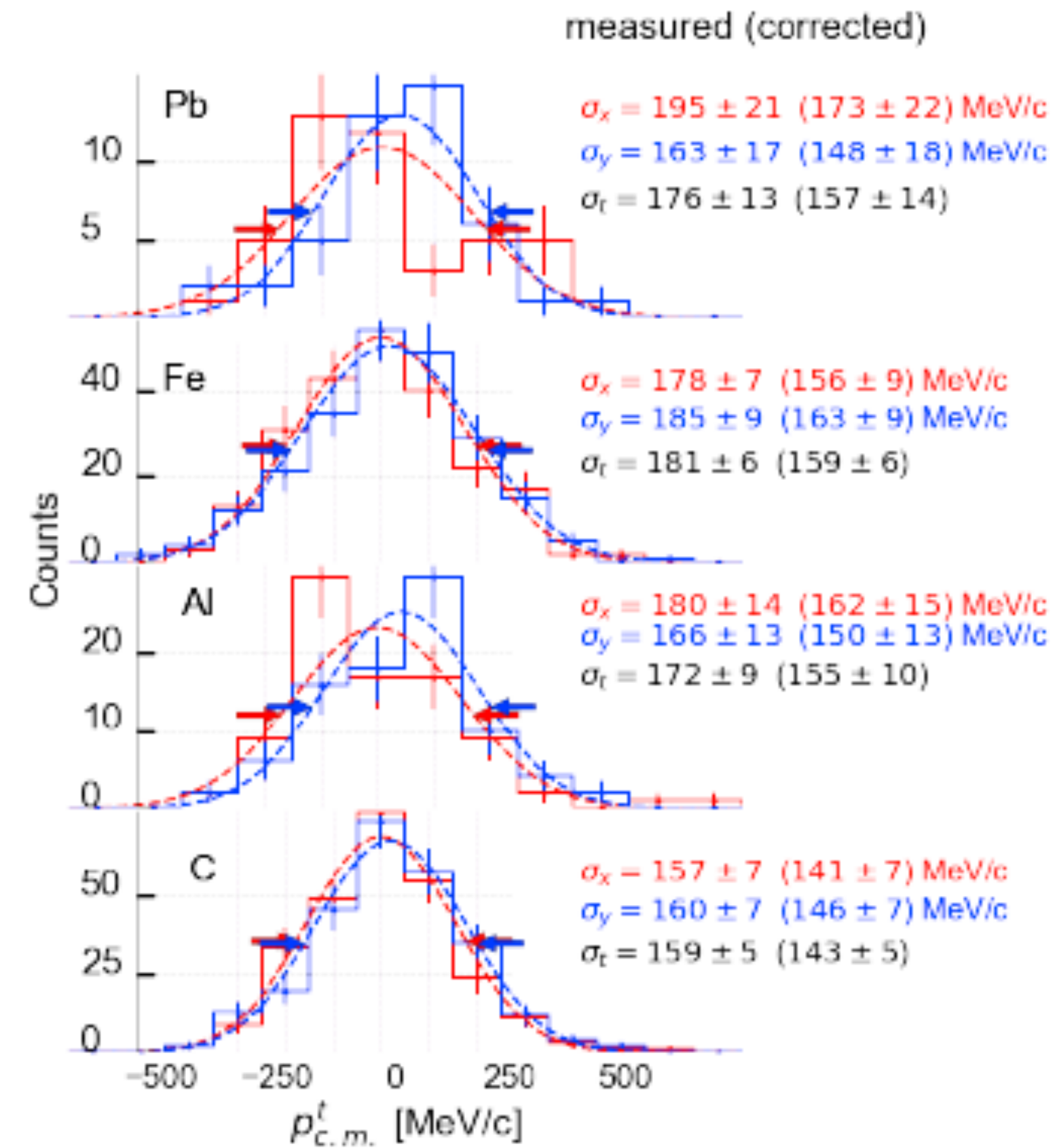


knocked out SRC nucleons fly out almost back-to-back (relative s-wave pairs)

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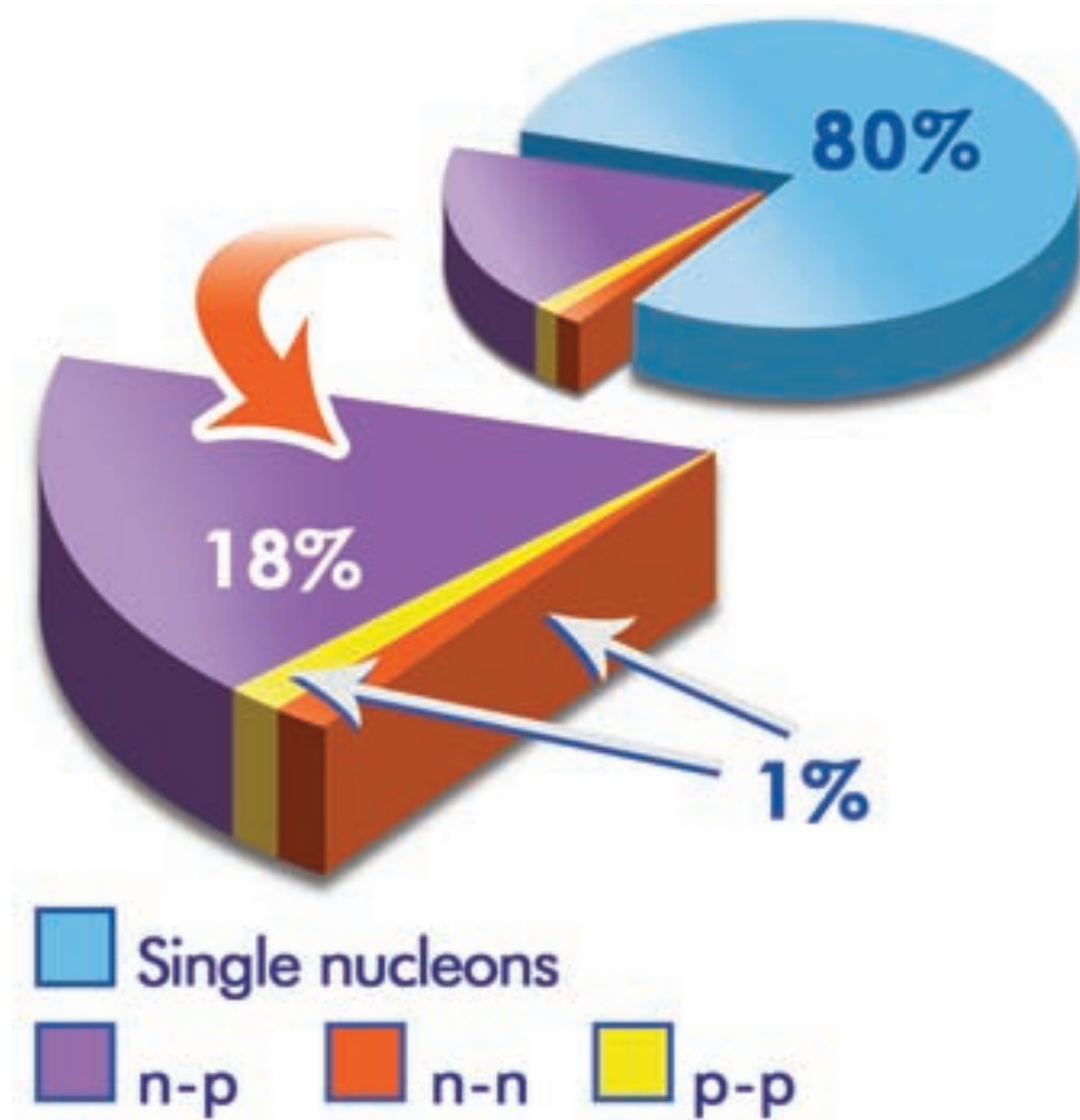


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pair CM momentum distribution gaussian of width  $\sim k_F$

## 3) np dominance at intermediate (300-500 MeV) relative momenta



**Fig. 3.** The average fraction of nucleons in the various initial-state configurations of  $^{12}\text{C}$ .

R. Subedi et al., Science (2008)

20% of nucleons in SRC pairs  
but mostly neutron-proton

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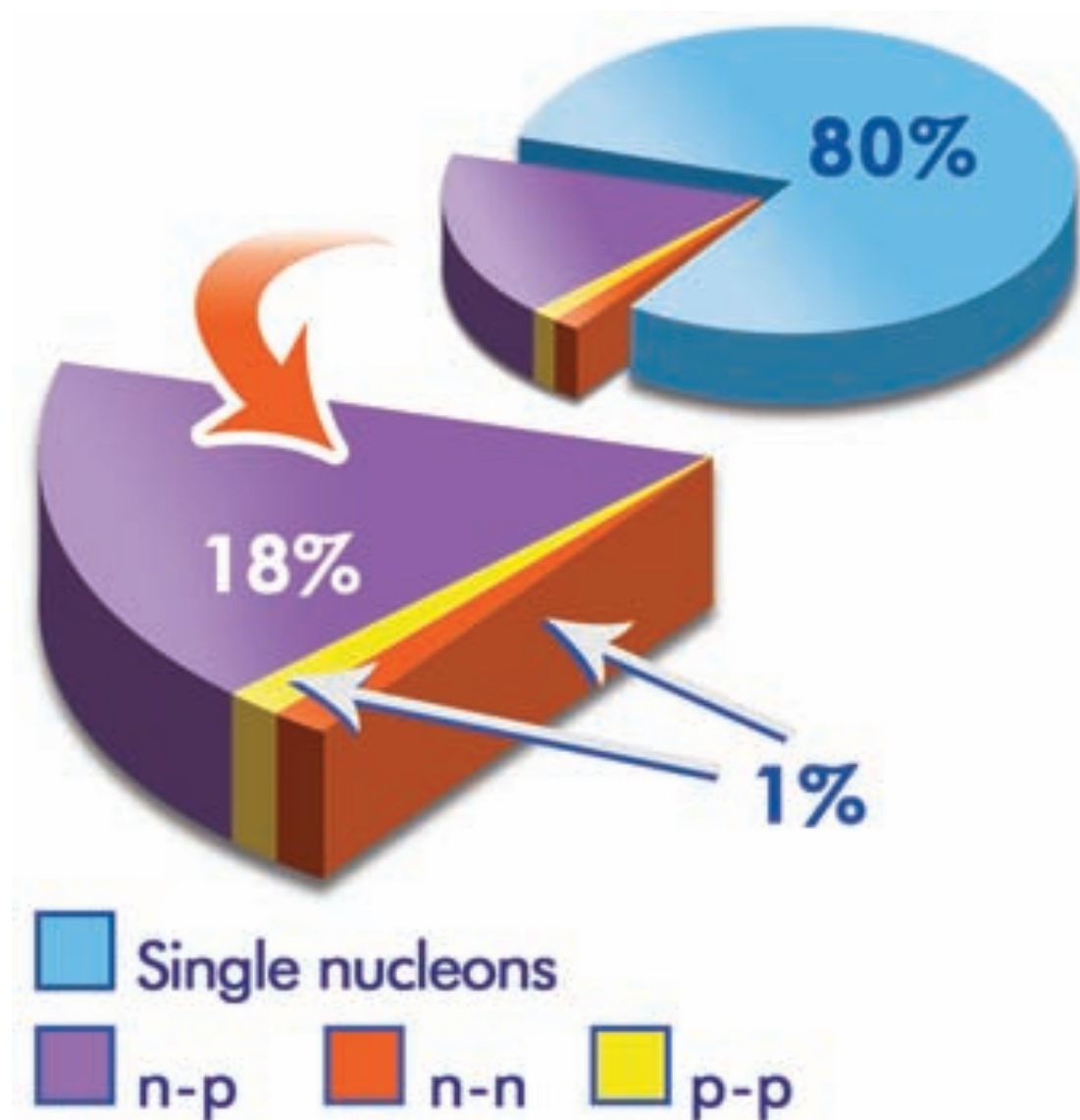
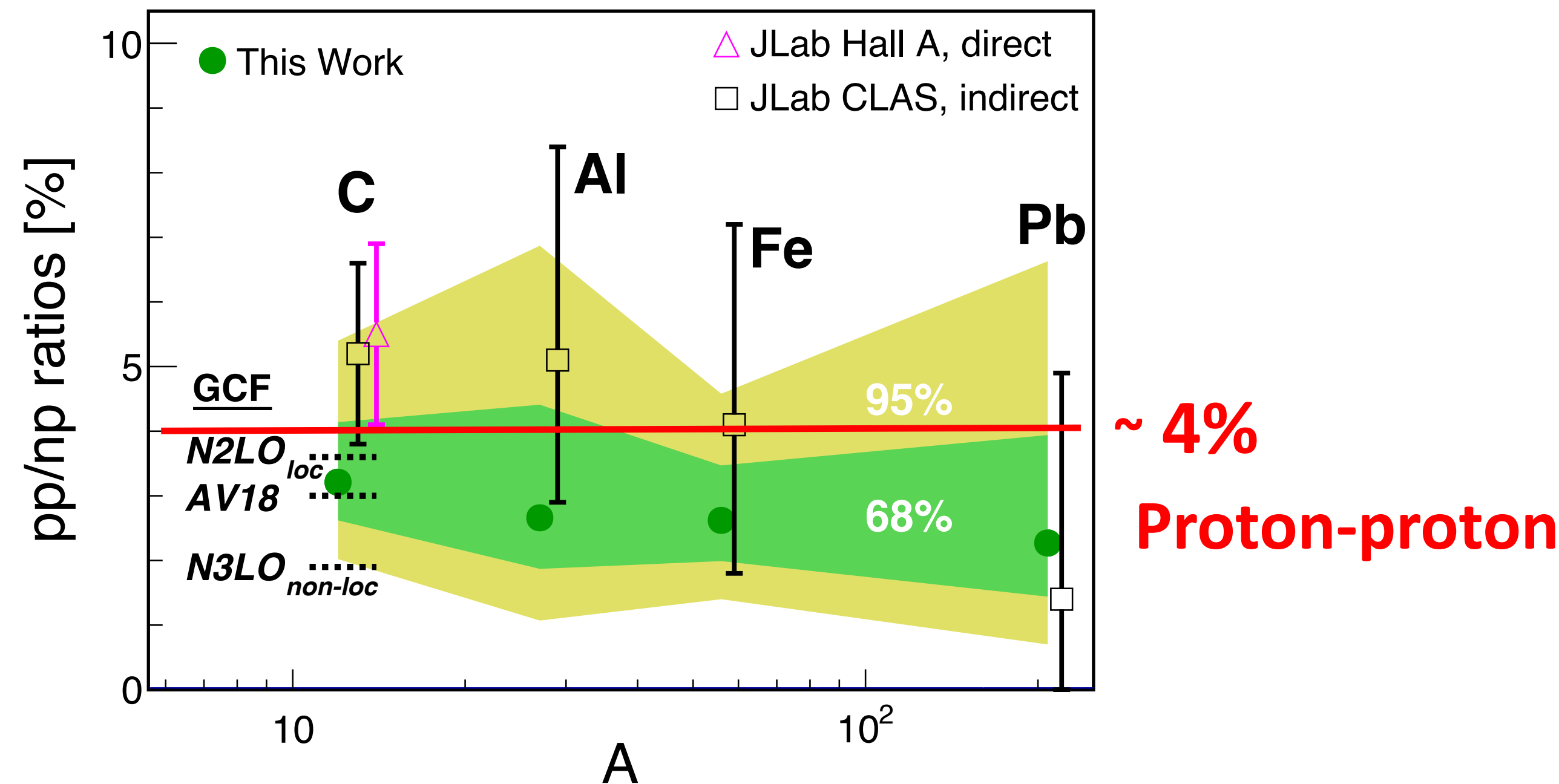


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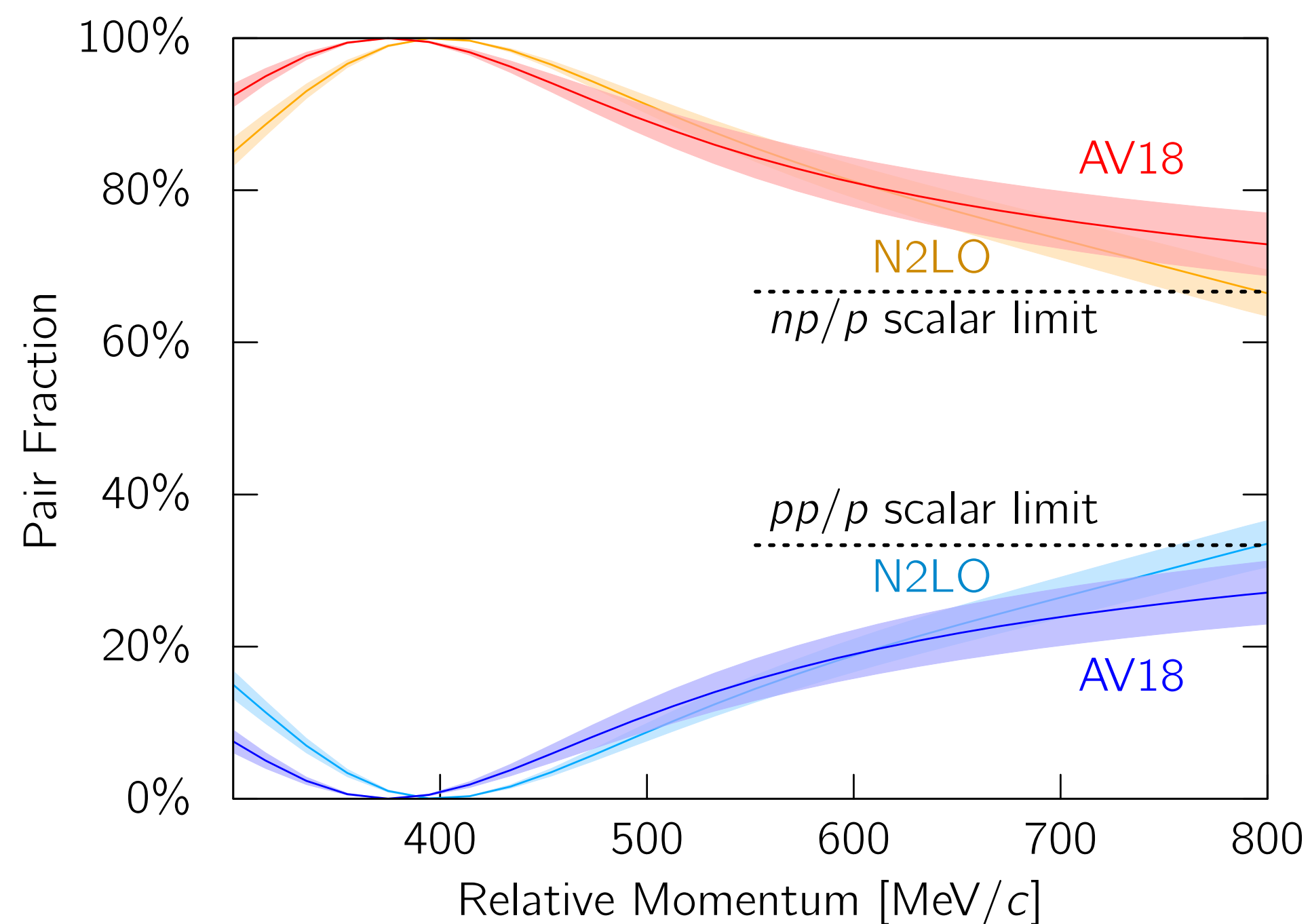
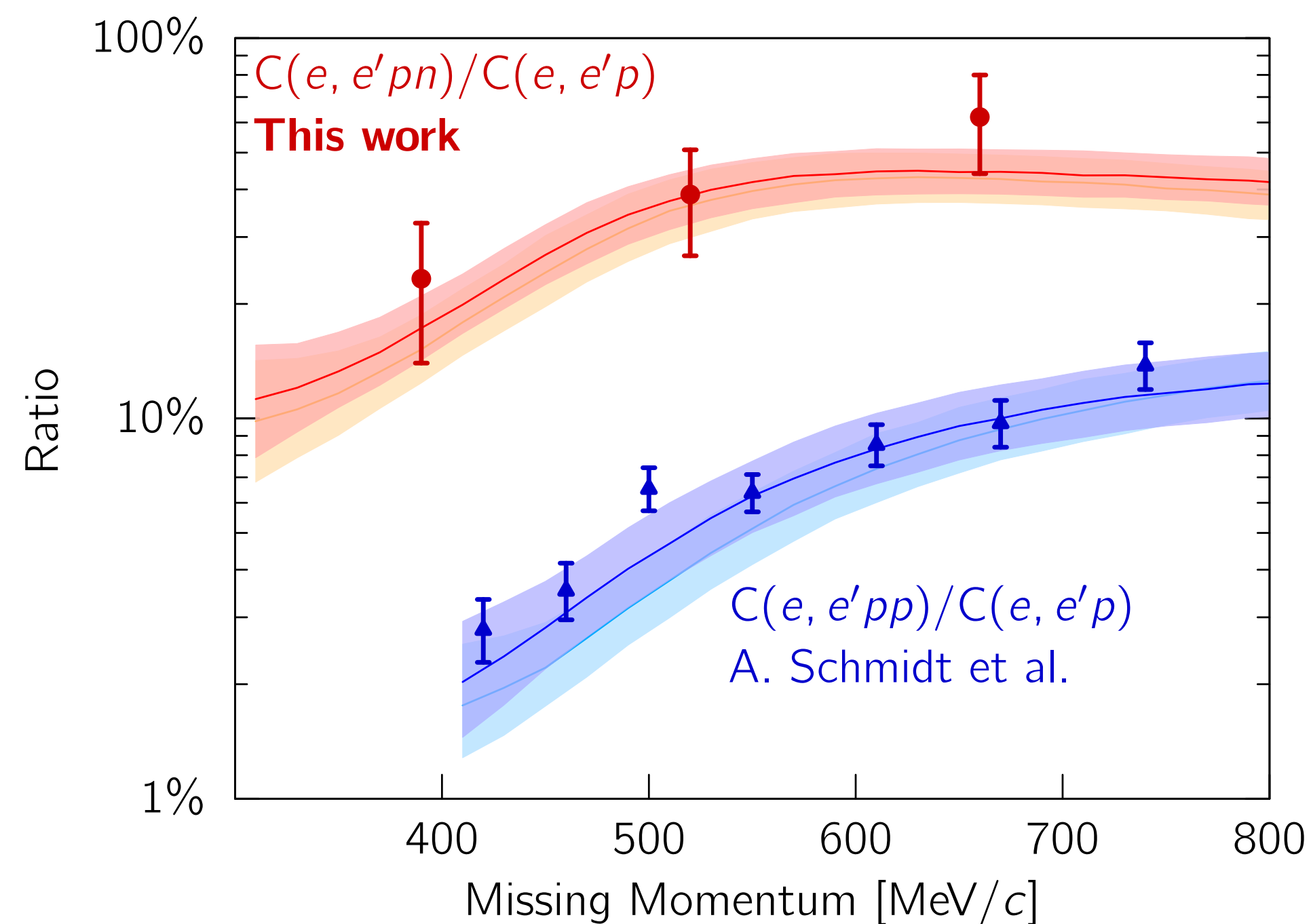
np pairs predominate



Duer, PRL (2019); Duer, Nature (2018); Hen, Science (2014); Korover, PRL (2014); Subedi, Science (2008); Shneor, PRL (2007); Piassetzky, PRL (2006); Tang, PRL (2003); Review: Hen RMP (2017);



## 4) transition to scalar counting at higher relative momentum



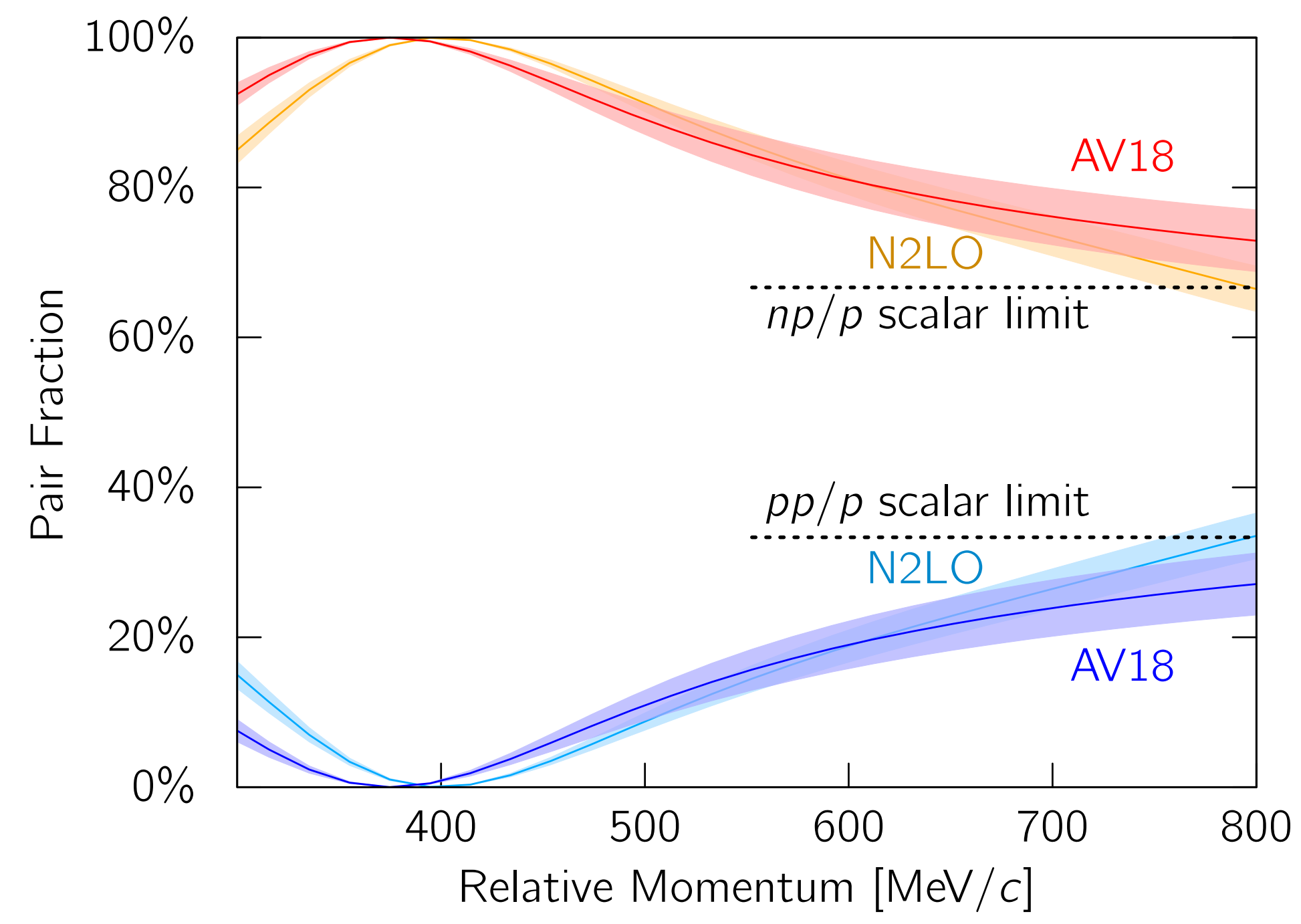
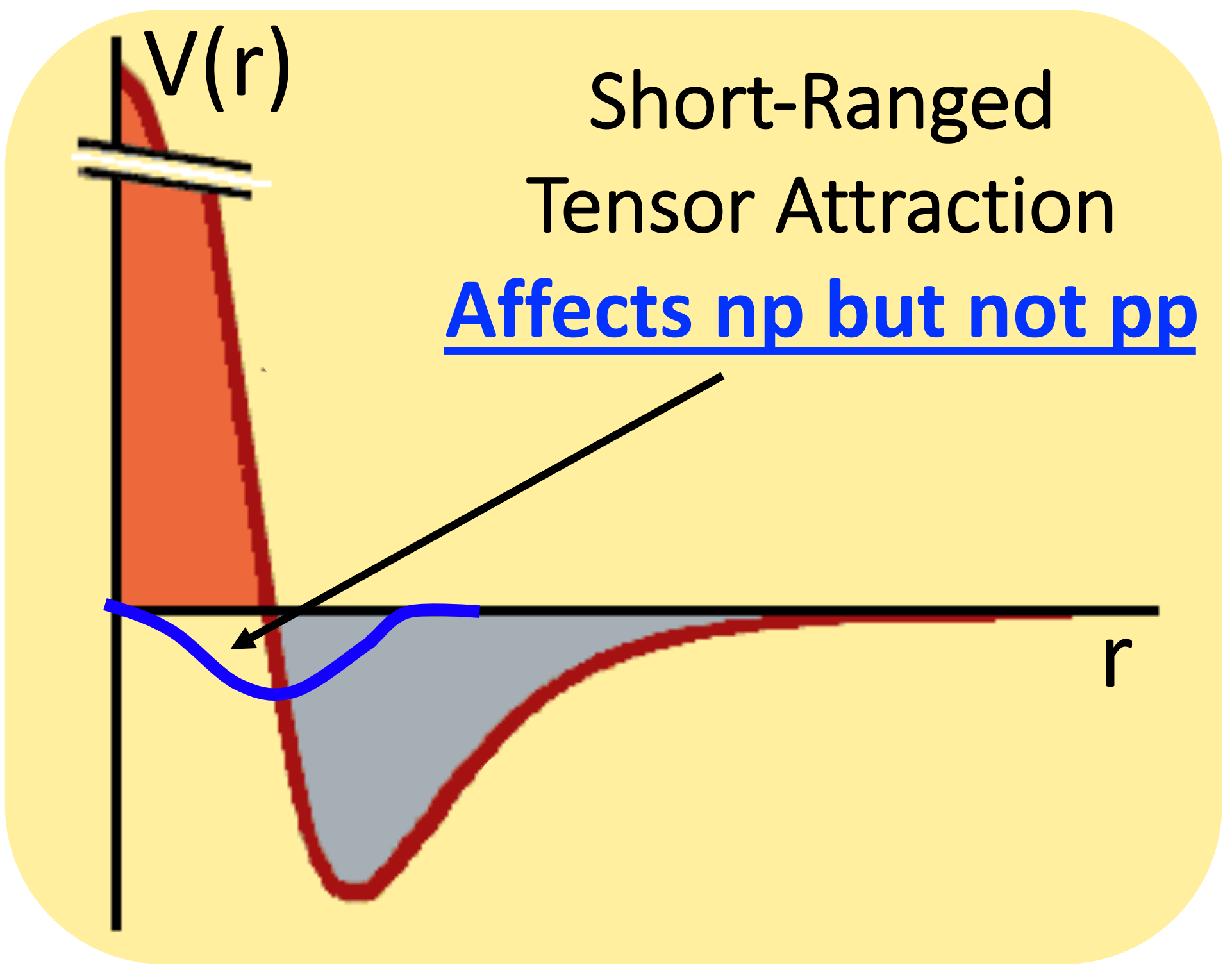
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## 5) Protons “speed up” in neutron rich nuclei

Experiments with increasingly neutron-rich nuclei:

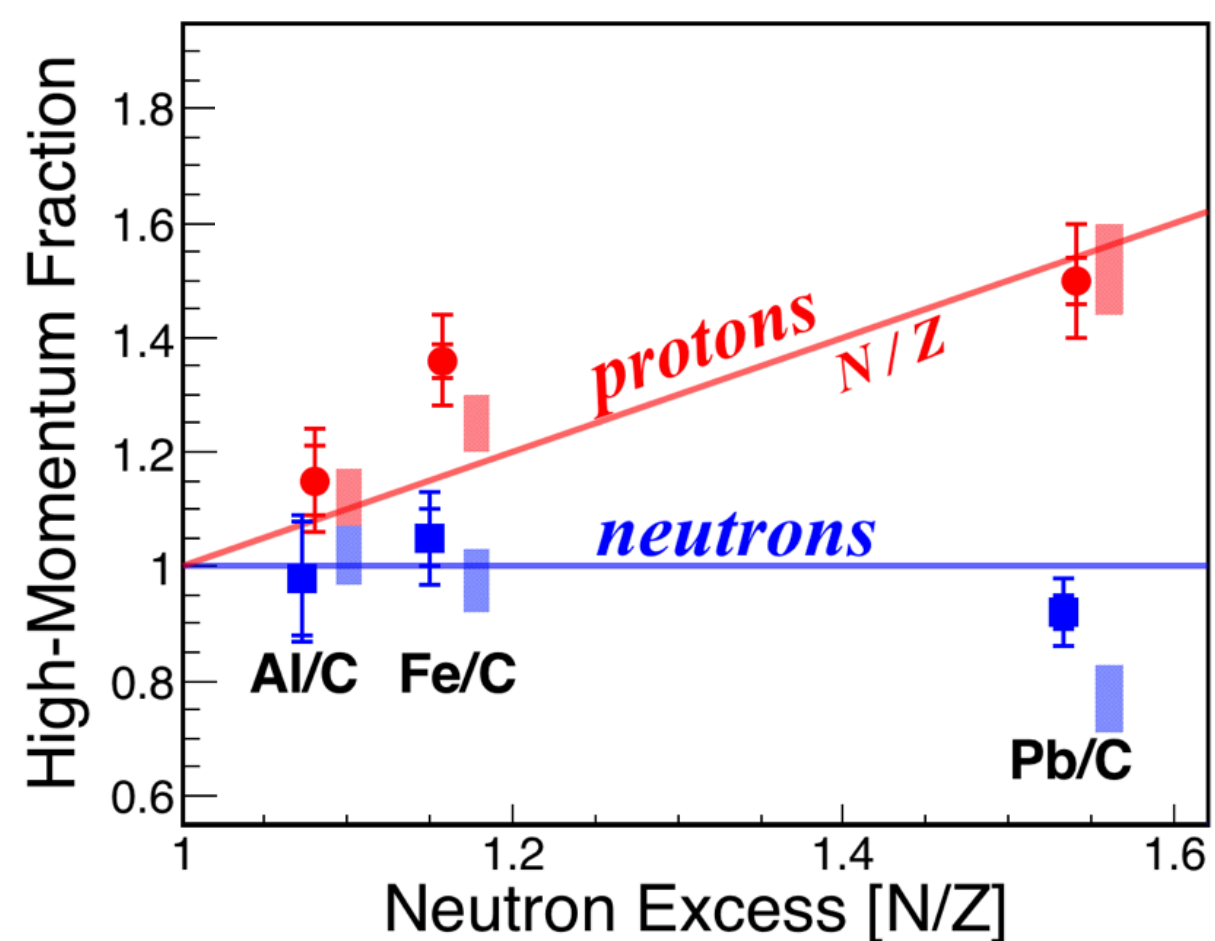
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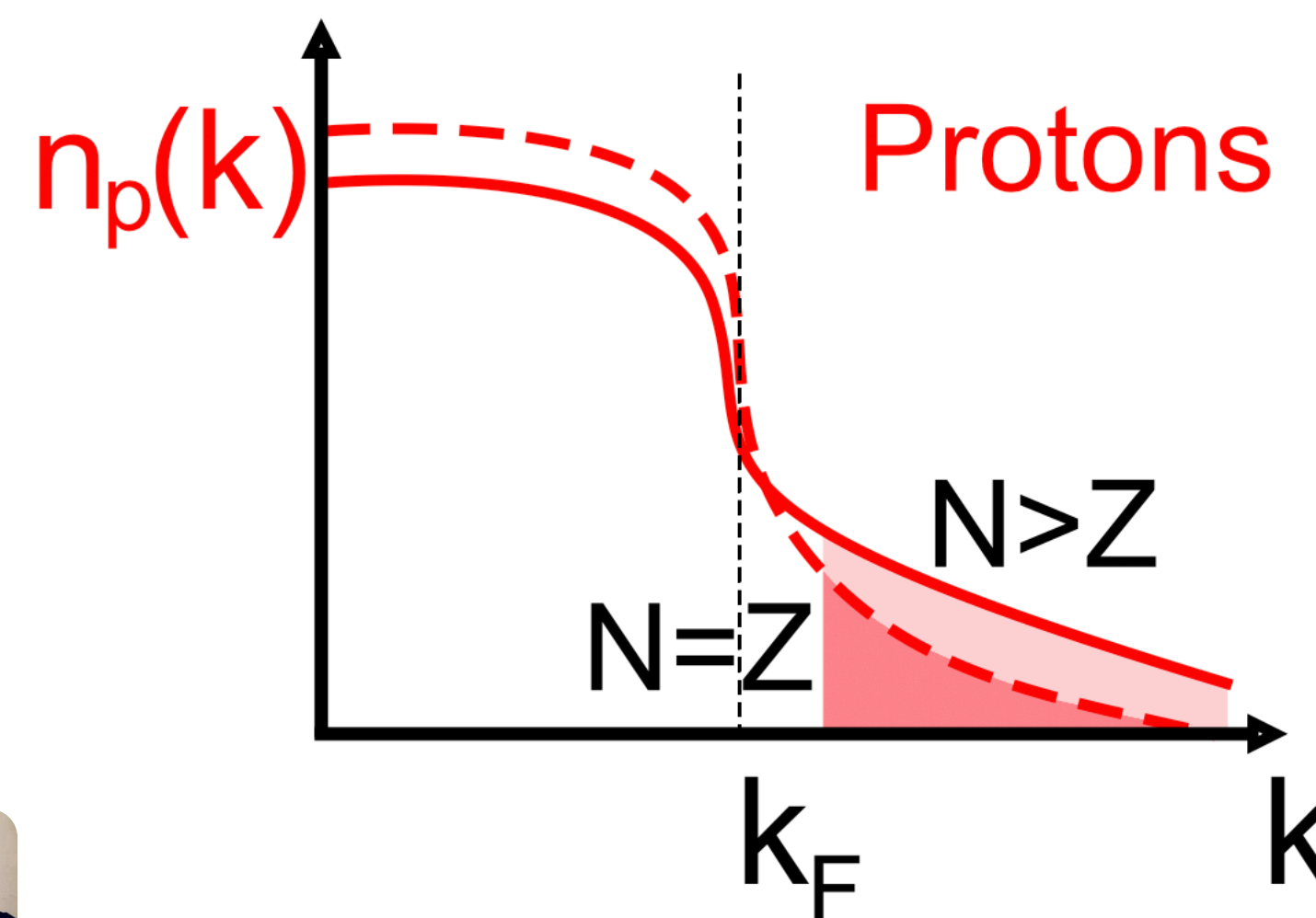
Experiments with increasingly neutron-rich nuclei:  
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Correlation Probability:  
 Neutrons saturate Protons grow



Duer Nature (2018)

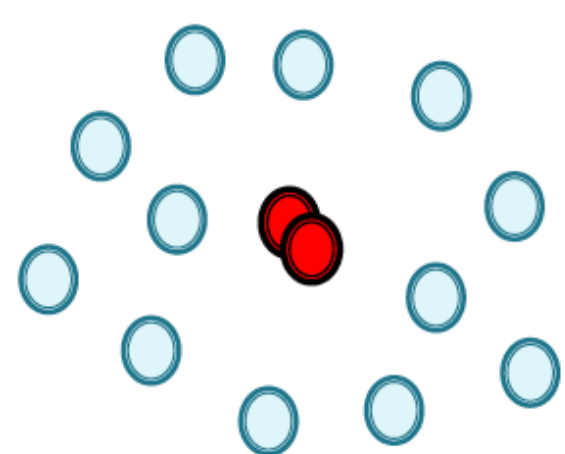
## Protons 'Speed-Up' In Neutron-Rich Nuclei



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## 6) Generalized Contact Formalism (GCF)

Cruz-Torres et al. arXiv:1907.03658  
and earlier papers of Weiss/Barnea/et al.



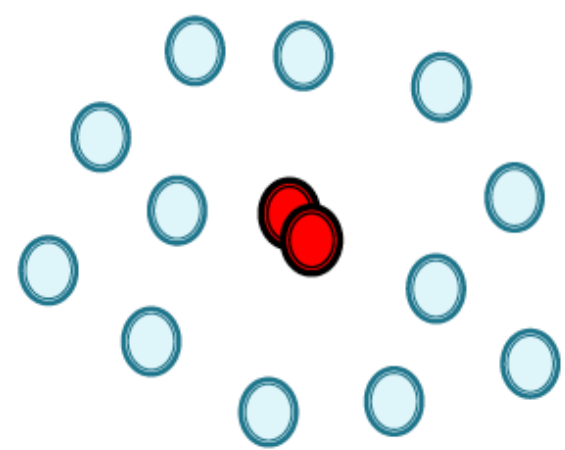
$\mathbf{r}_{12} \rightarrow 0$

GCF has *factorized* small-r / large-k approximation to many-body wave function:

$$\Psi_n^A(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A) \sim \phi(\mathbf{r}_{12}) \chi_n^A(\mathbf{R}_{12}, \mathbf{r}_3, \dots, \mathbf{r}_A) \quad \text{cf. Brueckner 1955, Tan 2005}$$

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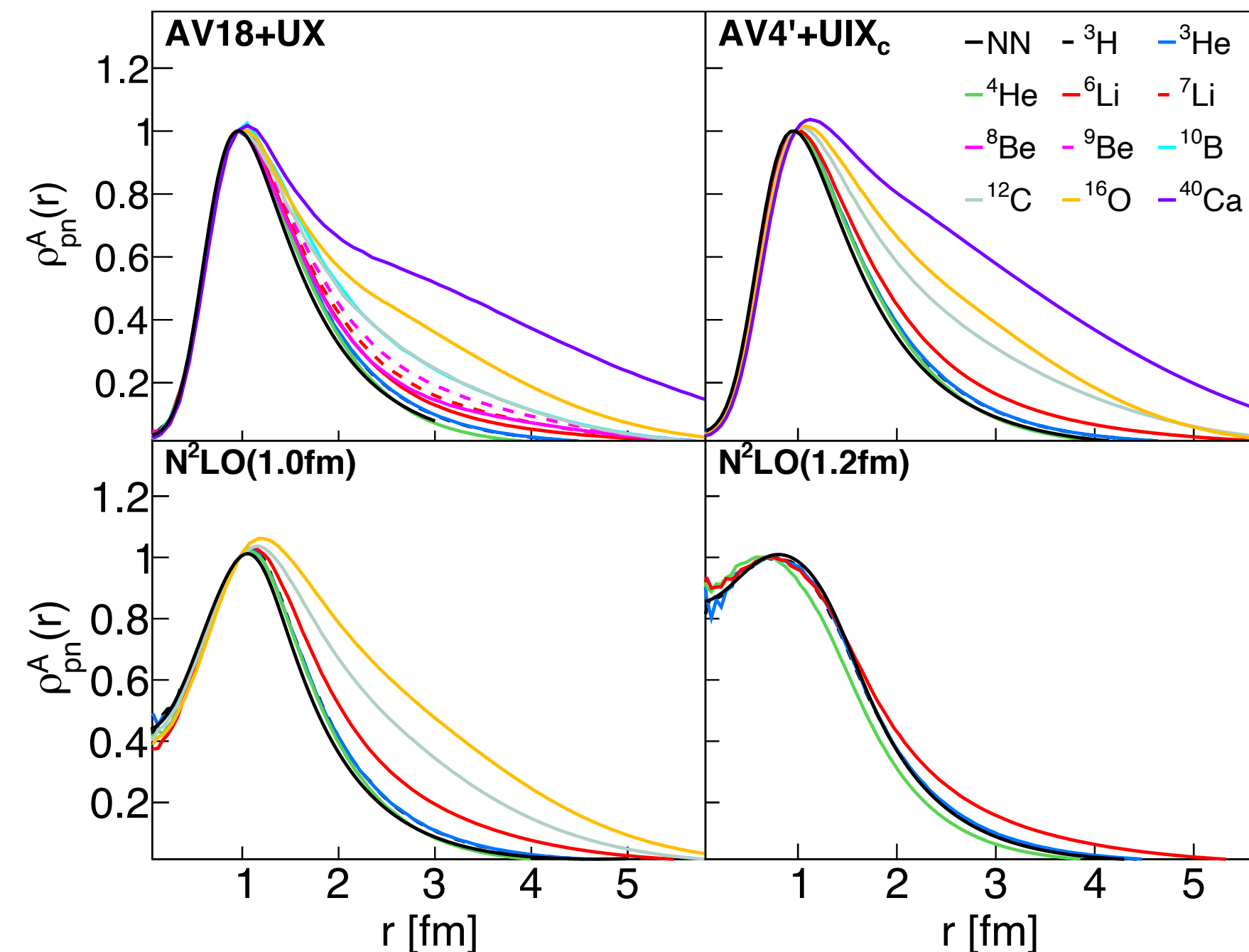
E.g., pair density

$$\rho_A^{NN,\alpha}(\mathbf{r}) = \langle \Psi | \sum_{i < j} \delta(\mathbf{r} - \mathbf{r}_{ij}) P_\alpha | \Psi \rangle$$

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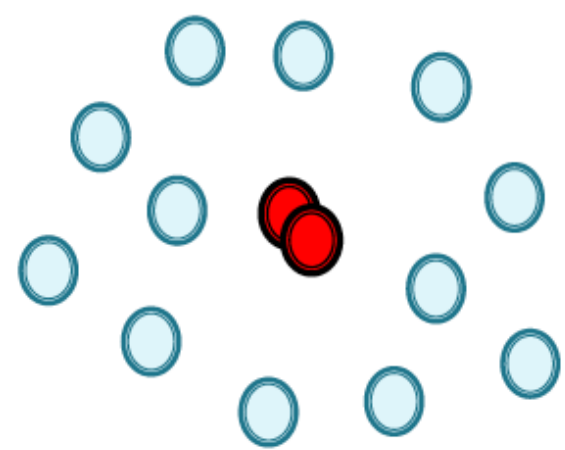
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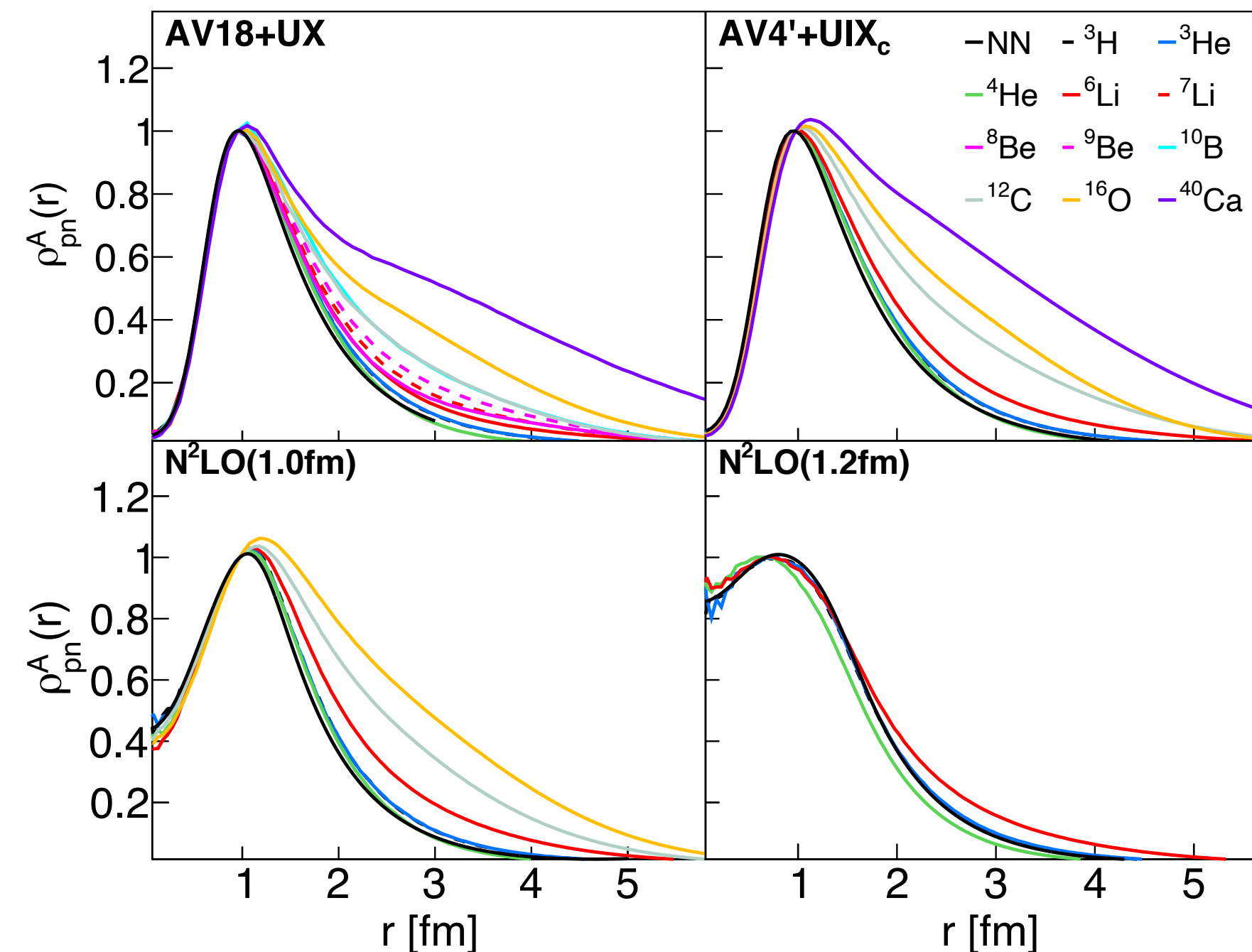
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cf. Brueckner 1955, Tan 2005



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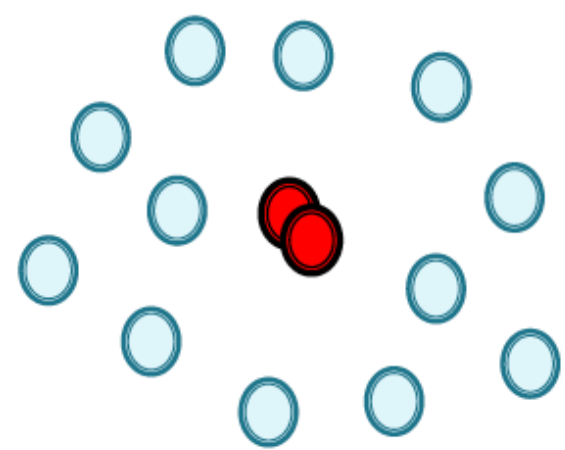
A-dep scale factors ("nuclear contacts")  $C_A \sim \langle \chi | \chi \rangle$

# Modern SRC Phenomenology (2000's - present)



## 6) Generalized Contact Formalism (GCF)

Cruz-Torres et al. arXiv:1907.03658  
and earlier papers of Weiss/Barnea/et al.



$\mathbf{r}_{12} \rightarrow 0$

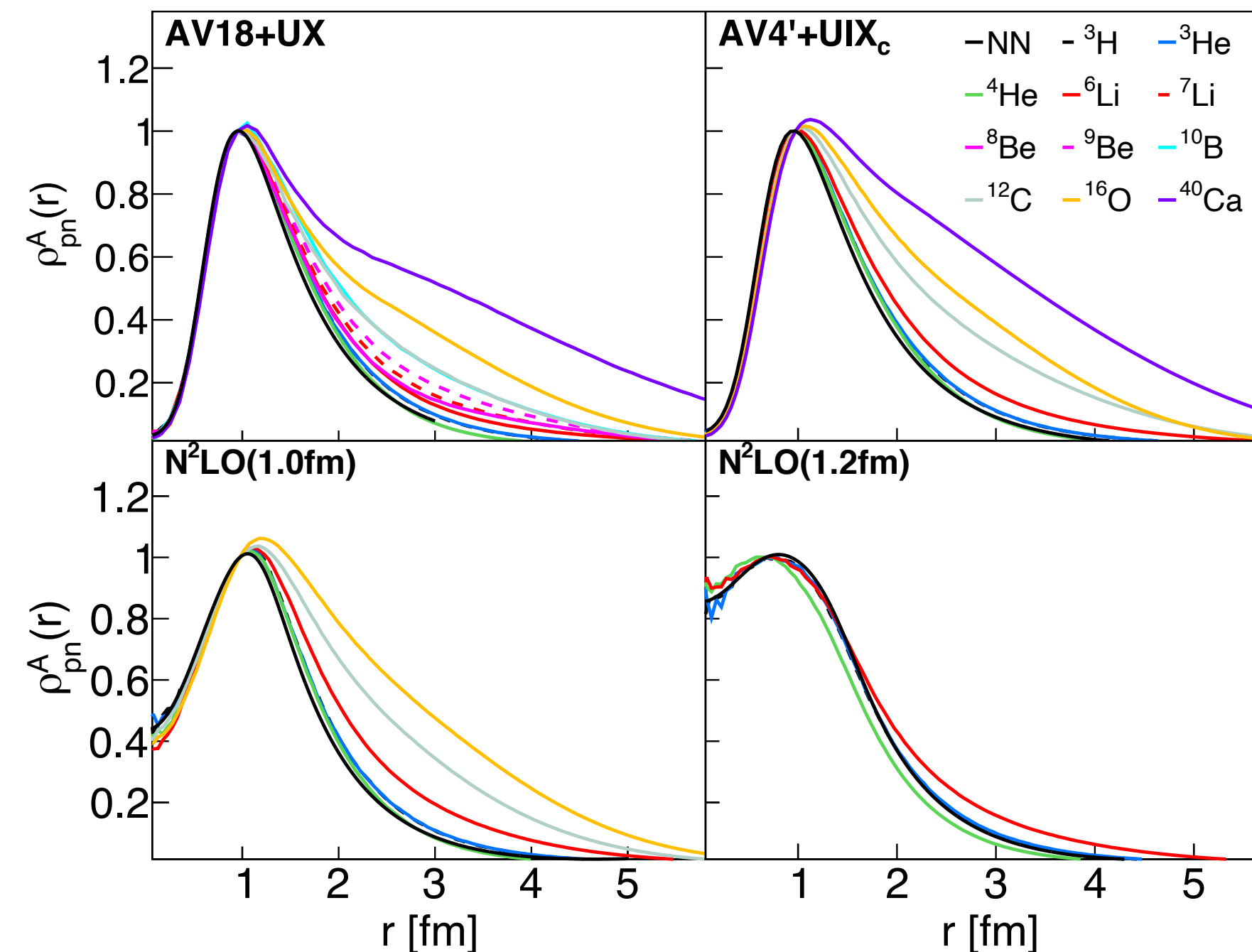
E.g., pair density

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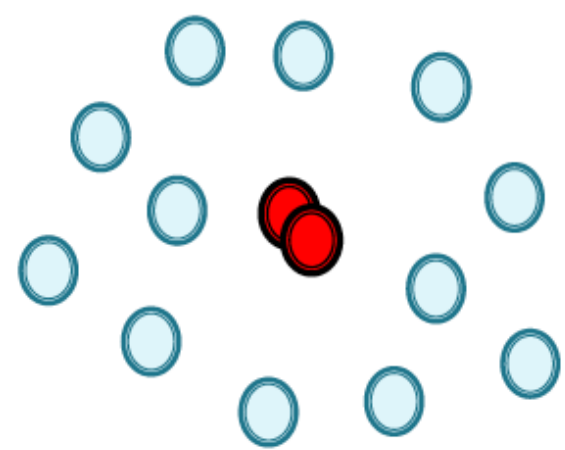
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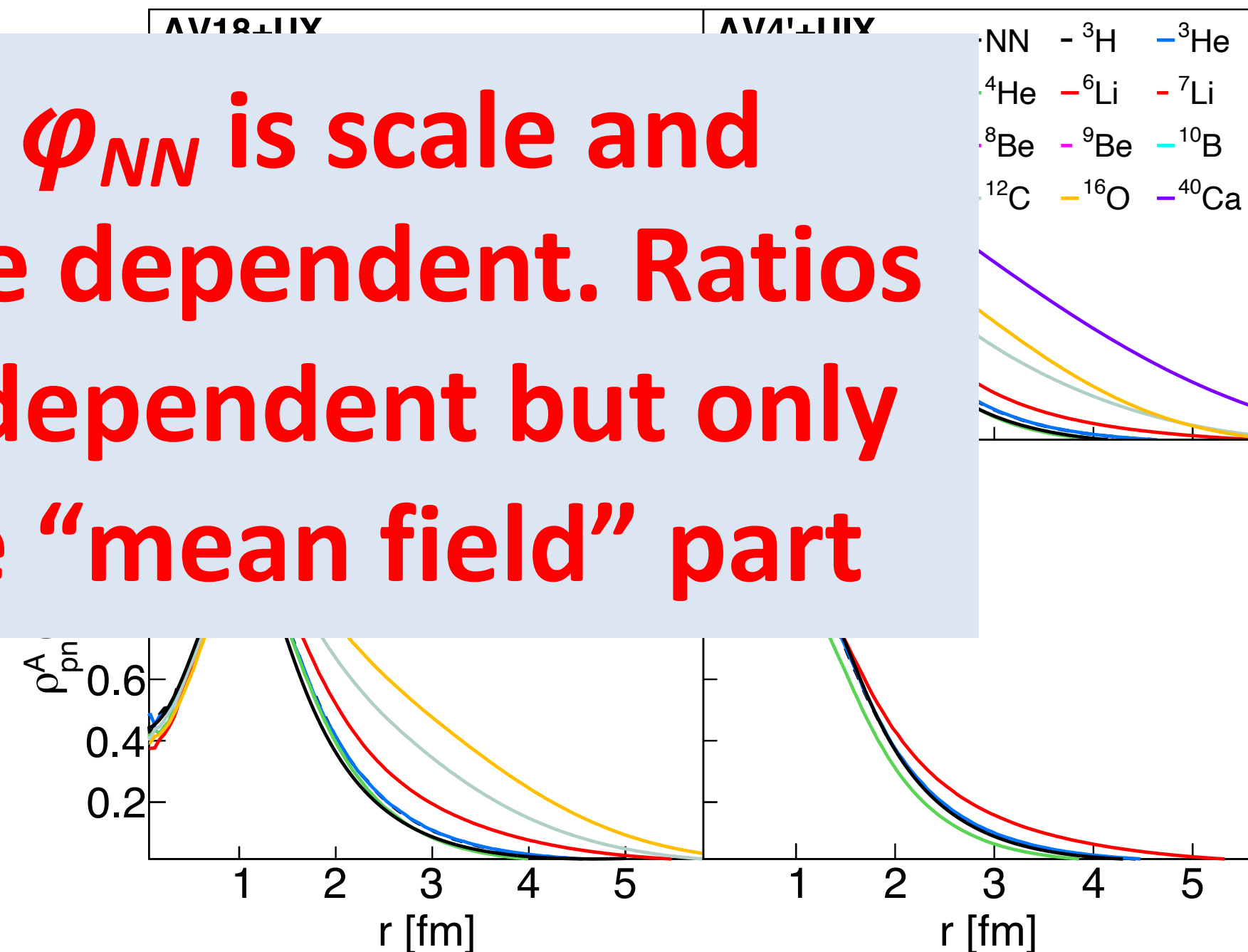


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**But  $\varphi_{NN}$  is scale and scheme dependent. Ratios are independent but only probe "mean field" part**



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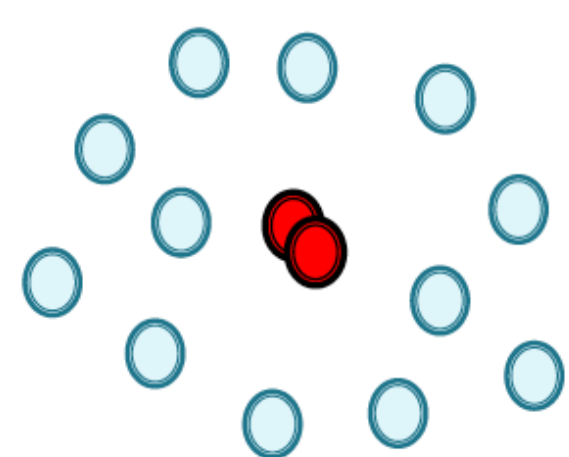
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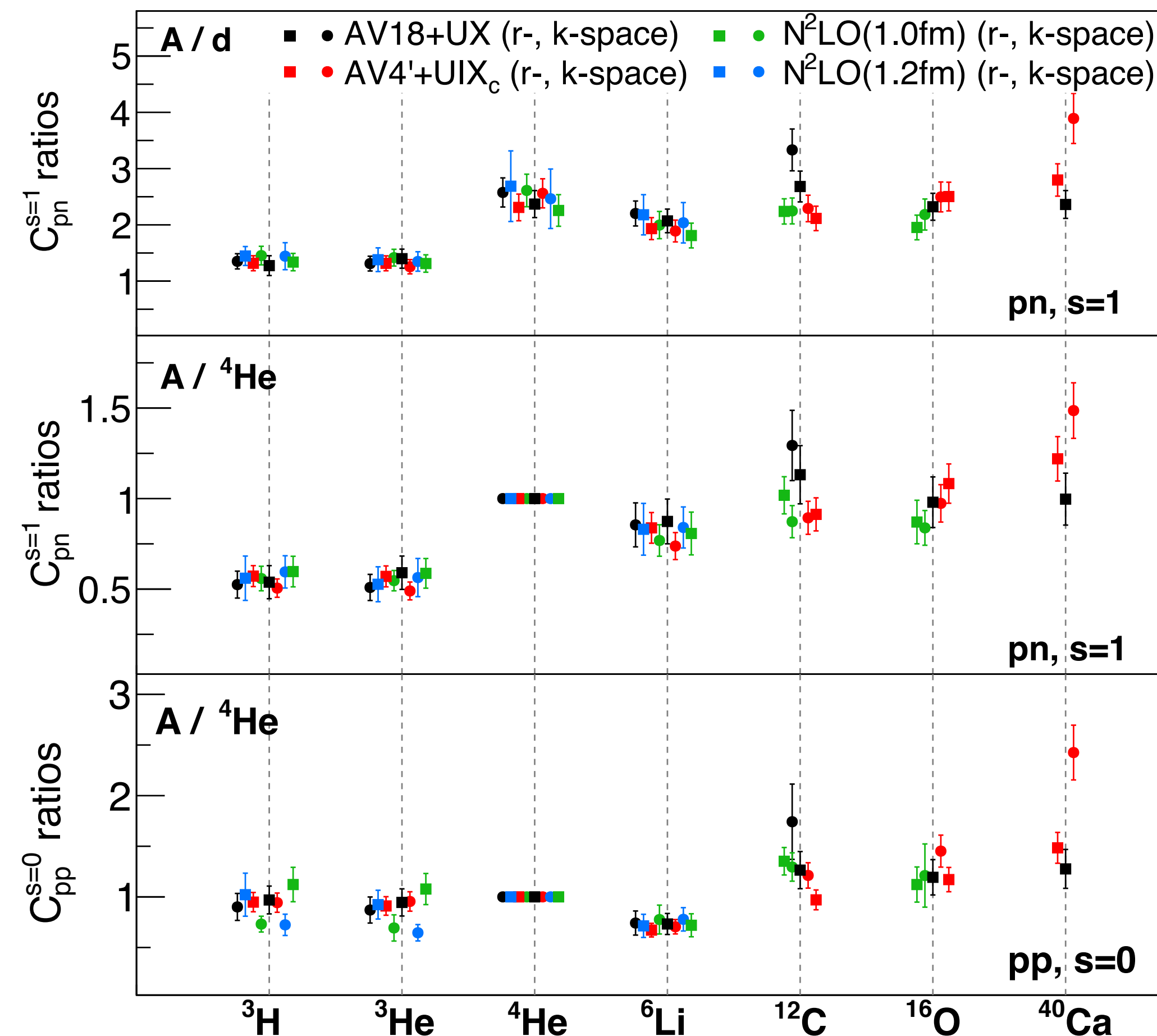
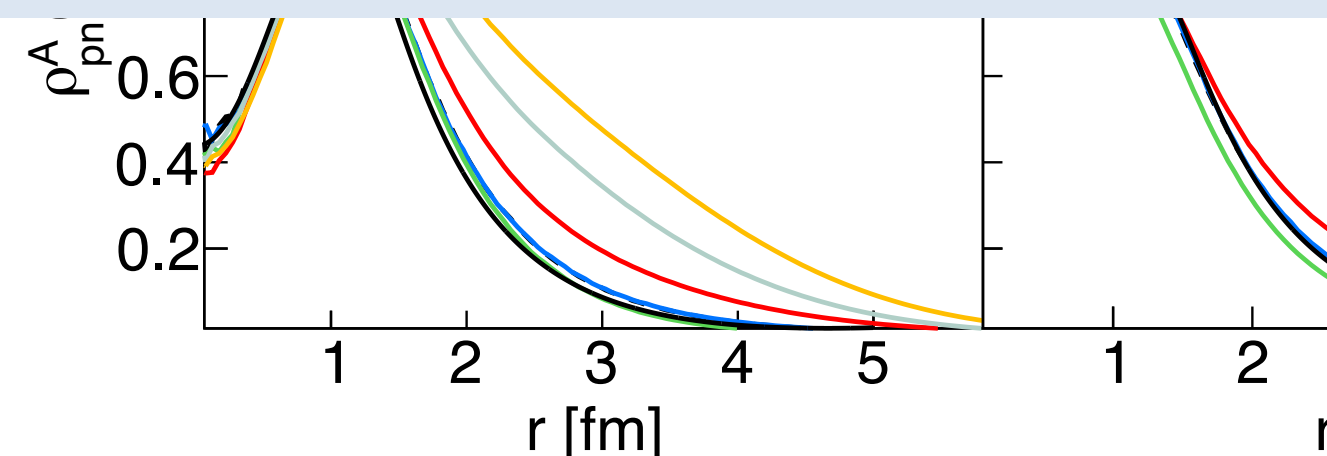
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$|\chi\rangle$

**SRC**

**EFT**

**20% High-p  
Tails**

**Q: How to reconcile/connect  
low-and high-resolution  
pictures?**

**Chiral  
Interactions**

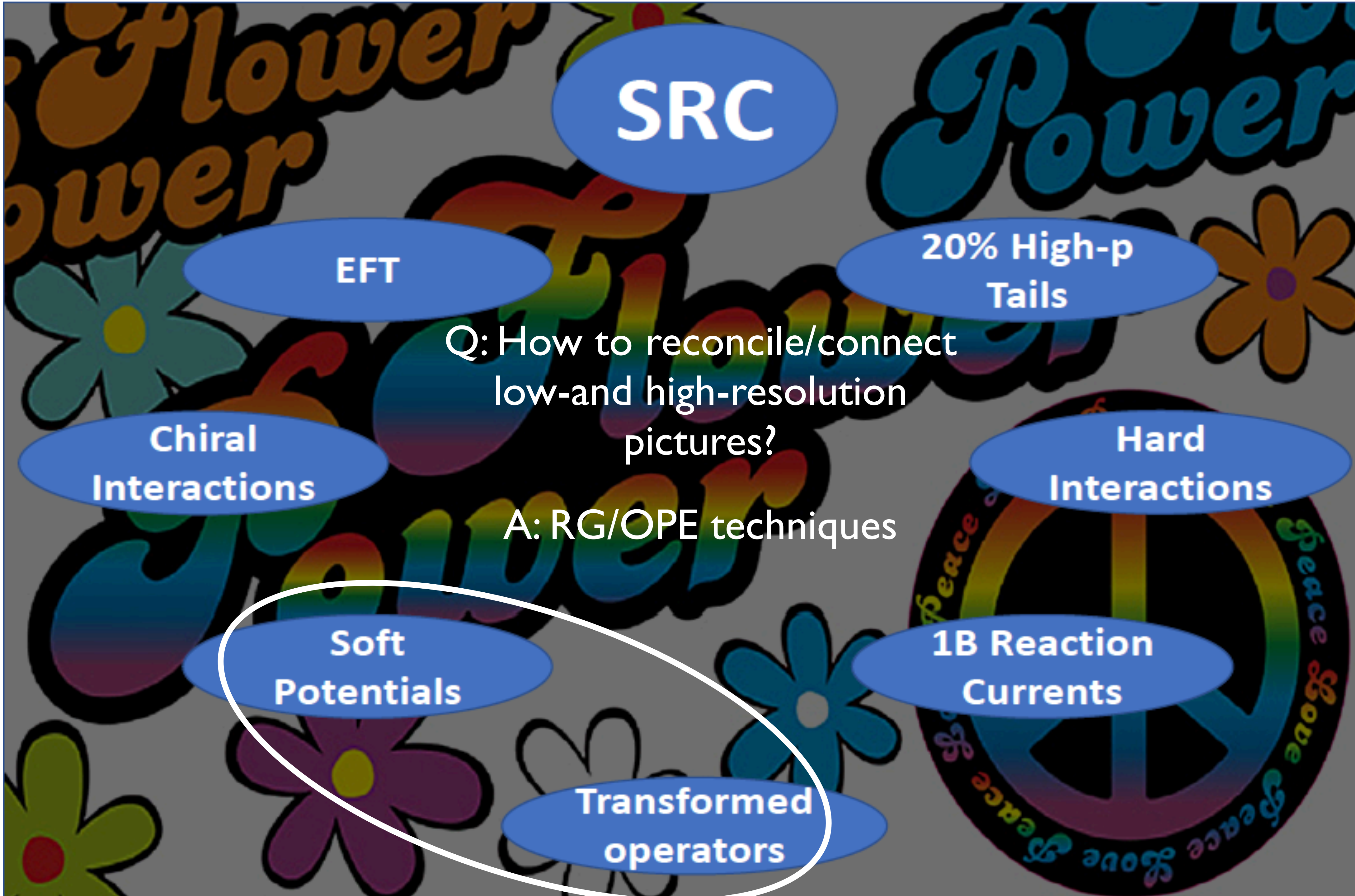
**Hard  
Interactions**

**A: RG/OPE techniques**

**Soft  
Potentials**

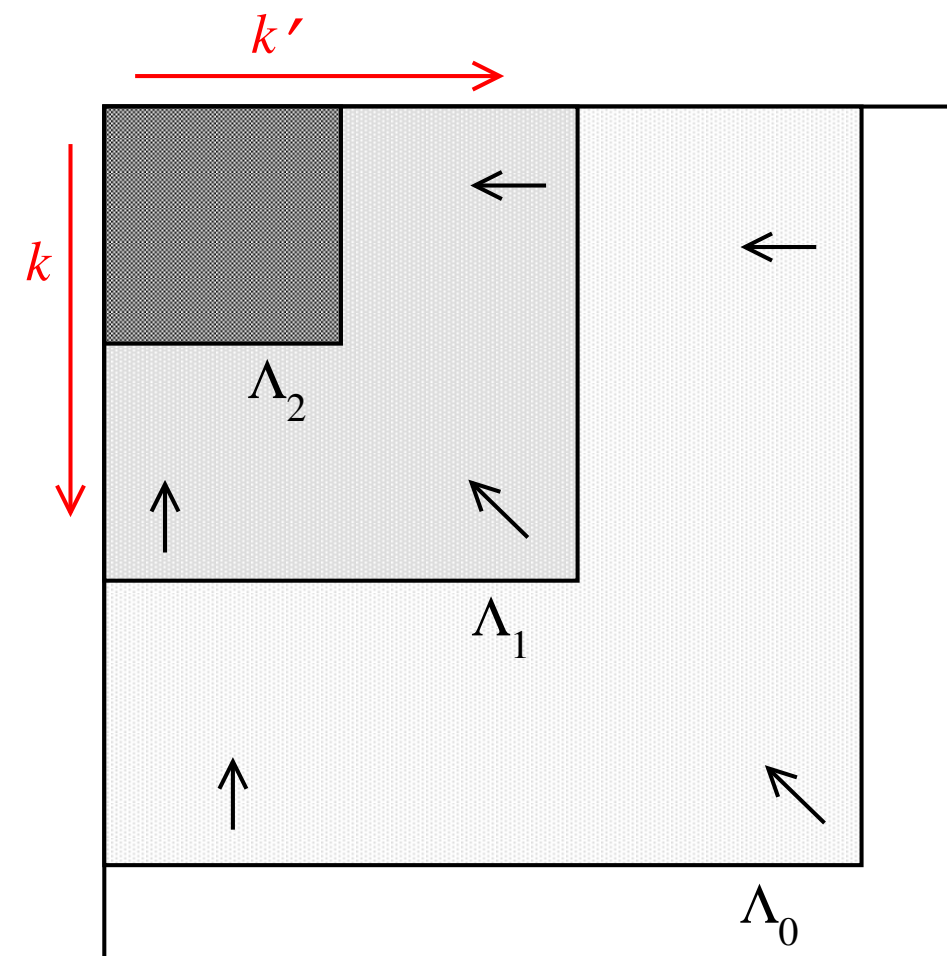
**1B Reaction  
Currents**

**Transformed  
operators**



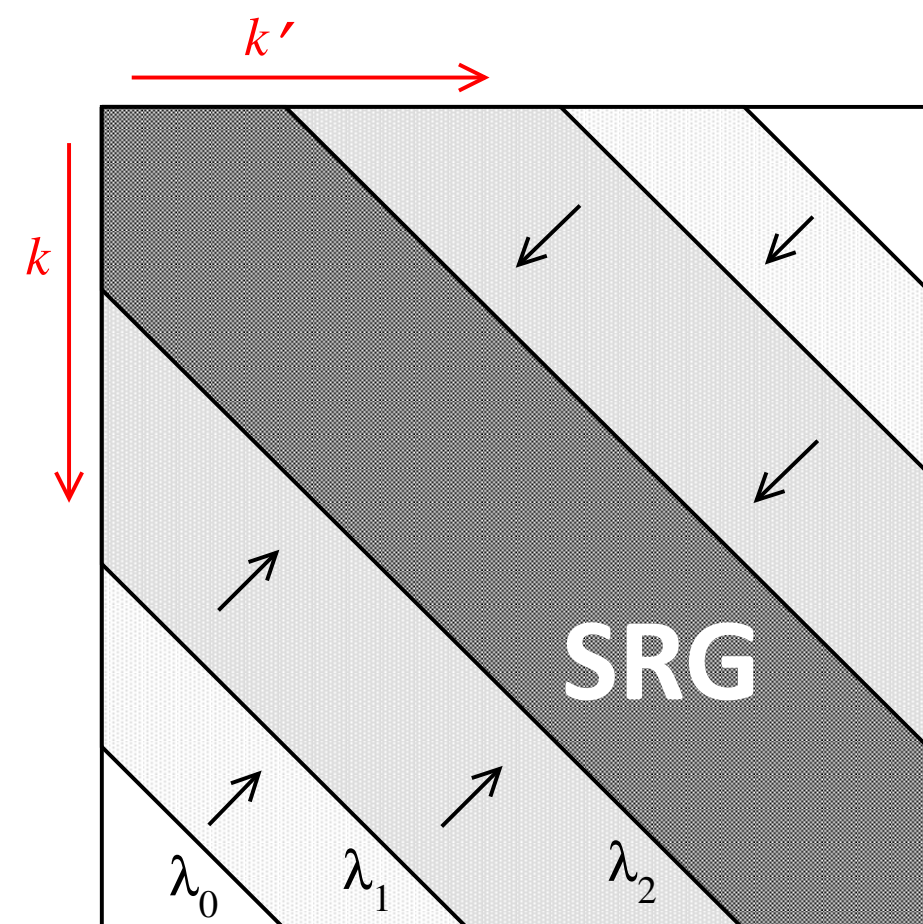
# RG in low energy nuclear physics

Bogner, Furnstahl, Schwenk Prog. Part. Nucl. Phys. 2010



Integrate out momenta  $k > \Lambda$

preserve physics **up to**  $\Lambda$



Unitary RG (“**S**imilarity **R**enormalization **G**roup”)

$$H(\lambda) = U(\lambda) H U^\dagger(\lambda) \quad O(\lambda) = U(\lambda) O U^\dagger(\lambda)$$

preserves all physics (unitary) if no approximations

low E states  $\Rightarrow k \gtrsim \lambda$  highly suppressed

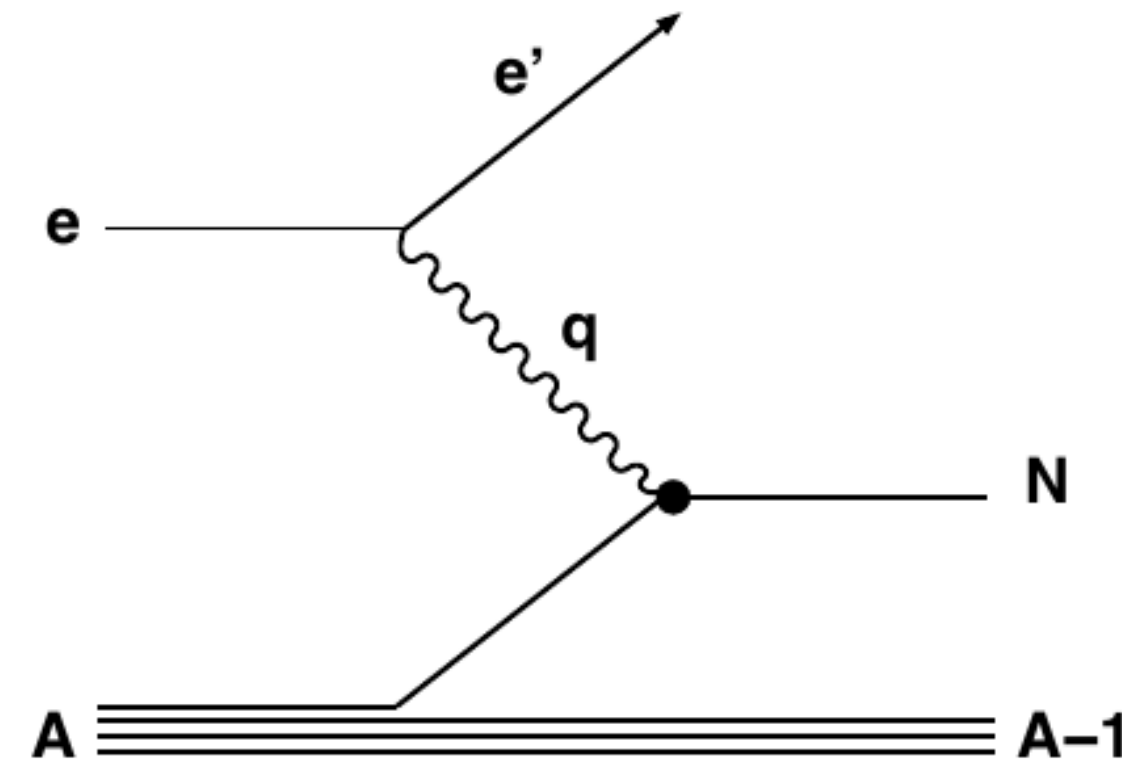


# Bridging structure and reactions

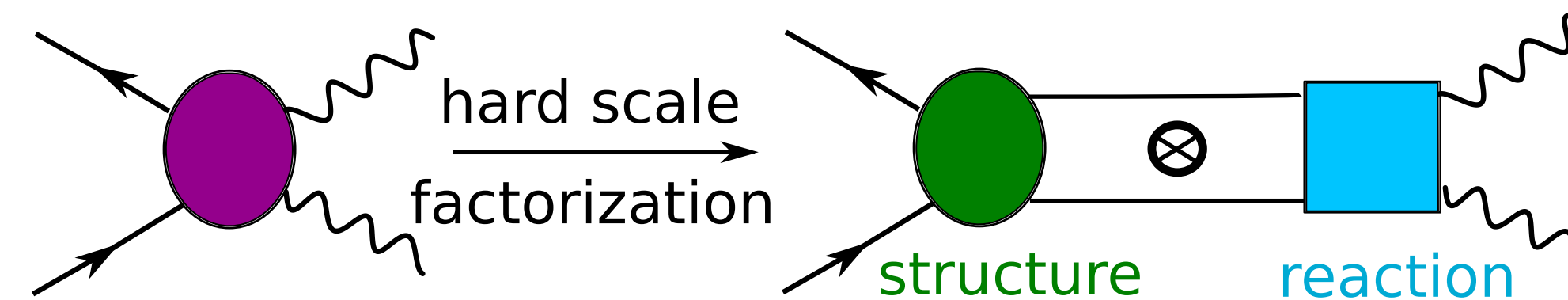
- Goal: Extract nuclear properties from experiments and predict them from theory

- $\frac{d\sigma}{d\Omega} \propto |\langle \psi_f | \hat{O}(q) | \psi_i \rangle|^2$

- Factorization to isolate components and extract process-independent properties



e.g., nucleon knockout reaction

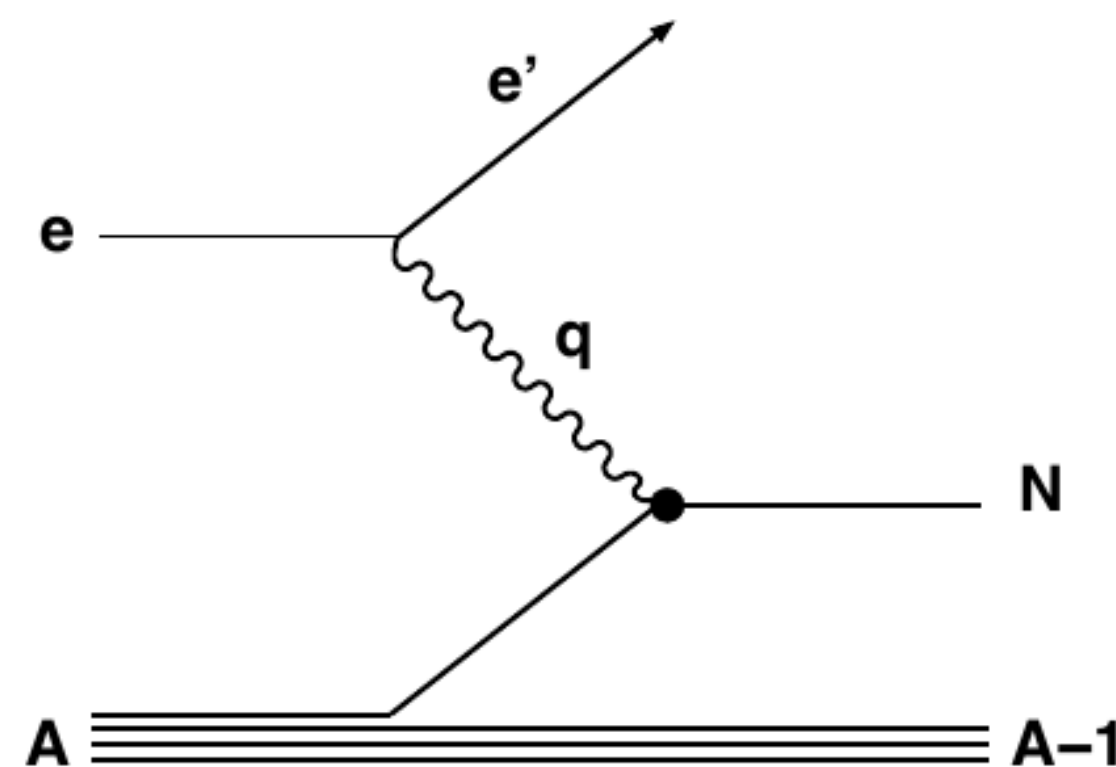


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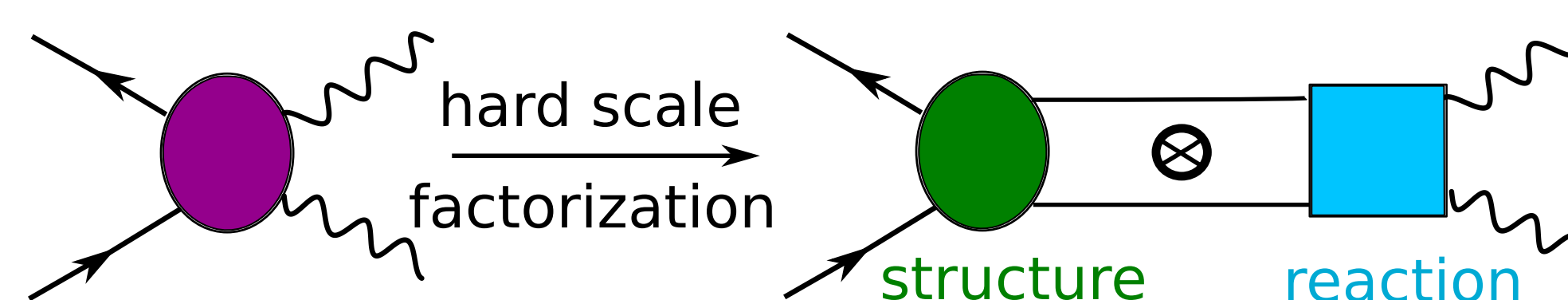
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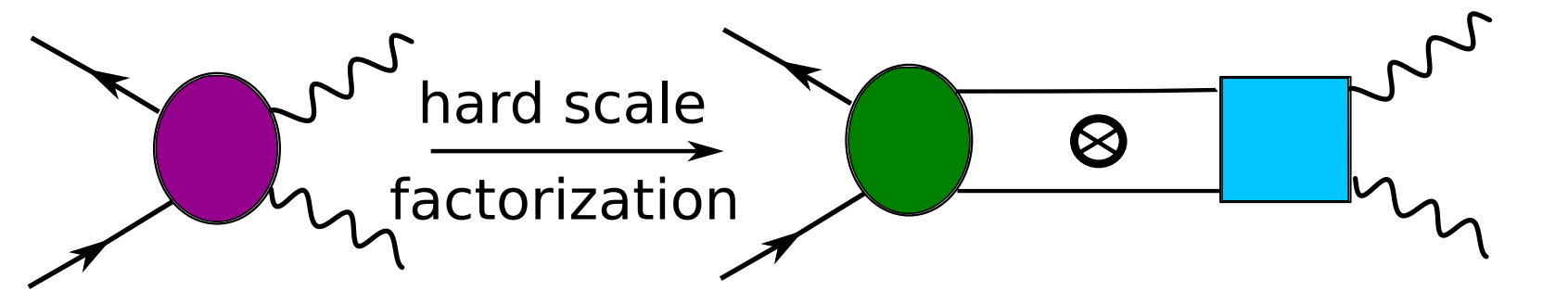
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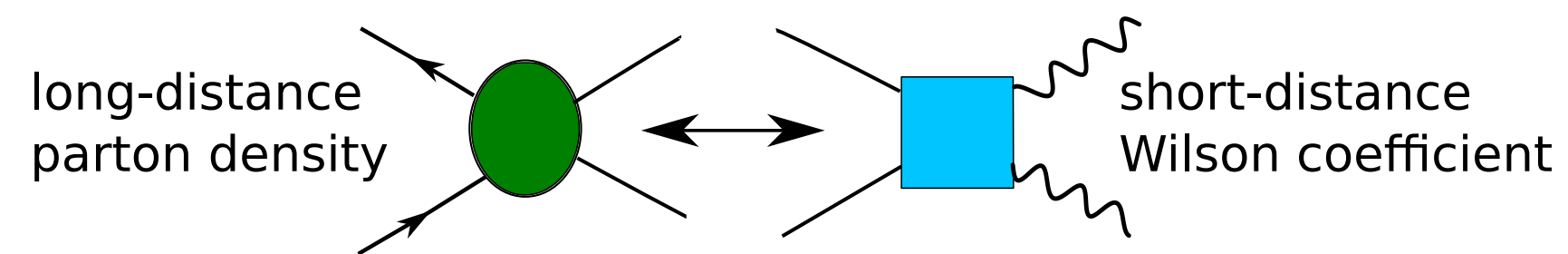
$$\underbrace{\langle \psi_f |}_{\text{structure}} \overbrace{\hat{O}(q)}^{\text{reaction}} \underbrace{|\psi_i \rangle}_{\text{structure}} = \langle \psi_f | U_\lambda U_\lambda^\dagger \hat{O}(q) U_\lambda U_\lambda^\dagger | \psi_i \rangle = \underbrace{\langle \psi_f^\lambda |}_{\text{structure}(\lambda)} \overbrace{\hat{O}^\lambda(q)}^{\text{reaction}(\lambda)} \underbrace{|\psi_i^\lambda \rangle}_{\text{structure}(\lambda)}$$

Factorization is scale-dependent (not unique)!!

## High-E QCD



$$F_2(x, Q^2) \sim \sum_a f_a(x, \mu_f) \otimes \hat{F}_2^a(x, Q/\mu_f)$$



## Low-E Nuclear

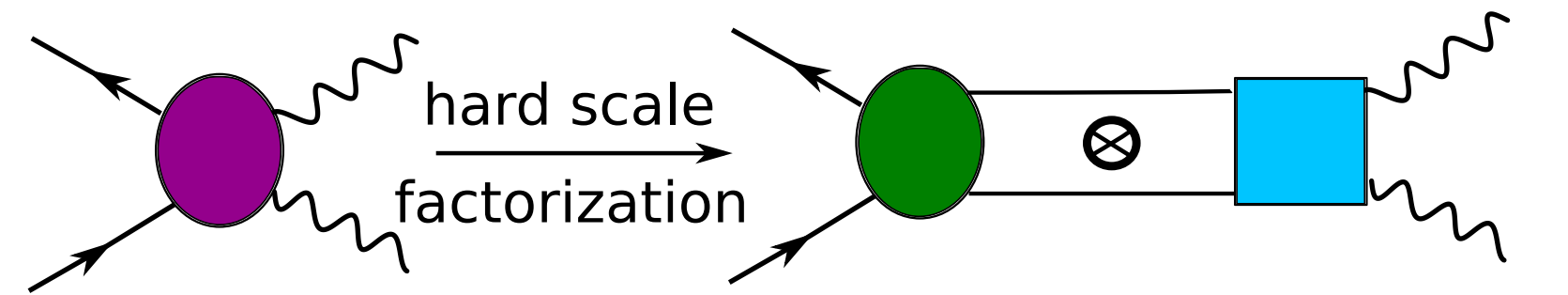
Observable: cross section

Structure model: spectroscopic factor

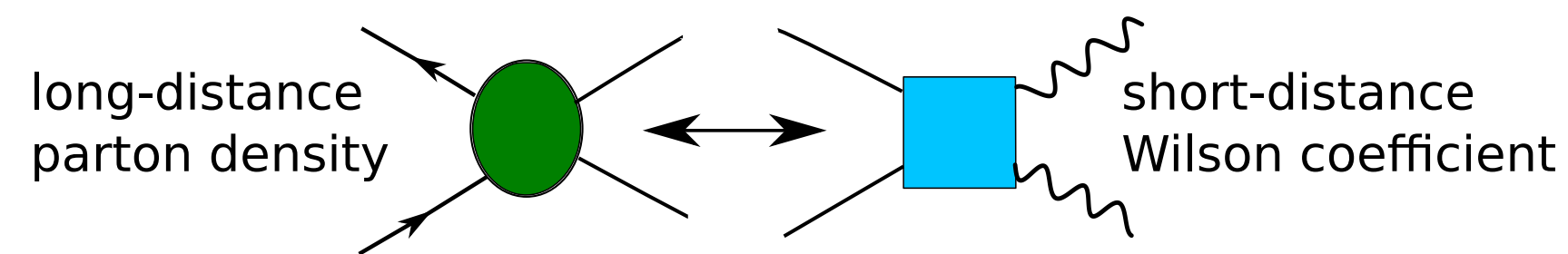
Reaction model: single-particle cross section

$$\sigma^{if} = \sum_{|J_i - J_f| \leq j \leq J_i + J_f} S_j^{if} \sigma_{sp}$$

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- Separation not unique, depends on the scale  $\mu_f$
- Form factor  $F_2$  independent of  $\mu_f$  but pieces not
- $f_a(x, \mu_f)$  runs with  $\mu_f^2 = Q^2$ , but is process independent

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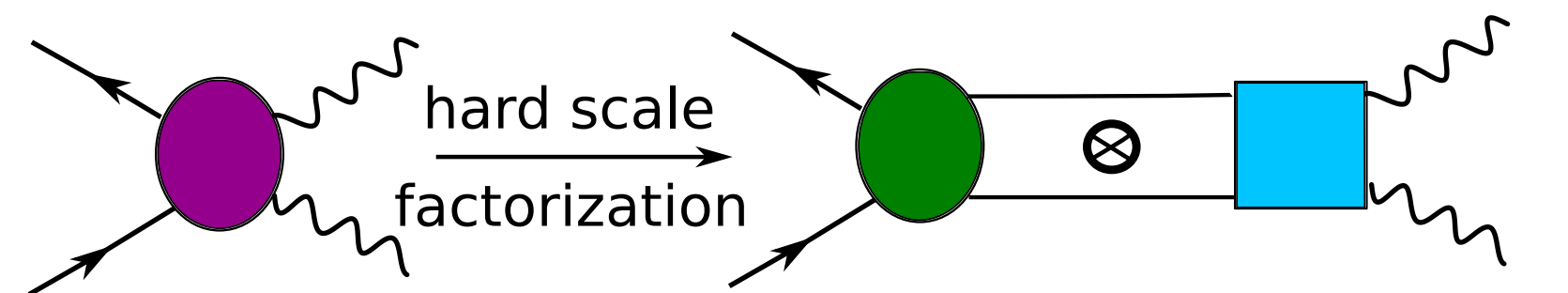
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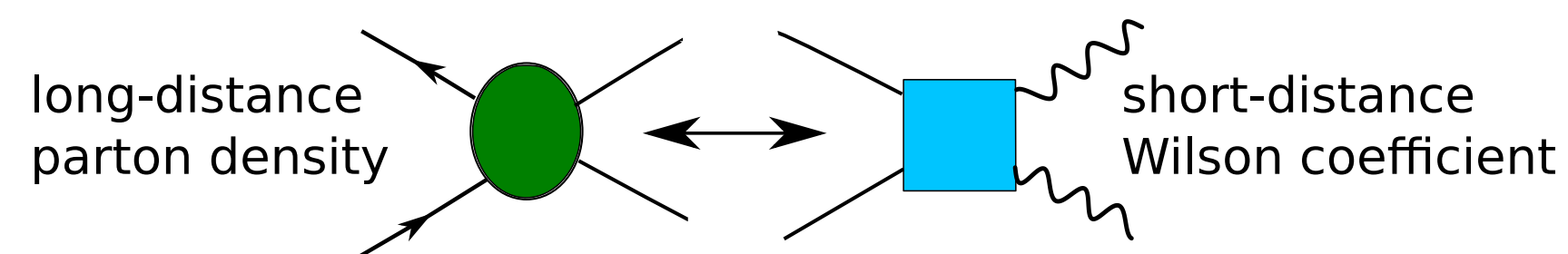
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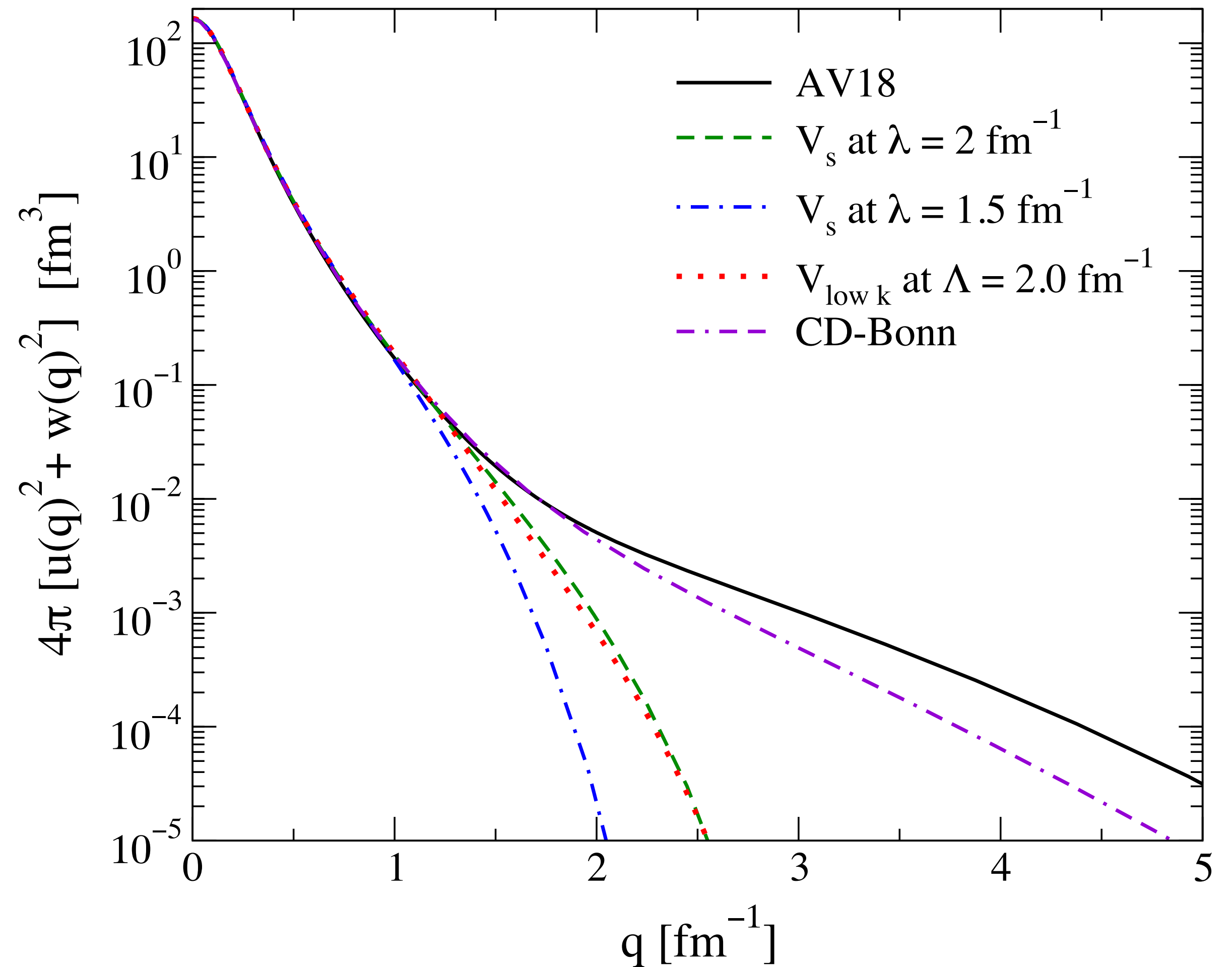
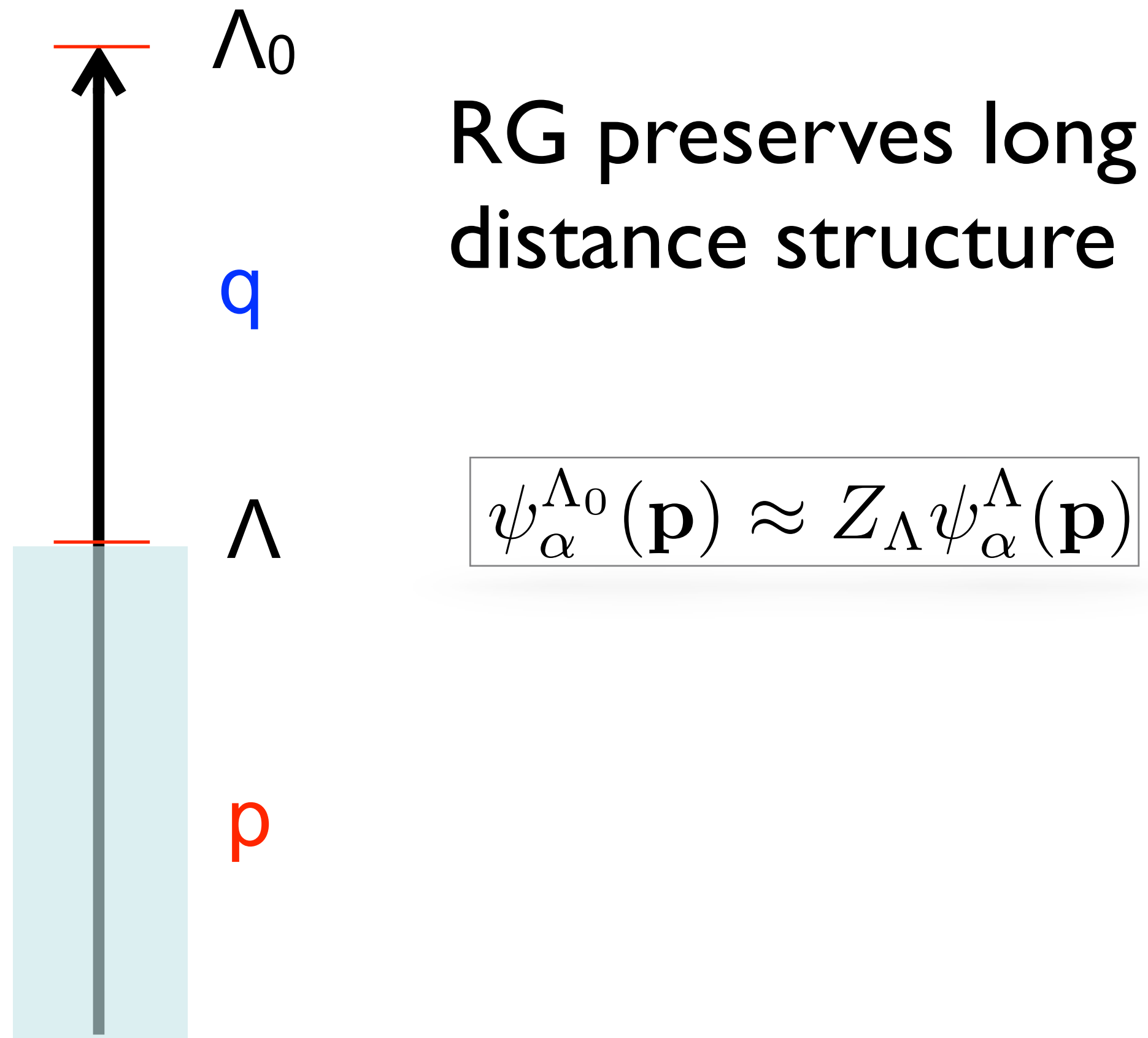
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## Open Questions

- What is the scale/scheme dependence of extracted props?
- Extract at one scale (e.g., to minimize FSI) and evolve to another?
- Scale/scheme dependence of interpretations? Are some better?
- **Structure of evolved operators?**

# High-q operators evolved to low-resolution scales

Consider **low-k** components of **low-E** wf's for  $A=2$ .

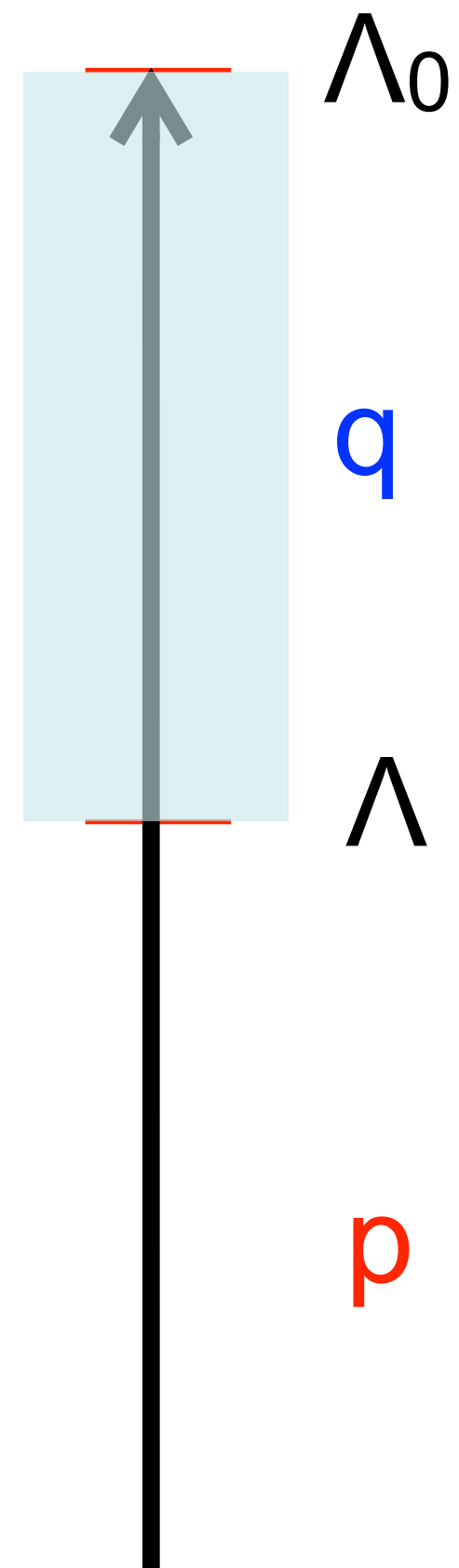


Anderson et al., PRC **82** (2010)

SKB and Roscher, PRC **86** (2012)

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Scale separation ( $E_\alpha \ll \Lambda^2 \ll q^2$ )

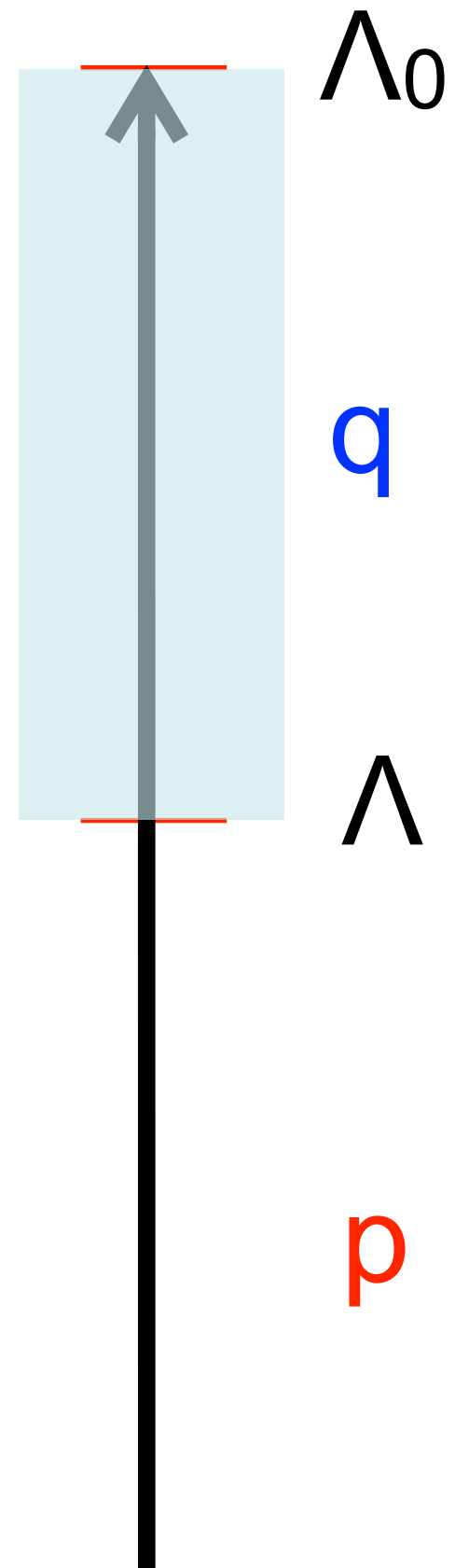
$$\psi_\alpha^{\Lambda_0}(\mathbf{q}) \approx \gamma(\mathbf{q}; \Lambda) \int_0^\Lambda d^3p Z_\Lambda \psi_\alpha^\Lambda(\mathbf{p}) + \eta(\mathbf{q}; \Lambda) \int_0^\Lambda d^3p \mathbf{p}^2 Z_\Lambda \psi_\alpha^\Lambda(\mathbf{p}) \dots$$

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**O**perator **P**roduct **E**xpansion  
of wave function a-la Lepage

$$\gamma(\mathbf{q}; \Lambda) = - \int_\Lambda^{\Lambda_0} d\mathbf{q}' \langle \mathbf{q} | \frac{1}{QH^{\Lambda_0}Q} | \mathbf{q}' \rangle V^{\Lambda_0}(\mathbf{q}', 0)$$

$$\beta(\mathbf{q}; \Lambda) = - \int_\Lambda^{\Lambda_0} d\mathbf{q}' \langle \mathbf{q} | \frac{1}{QH^{\Lambda_0}Q} | \mathbf{q}' \rangle \frac{\partial^2}{\partial p^2} V^{\Lambda_0}(\mathbf{q}', \mathbf{p}) \Big|_{\mathbf{p}=0}$$

State-independent  
Wilson Coefficients

Anderson et al., PRC **82** (2010)  
SKB and Roscher, PRC **86** (2012)

# High-q operators evolved to low-resolution scales



$$\begin{aligned} \langle \psi_{\alpha}^{\Lambda_0} | \widehat{O}_{\Lambda_0} | \psi_{\alpha}^{\Lambda_0} \rangle &= \int_0^{\Lambda} dp \int_0^{\Lambda} dp' \psi_{\alpha}^{\Lambda_0*}(p) O(p, p') \psi_{\alpha}^{\Lambda_0}(p') + \int_0^{\Lambda} dp \int_{\Lambda}^{\Lambda_0} dq \psi_{\alpha}^{\Lambda_0*}(p) O(p, q) \psi_{\alpha}^{\Lambda_0}(q) \\ &+ \int_{\Lambda}^{\Lambda_0} dq \int_0^{\Lambda} dp \psi_{\alpha}^{\Lambda_0*}(q) O(q, p) \psi_{\alpha}^{\Lambda_0}(p) + \int_{\Lambda}^{\Lambda_0} dq \int_{\Lambda}^{\Lambda_0} dq' \psi_{\alpha}^{\Lambda_0*}(q) O(q, q') \psi_{\alpha}^{\Lambda_0}(q') \end{aligned}$$

# High-q operators evolved to low-resolution scales



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Now use:

$$\psi_{\alpha}^{\Lambda_0}(\mathbf{q}) \approx \gamma(\mathbf{q}; \Lambda) \int_0^{\Lambda} d^3p Z_{\Lambda} \psi_{\alpha}^{\Lambda}(\mathbf{p}) + \dots \quad \text{OPE for w.f.'s}$$

$$\psi_{\alpha}^{\Lambda_0}(\mathbf{p}) \approx Z_{\Lambda} \psi_{\alpha}^{\Lambda}(\mathbf{p}) \quad \text{IR structure unaltered}$$

$$O(q, p) \approx O(q, 0) + \dots \quad \text{Scale separation}$$

# High-q operators evolved to low-resolution scales



$$\langle \psi_{\alpha}^{\Lambda_0} | \hat{O}_{\Lambda_0} | \psi_{\alpha}^{\Lambda_0} \rangle \approx Z_{\Lambda}^2 \langle \psi_{\alpha}^{\Lambda} | \hat{O}_{\Lambda_0} | \psi_{\alpha}^{\Lambda} \rangle + g^{(0)}(\Lambda) \langle \psi_{\alpha}^{\Lambda} | \delta^{(3)}(\mathbf{r}) | \psi_{\alpha}^{\Lambda} \rangle + \dots$$

state-independent  
high-q physics  
depends on operator

state dependent  
soft m.e. (low-k)  
same for all high-q operators

# High-q operators evolved to low-resolution scales

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state-independent  
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E.g.,

$$g^{(0)}(\Lambda) \equiv 2Z_{\Lambda}^2 \int_{\Lambda}^{\Lambda_0} d\tilde{q} O(0, q) \gamma(q; \Lambda) + Z_{\Lambda}^2 \int_{\Lambda}^{\Lambda_0} d\tilde{q} \int_{\Lambda}^{\Lambda_0} d\tilde{q}' \gamma^*(q; \Lambda) O(q, q') \gamma(q'; \Lambda)$$

Generically:

$$\hat{O}_{\Lambda} = Z_{\Lambda}^2 \hat{O}_{\Lambda_0} + g^{(0)}(\Lambda) \delta(\mathbf{r}) + g^{(2)}(\Lambda) \nabla^2 \delta(\mathbf{r}) + \dots$$



# High-q operators evolved to low-resolution scales



How does an operator that probes high-momentum w.f. components look in a low-momentum effective theory?

$$\langle \psi_{\alpha}^{\Lambda_0} | \hat{O}_{\Lambda_0} | \psi_{\alpha}^{\Lambda_0} \rangle \approx Z_{\Lambda}^2 \langle \psi_{\alpha}^{\Lambda} | \hat{O}_{\Lambda_0} | \psi_{\alpha}^{\Lambda} \rangle + g^{(0)}(\Lambda) \langle \psi_{\alpha}^{\Lambda} | \delta^{(3)}(\mathbf{r}) | \psi_{\alpha}^{\Lambda} \rangle + \dots$$

$= 0$  since  $P_{\Lambda} O_{\Lambda_0} P_{\Lambda} = 0$

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 = 0 since  $P_{\Lambda} O_{\Lambda_0} P_{\Lambda} = 0$

E.g., momentum distribution for  $q \gg \Lambda$

$$\langle \psi_{\alpha}^{\Lambda_0} | a_{\mathbf{q}}^{\dagger} a_{\mathbf{q}} | \psi_{\alpha}^{\Lambda_0} \rangle \approx \gamma^2(\mathbf{q}; \Lambda) Z_{\Lambda}^2 |\langle \psi_{\alpha}^{\Lambda} | \delta(\mathbf{r}) | \psi_{\alpha}^{\Lambda} \rangle|^2$$

low-E states have the same large-q tails if leading OPE term dominates

Generalize to arbitrary **A-body** states

# High-q operators evolved to low-resolution scales



Ex1: momentum distribution ( $\Lambda \ll q < \Lambda_0$ ):

$$\langle \psi_{\alpha,A}^{\Lambda_0} | a_{\mathbf{q}}^\dagger a_{\mathbf{q}} | \psi_{\alpha,A}^{\Lambda_0} \rangle \approx \gamma^2(\mathbf{q}; \Lambda) \times \sum_{\mathbf{k}, \mathbf{k}', \mathbf{K}}^{\Lambda} Z_{\Lambda}^2 \langle \psi_{\alpha,A}^{\Lambda} | a_{\frac{\mathbf{K}}{2} + \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}'} a_{\frac{\mathbf{K}}{2} + \mathbf{k}'} | \psi_{\alpha,A}^{\Lambda} \rangle$$

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Ex2: static structure factor ( $\Lambda \ll q < \Lambda_0$ ):

$$\langle \psi_{\alpha,A}^{\Lambda_0} | \widehat{S}(\mathbf{q}) | \psi_{\alpha,A}^{\Lambda_0} \rangle \approx \left\{ 2\gamma(\mathbf{q}; \Lambda) + \sum_{\mathbf{P}} \gamma(\mathbf{P} + \mathbf{q}; \Lambda) \gamma(\mathbf{P}; \Lambda) \right\} \\ \times \sum_{\mathbf{K}, \mathbf{k}, \mathbf{k}'}^{\Lambda} Z_{\Lambda}^2 \langle \psi_{\alpha,A}^{\Lambda} | a_{\frac{\mathbf{K}}{2} + \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}'} a_{\frac{\mathbf{K}}{2} + \mathbf{k}'} | \psi_{\alpha,A}^{\Lambda} \rangle$$

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- hard (high q) physics
- Universal (state-indep)
- depends on probe
- fixed from few-body

**X**

- soft (low-k) m.e.
- same for all high-q probes
- A-dependent scale factor

# High- $q$ operators evolved to low-resolution scales



links few- and  $A$ -body systems (“derives” the GCF)

Correlations/scaling for 2 observables w/same leading OPE

Subleading OPE  $\Rightarrow$  deviations from scaling calculable in principle?

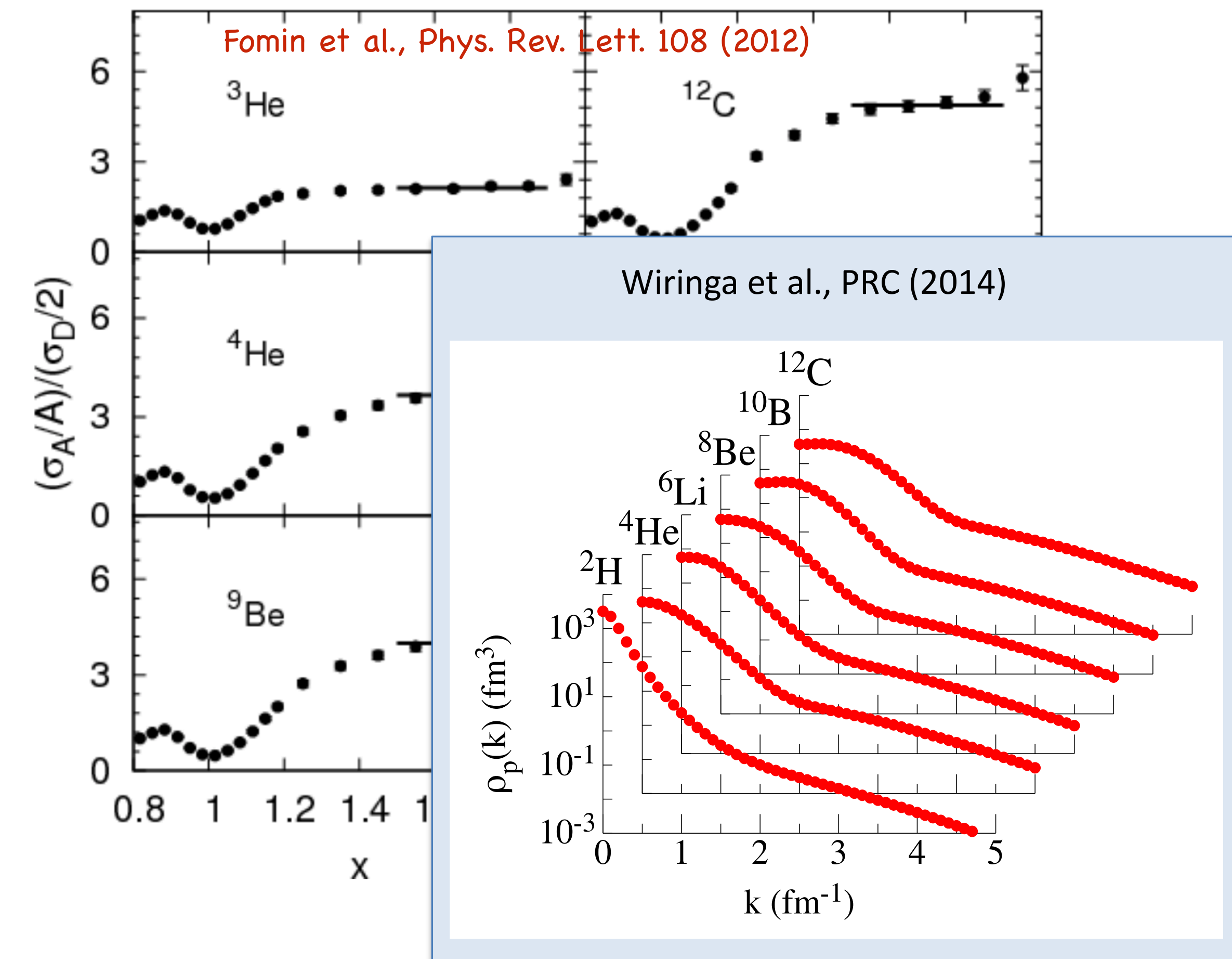
- fixed from few-body

# SRC phenomenology revisited (low-res picture)

Tropiano, SKB, Furnstahl (in progress)

## I) Universal high-momentum tails

inclusive ratios

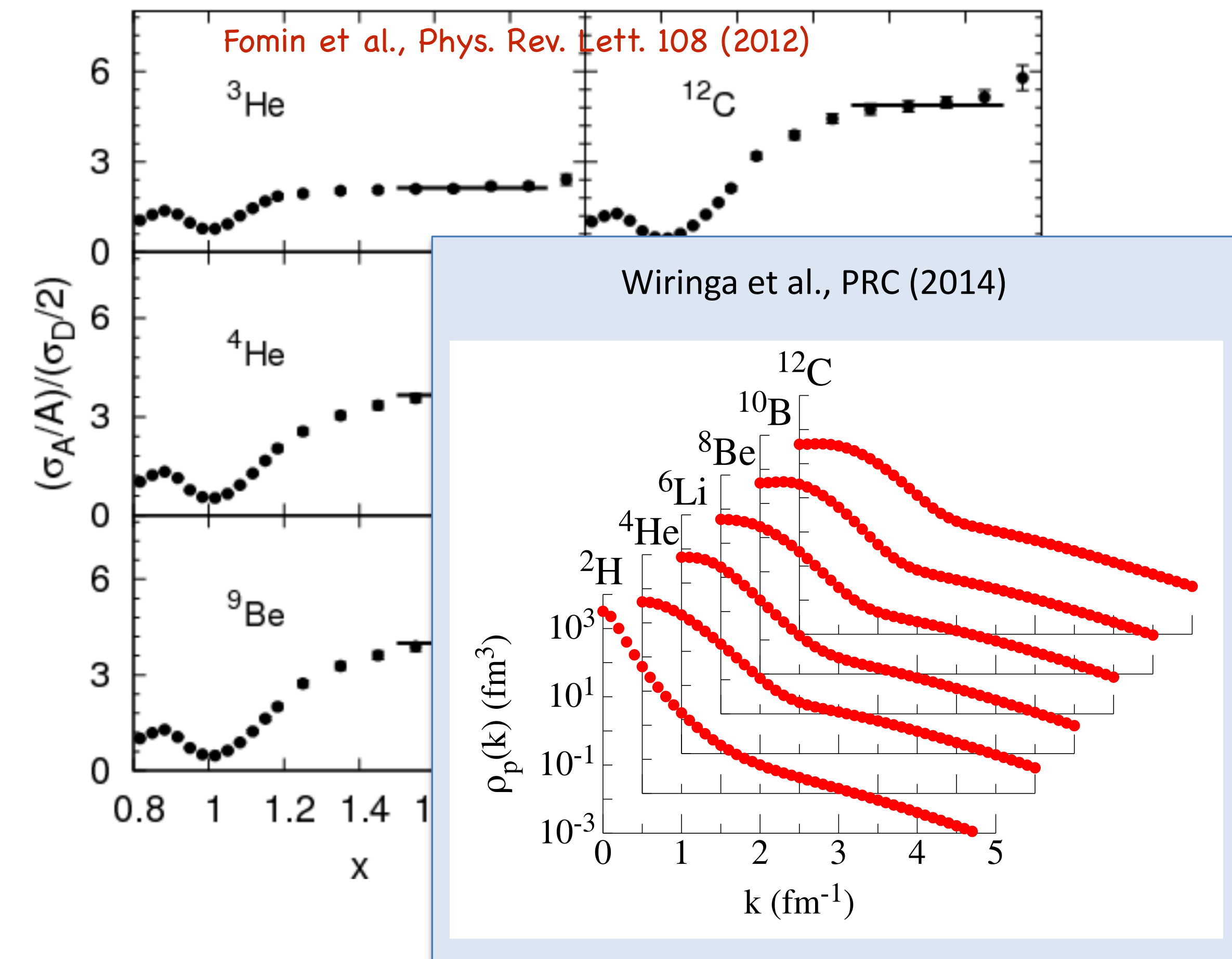


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$$\langle \psi_{\alpha,A}^{\Lambda_0} | a_{\mathbf{q}}^\dagger a_{\mathbf{q}} | \psi_{\alpha,A}^{\Lambda_0} \rangle \approx \gamma^2(\mathbf{q}; \Lambda) \times \sum_{\mathbf{k}, \mathbf{k}', \mathbf{K}}^{\Lambda} Z_{\Lambda}^2 \langle \psi_{\alpha,A}^{\Lambda} | a_{\frac{\mathbf{K}}{2} + \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}'} a_{\frac{\mathbf{K}}{2} + \mathbf{k}'} | \psi_{\alpha,A}^{\Lambda} \rangle$$

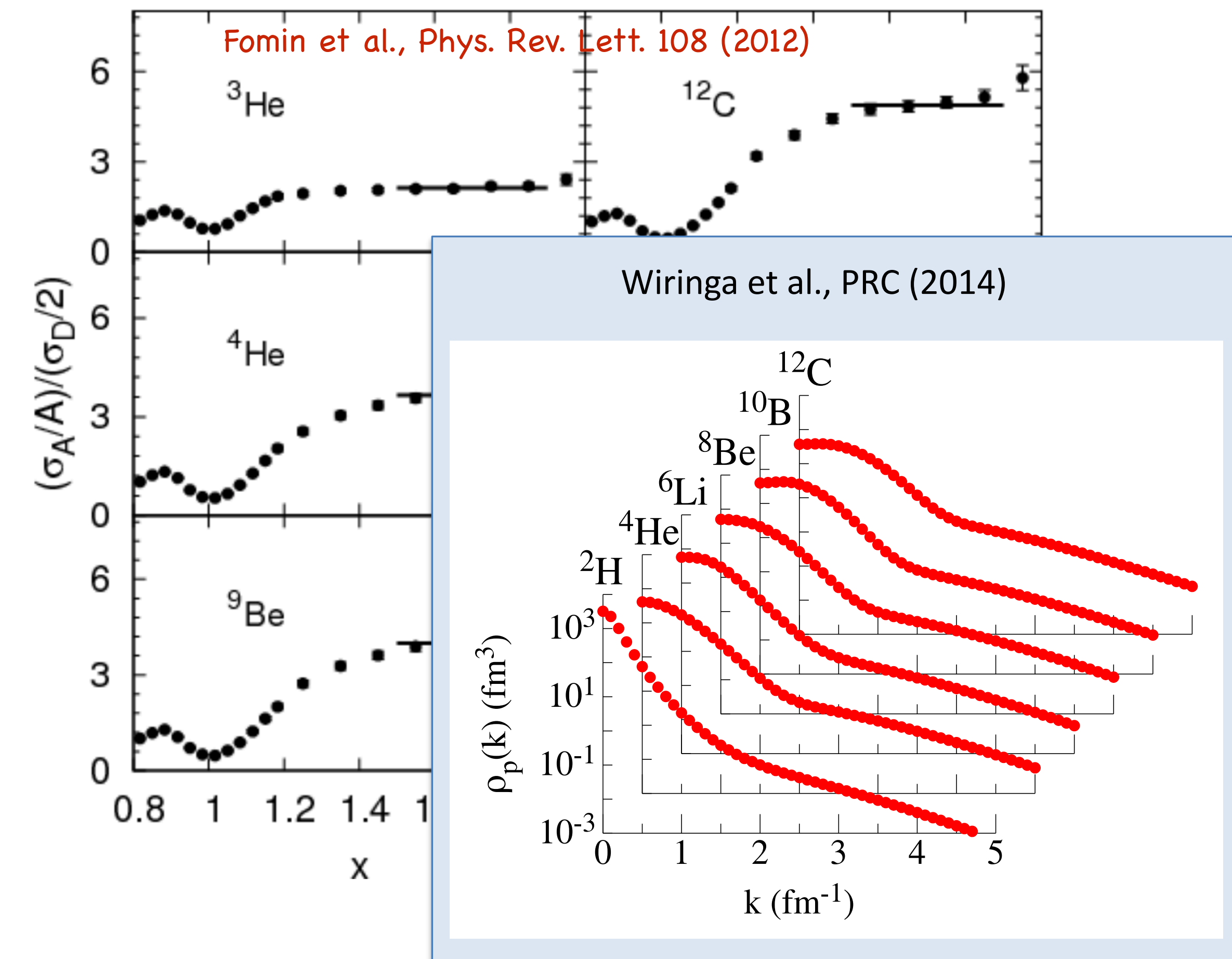


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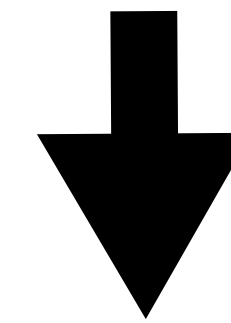
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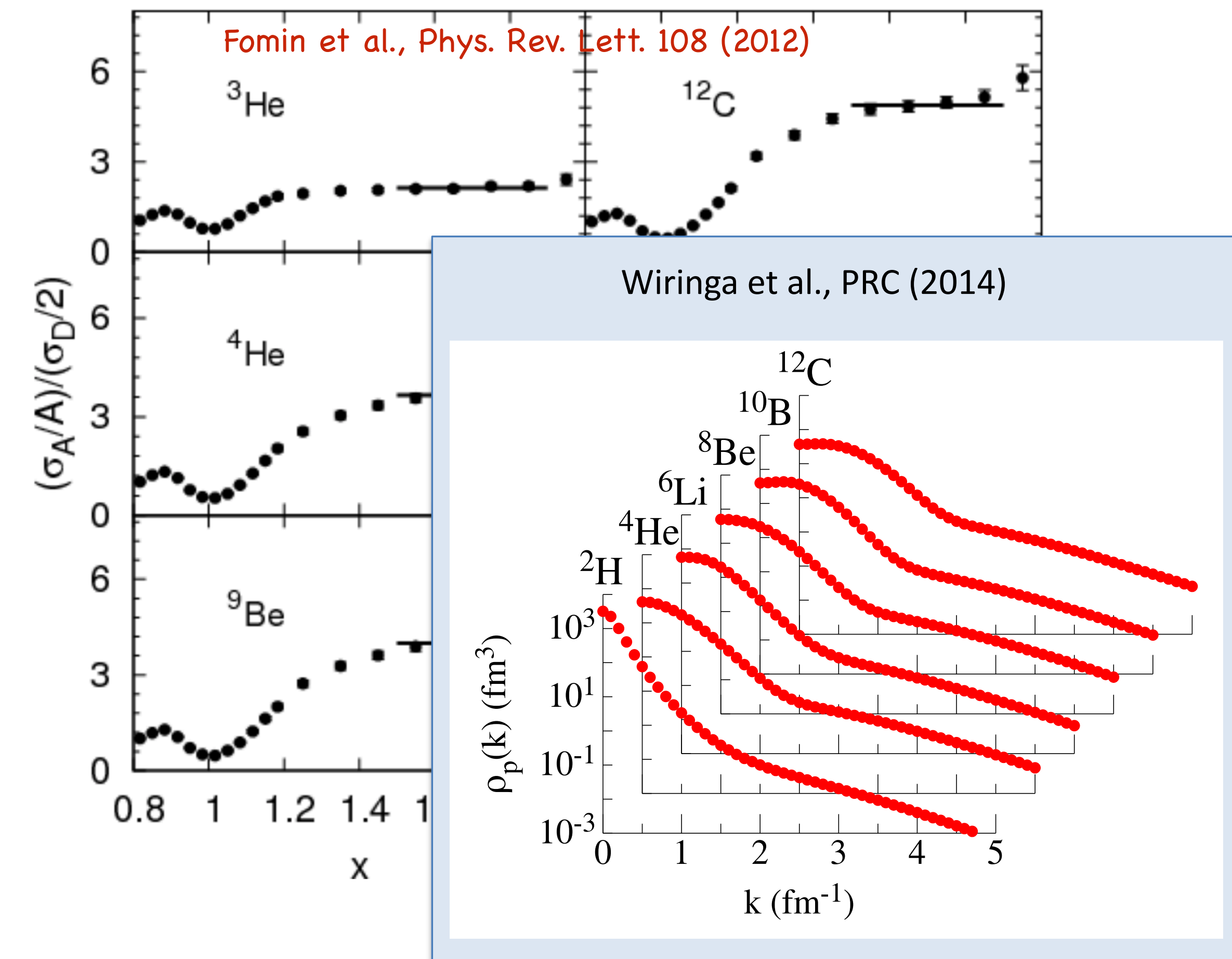
$$C(A, 2) \equiv \frac{n_A(\mathbf{q})}{n_D(\mathbf{q})} \sim \frac{\sum_{\mathbf{k}, \mathbf{k}', \mathbf{K}} \langle \psi_{\alpha,A}^\Lambda | a_{\frac{\mathbf{K}}{2} + \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}'} a_{\frac{\mathbf{K}}{2} + \mathbf{k}'} | \psi_{\alpha,A}^\Lambda \rangle}{\sum_{\mathbf{k}, \mathbf{k}', \mathbf{K}} \langle \psi_{\alpha,D}^\Lambda | a_{\frac{\mathbf{K}}{2} + \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}}^\dagger a_{\frac{\mathbf{K}}{2} - \mathbf{k}'} a_{\frac{\mathbf{K}}{2} + \mathbf{k}'} | \psi_{\alpha,D}^\Lambda \rangle}$$

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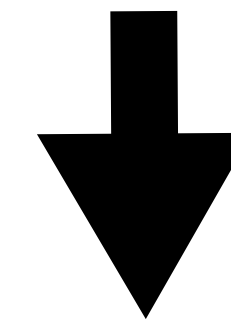
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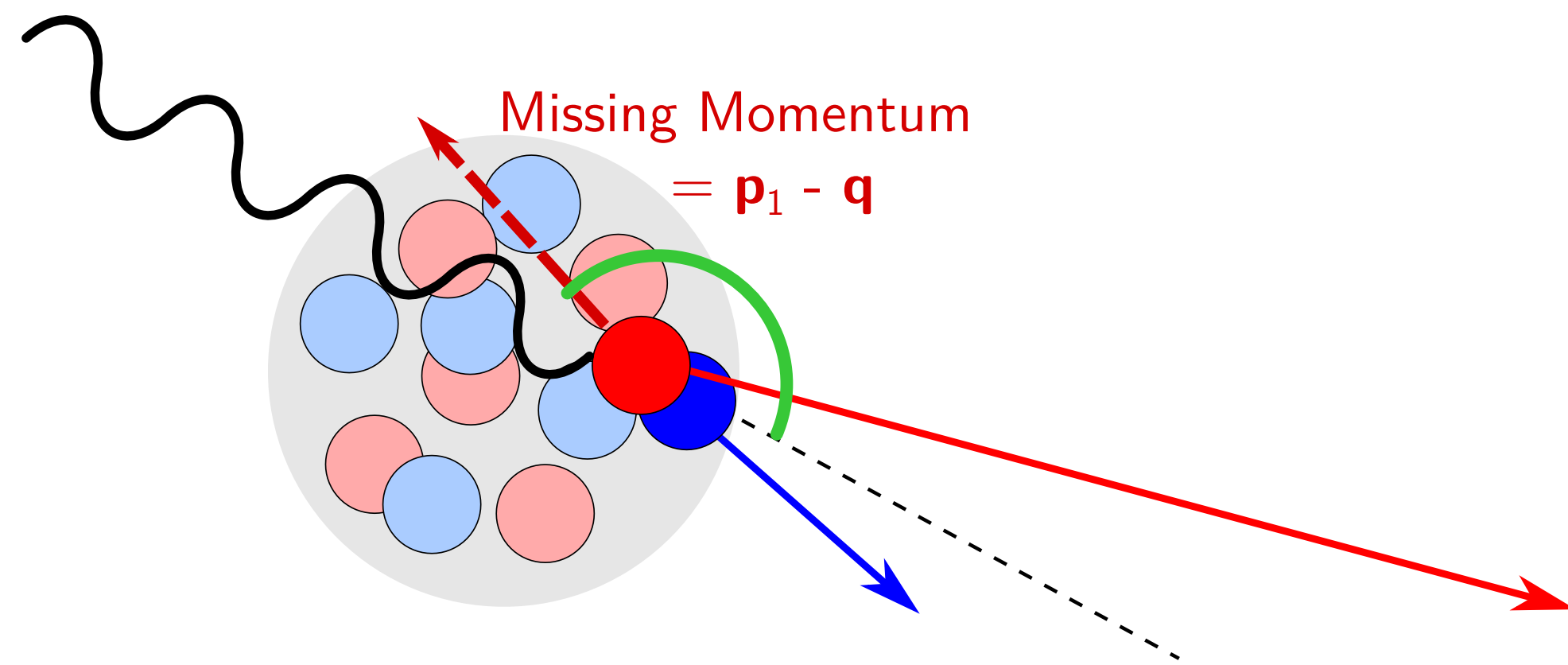


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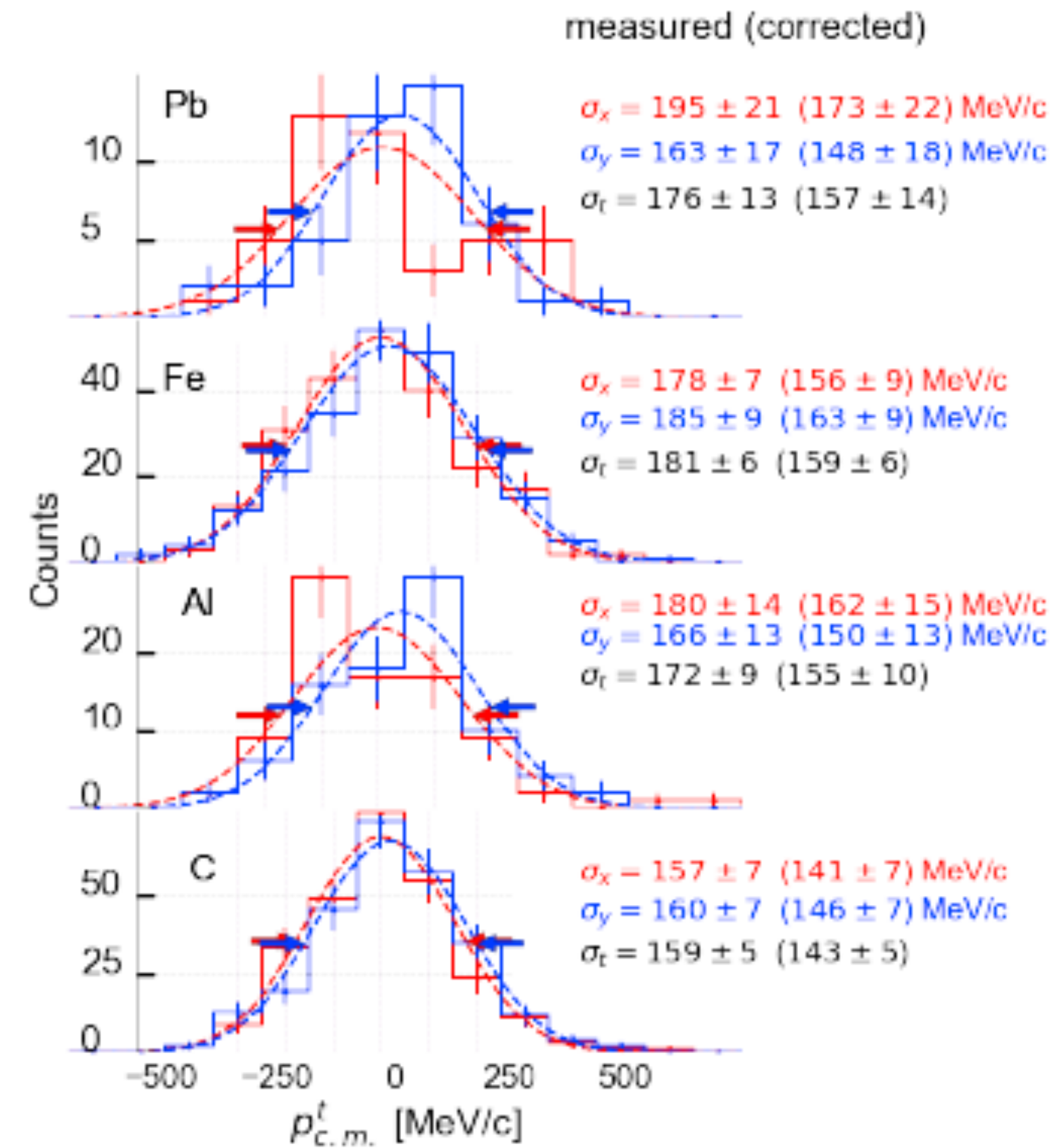
low-momentum/mean field physics => scale/scheme indeed to leading order

supports/explains GCF conclusions about inclusive ratios

## 2) Kinematics of knocked-out nucleons



knocked out SRC nucleons fly out almost back-to-back (relative s-wave pairs)



pair CM momentum distribution gaussian of width  $\sim k_F$

## 2) Kinematics

evolved pair momentum distribution ( $\lambda \sim k_F \ll q$ )

$$\rho_{NN,\alpha}(Q, q) \sim \gamma_\alpha^2(q; \Lambda) \sum_{k,k'} |\langle \psi^A(\Lambda) | [a_{\frac{Q}{2}+k}^\dagger a_{\frac{Q}{2}-k}^\dagger a_{\frac{Q}{2}-k'} a_{\frac{Q}{2}+k'}]_\alpha | \psi^A(\Lambda) \rangle$$

known  
almost  
(re

solution

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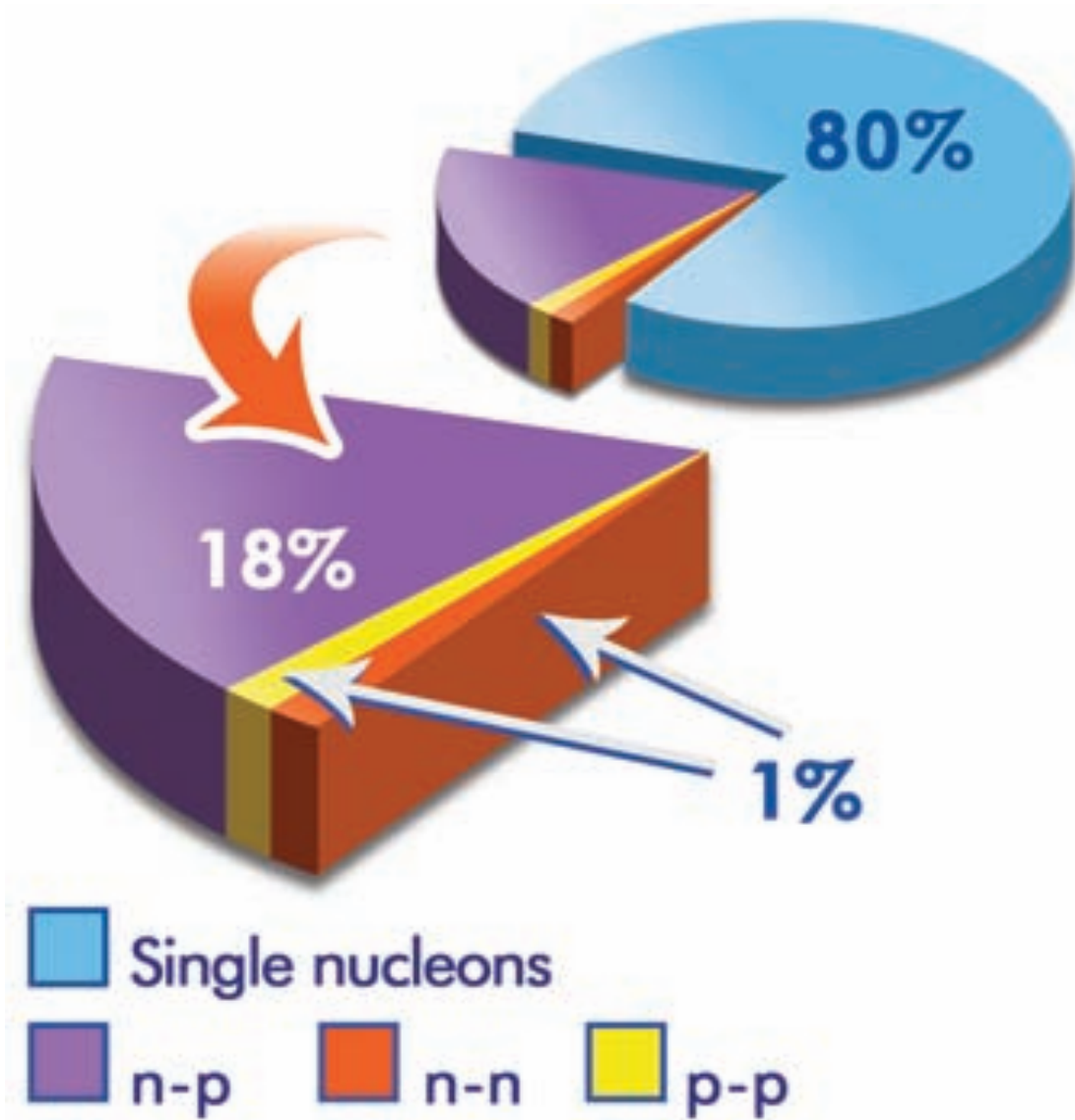
m.e. of smeared contact operator  $\implies$   
 high  $q$  pairs dominated relative s-waves

evolved  $\psi(\Lambda)$  “soft”, dominated by MFT configs  $\implies$   
 CM  $Q$  distribution smooth/gaussian with width  $\sim k_F$

know  
 alm  
 (re

olution

## 3) np dominance at intermediate (300-500 MeV) relative momenta

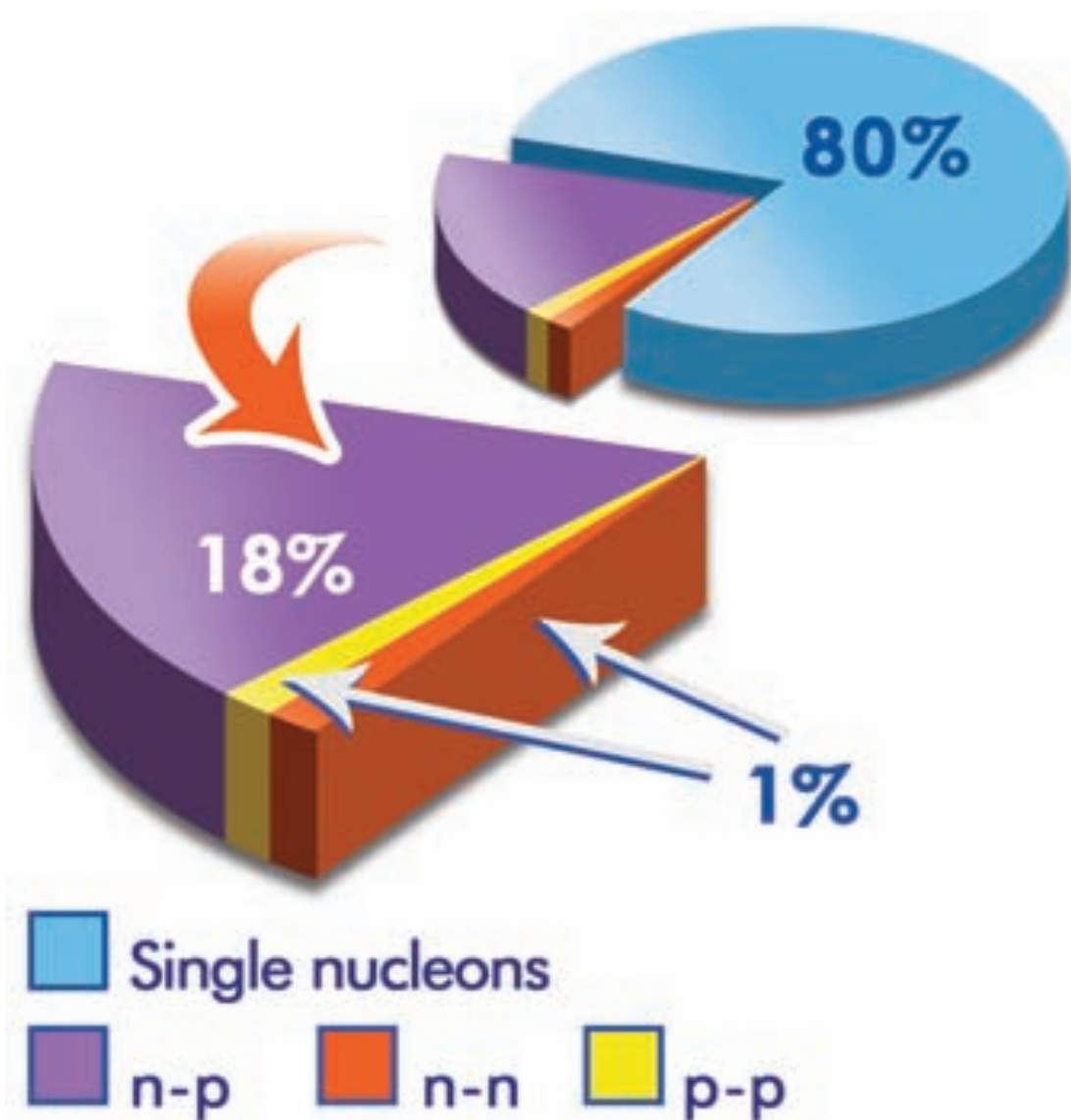


**Fig. 3.** The average fraction of nucleons in the various initial-state configurations of  $^{12}\text{C}$ .

R. Subedi et al., Science (2008)

20% of nucleons in SRC pairs  
but mostly neutron-proton

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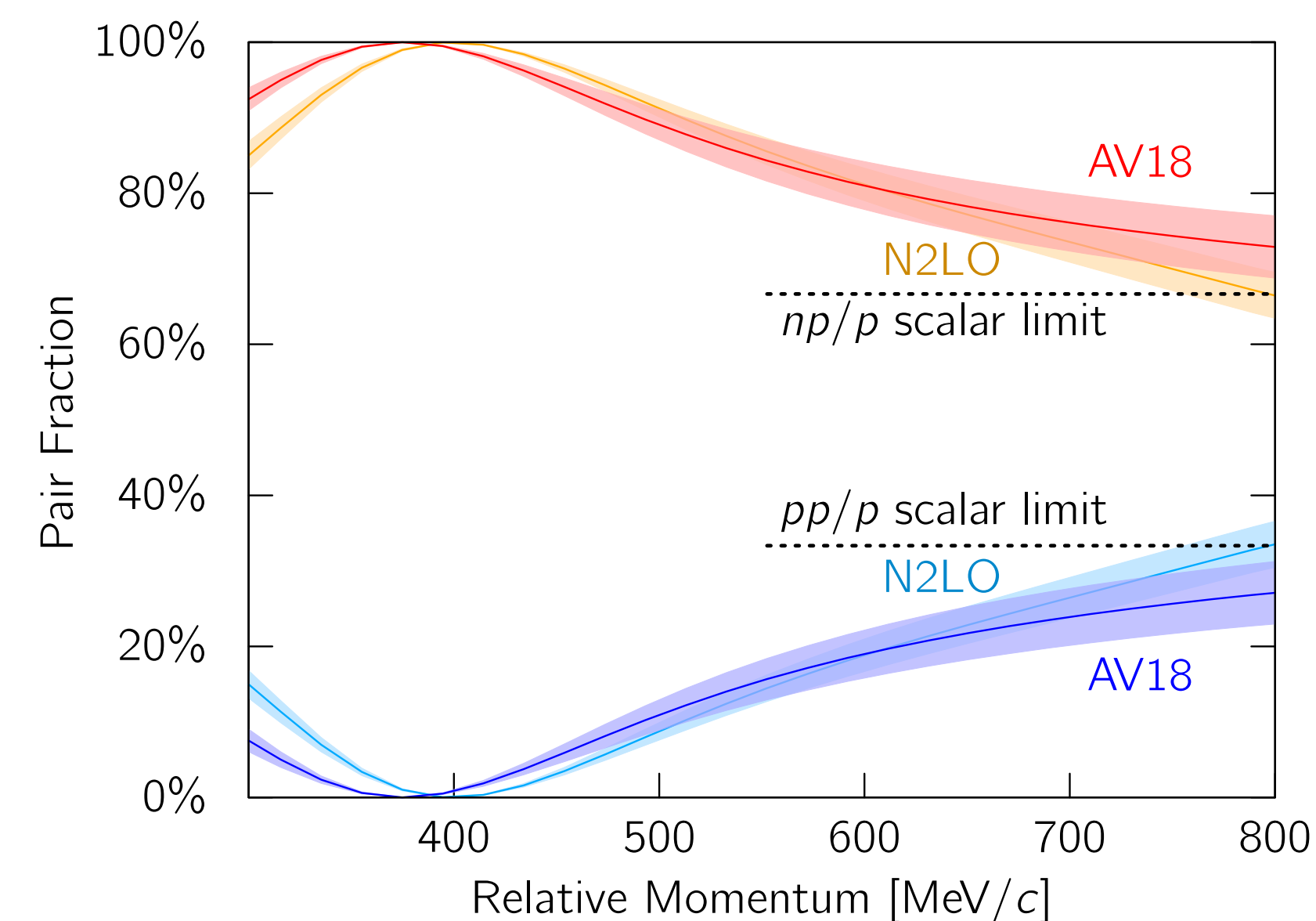


**Fig. 3.** The average fraction of nucleons in the various initial-state configurations of  $^{12}\text{C}$ .

## 4) transition to scalar counting at higher relative momentum

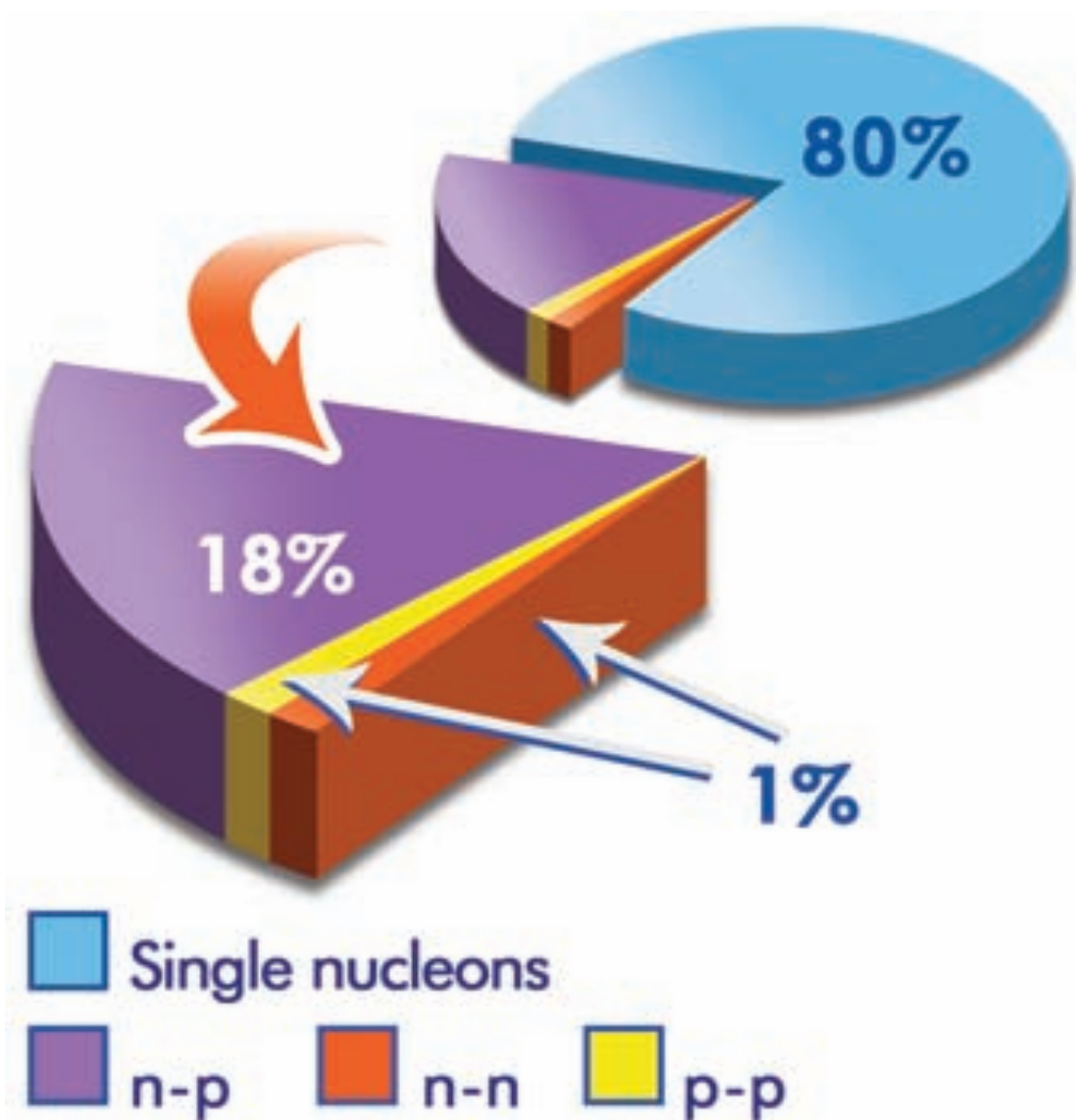
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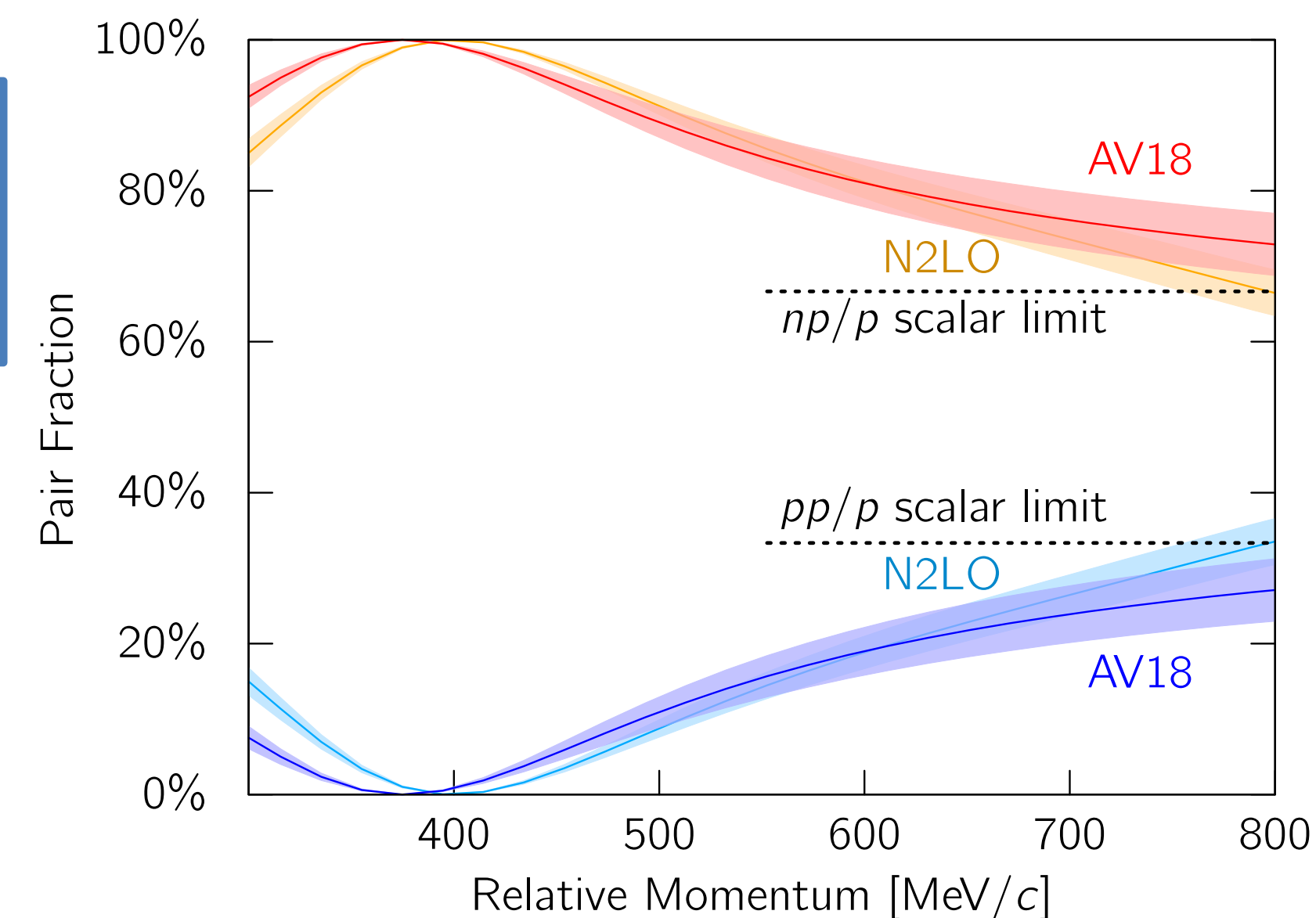
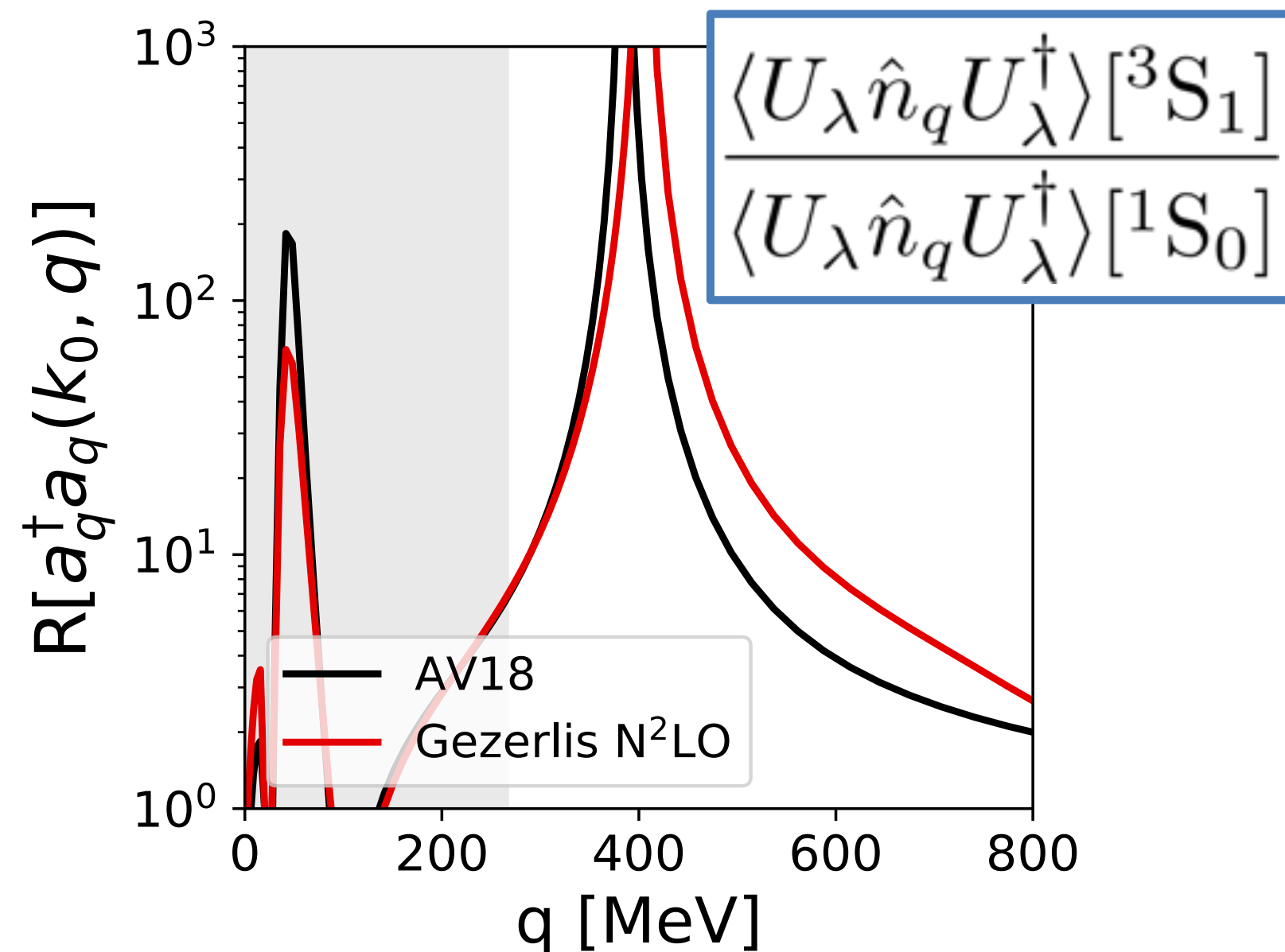
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Ratio of *evolved* high-mom. distributions in a low-mom. state (insensitive to details!)





## 6) Generalized Contact Formalism (GCF)

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$$\rho_A^{NN,\alpha}(r) = C_A^{NN,\alpha} \times |\varphi_{NN}^\alpha(r)|^2$$

$$n_A^{NN,\alpha}(q) = C_A^{NN,\alpha} \times |\varphi_{NN}^\alpha(q)|^2$$

A-dep scale factors (“nuclear contacts”)  $C_A \sim \langle \chi | \chi \rangle$

Universal (same all A, **not**  $V_{NN}$ ) shape from  
two-body zero energy wf  $\phi$

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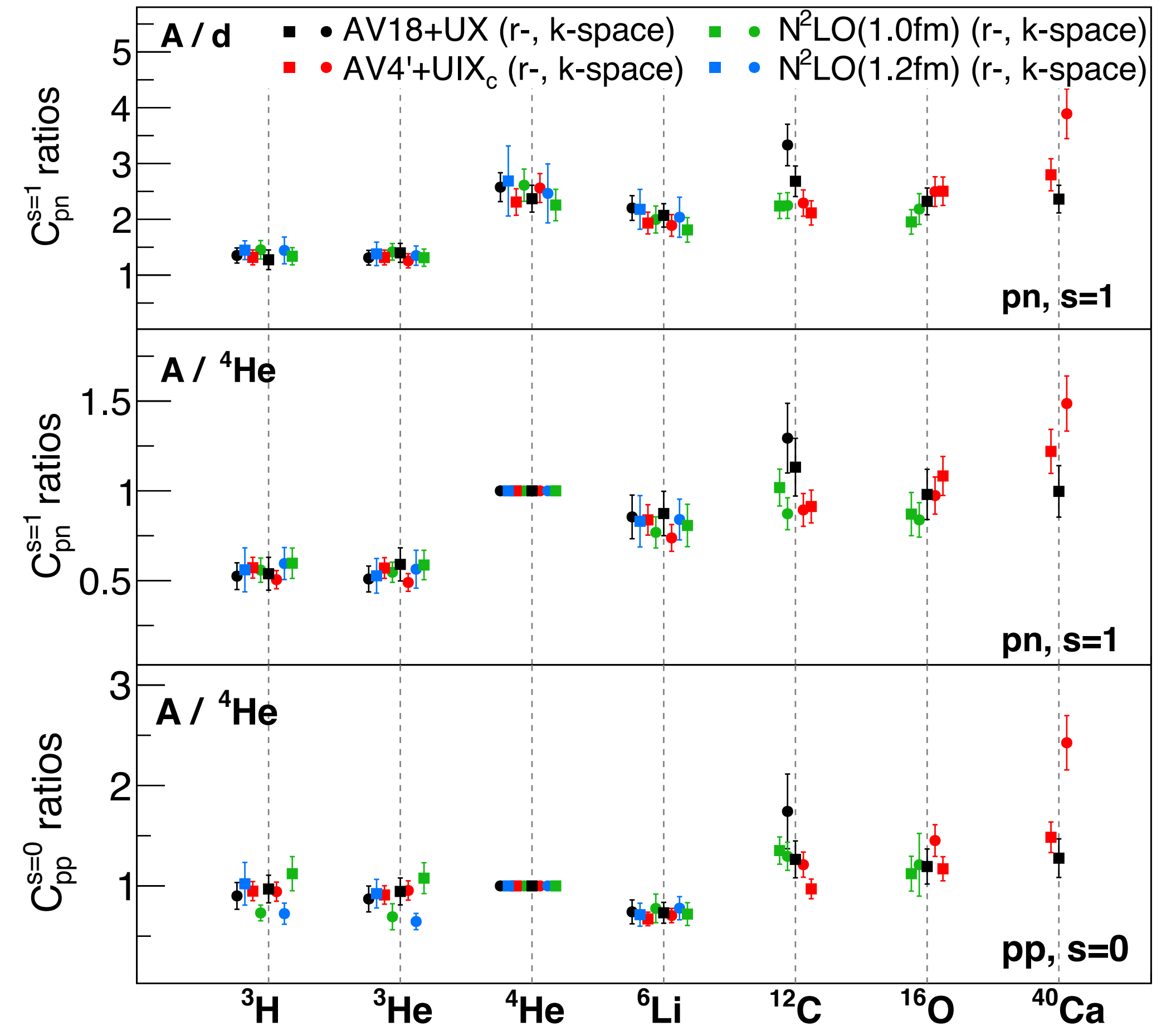
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**But  $\varphi_{NN}$  is scale and scheme dependent. Ratios are independent but only probe “mean field” part**



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## 6) General

$\rho_A^{NN,\alpha}$   
 $n_A^{NN,\alpha}$

A-dep scale

Universal (s  
two-body z

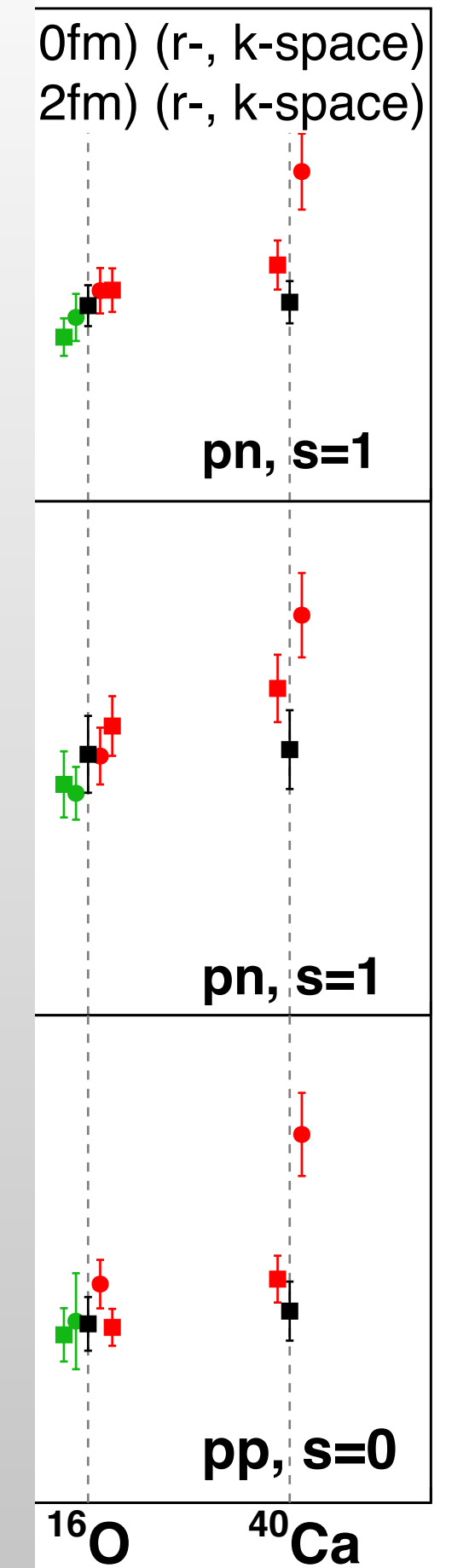
But  
schem  
are in  
probe

### Contacts **not** RG invariant

$$C_A = \sum_{K,k',k}^{\Lambda_0} \langle \psi_{\Lambda_0}^A | a_{\frac{K}{2}+k}^\dagger a_{\frac{K}{2}-k}^\dagger a_{\frac{K}{2}-k'} a_{\frac{K}{2}+k'} | \psi_{\Lambda_0}^A \rangle \Rightarrow f(\Lambda) \sum_{K,k',k}^{\Lambda} \langle \psi_{\Lambda}^A | a_{\frac{K}{2}+k}^\dagger a_{\frac{K}{2}-k}^\dagger a_{\frac{K}{2}-k'} a_{\frac{K}{2}+k'} | \psi_{\Lambda}^A \rangle$$

A-independent

...But ratios in different A approx. RG invariant



# Scale dependence of deuteron electrodisintegration

S. More, SKB, R.J. Furnstahl, Phys. Rev. C **96** 054004, (2017)

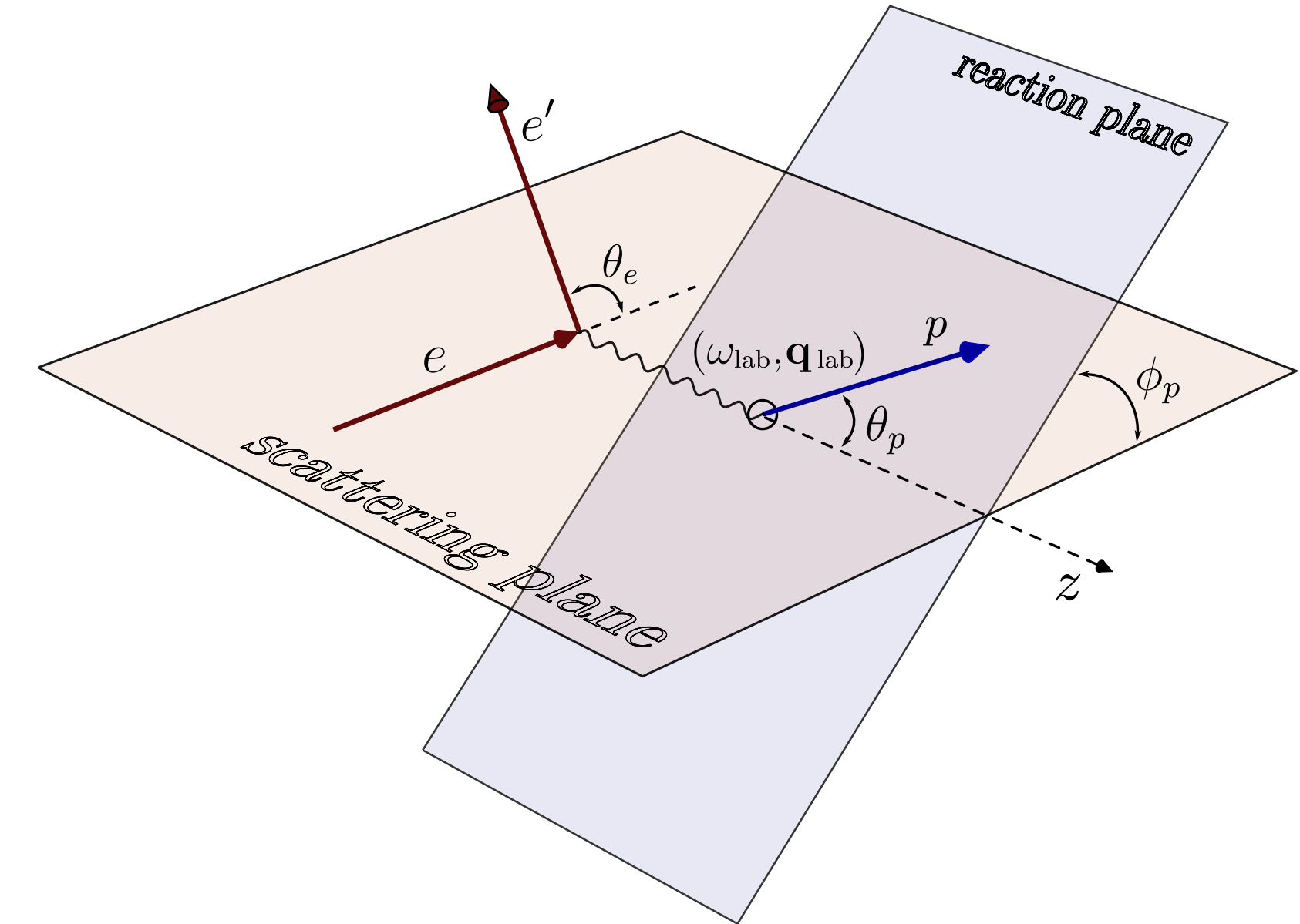
# Test ground: ${}^2\text{H}(e,e'p)n$

- Simplest knockout process (no induced 3N forces/currents)
- Focus on longitudinal structure function  $f_L$

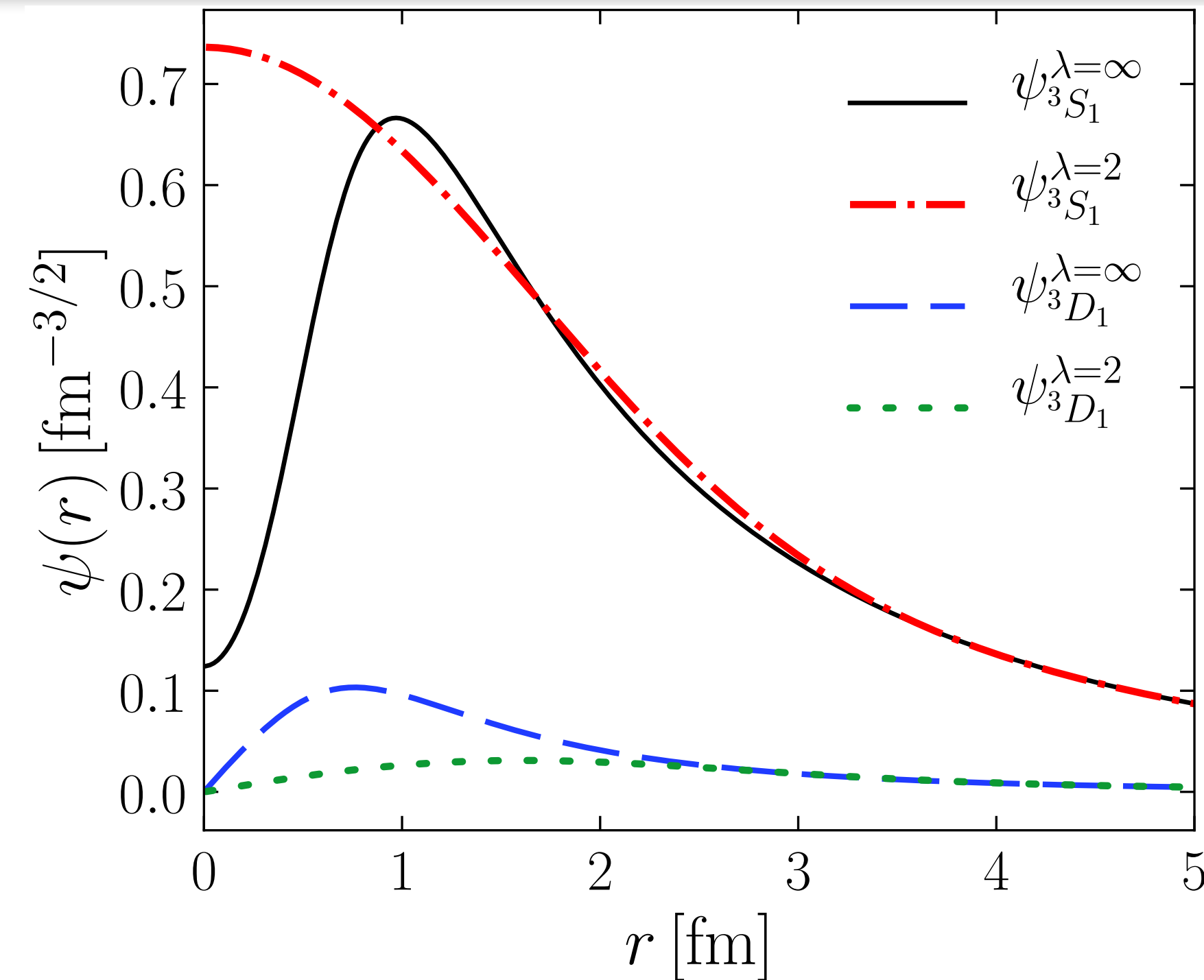
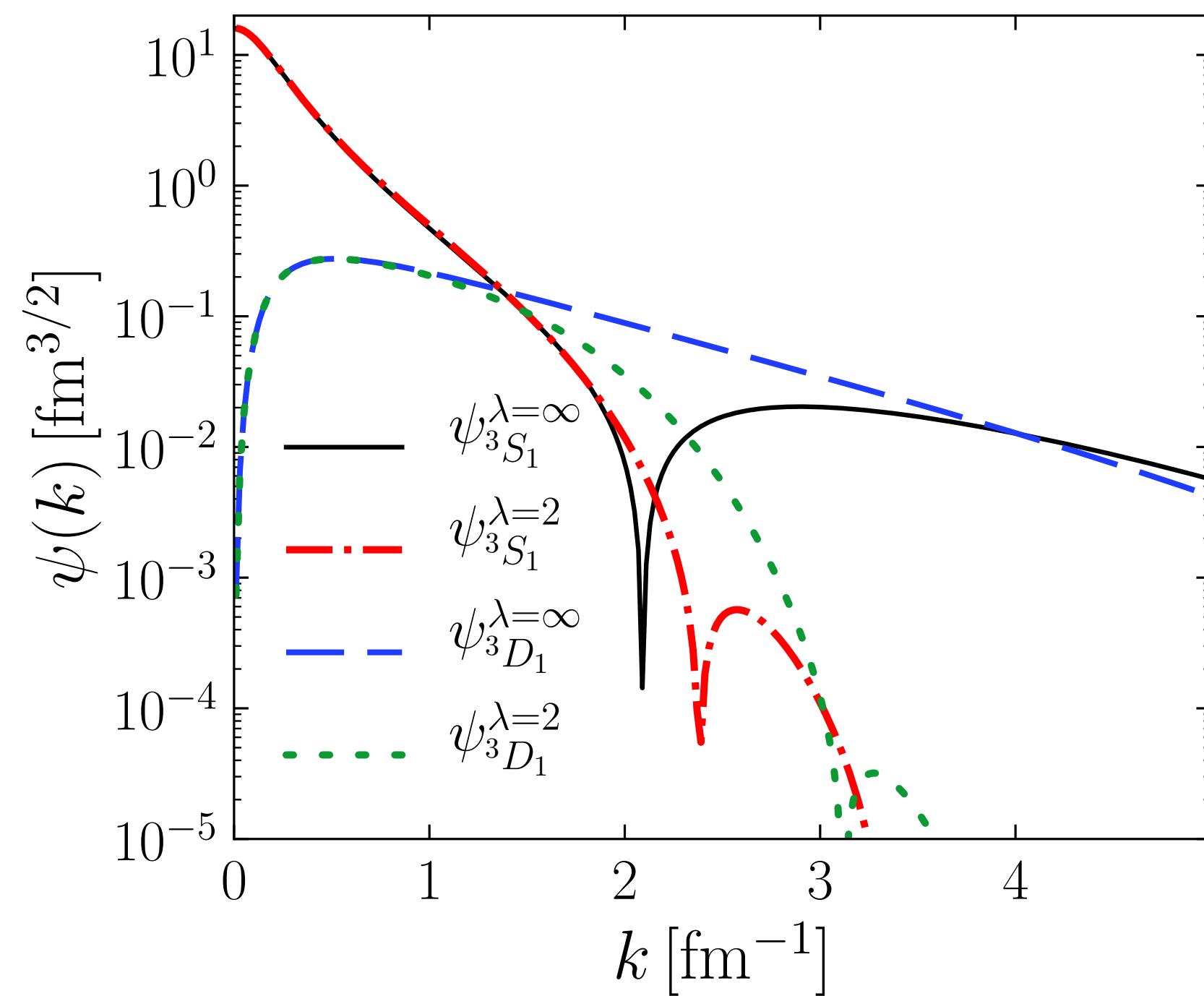
$$f_L \sim \sum_{m_s, m_J} |\langle \psi_f | J_0 | \psi_i \rangle|^2$$

$$\bullet \quad f_L^\lambda \sim \left| \underbrace{\langle \psi_f |}_{\psi_f^\lambda} U_\lambda^\dagger \underbrace{U_\lambda J_0 U_\lambda^\dagger}_{J_0^\lambda} \underbrace{U_\lambda | \psi_i \rangle}_{\psi_i^\lambda} \right|^2; \quad U_\lambda^\dagger U_\lambda = I; \quad f_L^\lambda = f_L$$

- Components (deuteron wf, transition operator, FSI) scale-dependent, total is not.
- Are some resolutions “better” than others? E.g., in a given kinematics, can FSI be minimized with different choices of  $\lambda$ ? How do interpretations change with scale?



# Deuteron wave function evolution

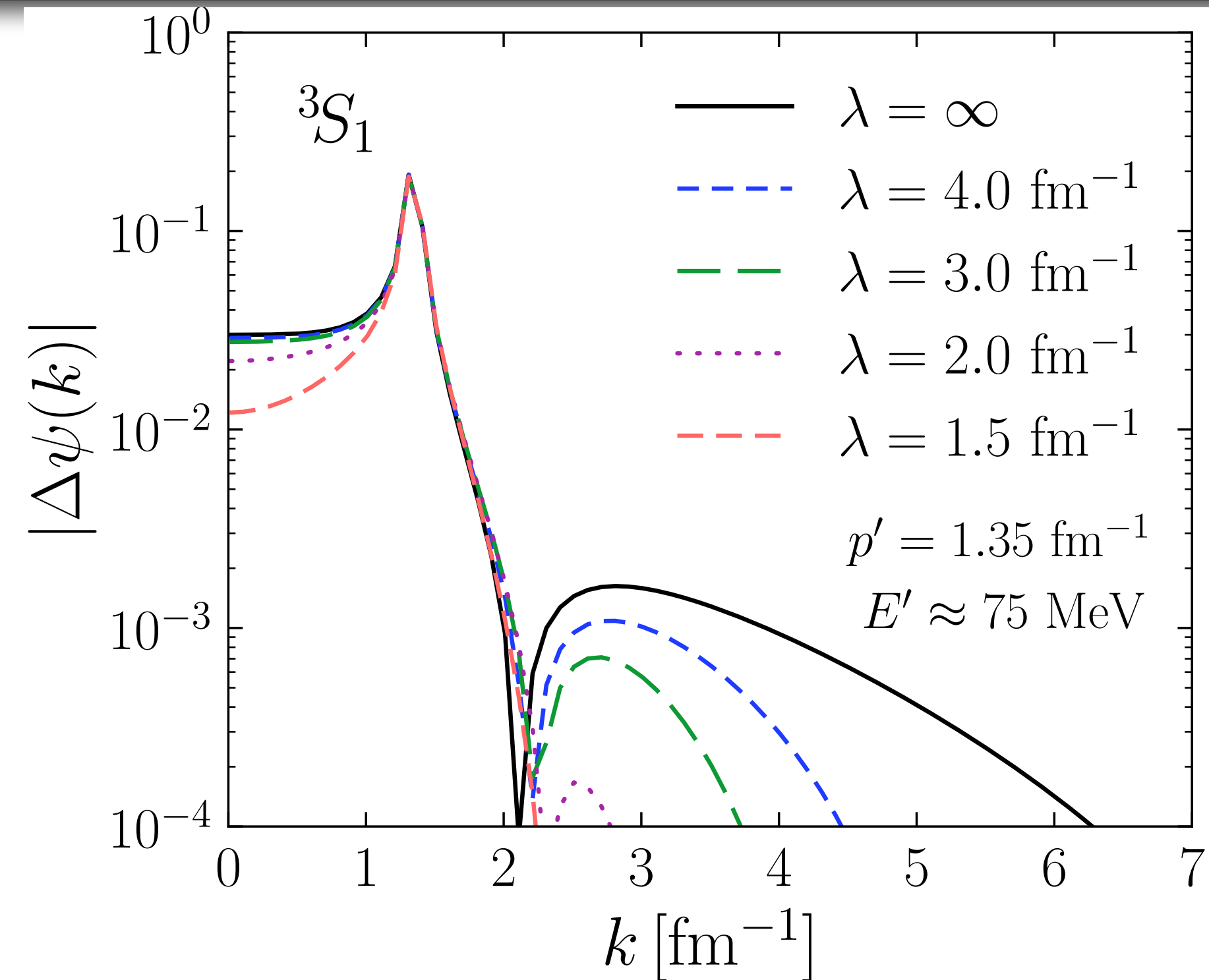
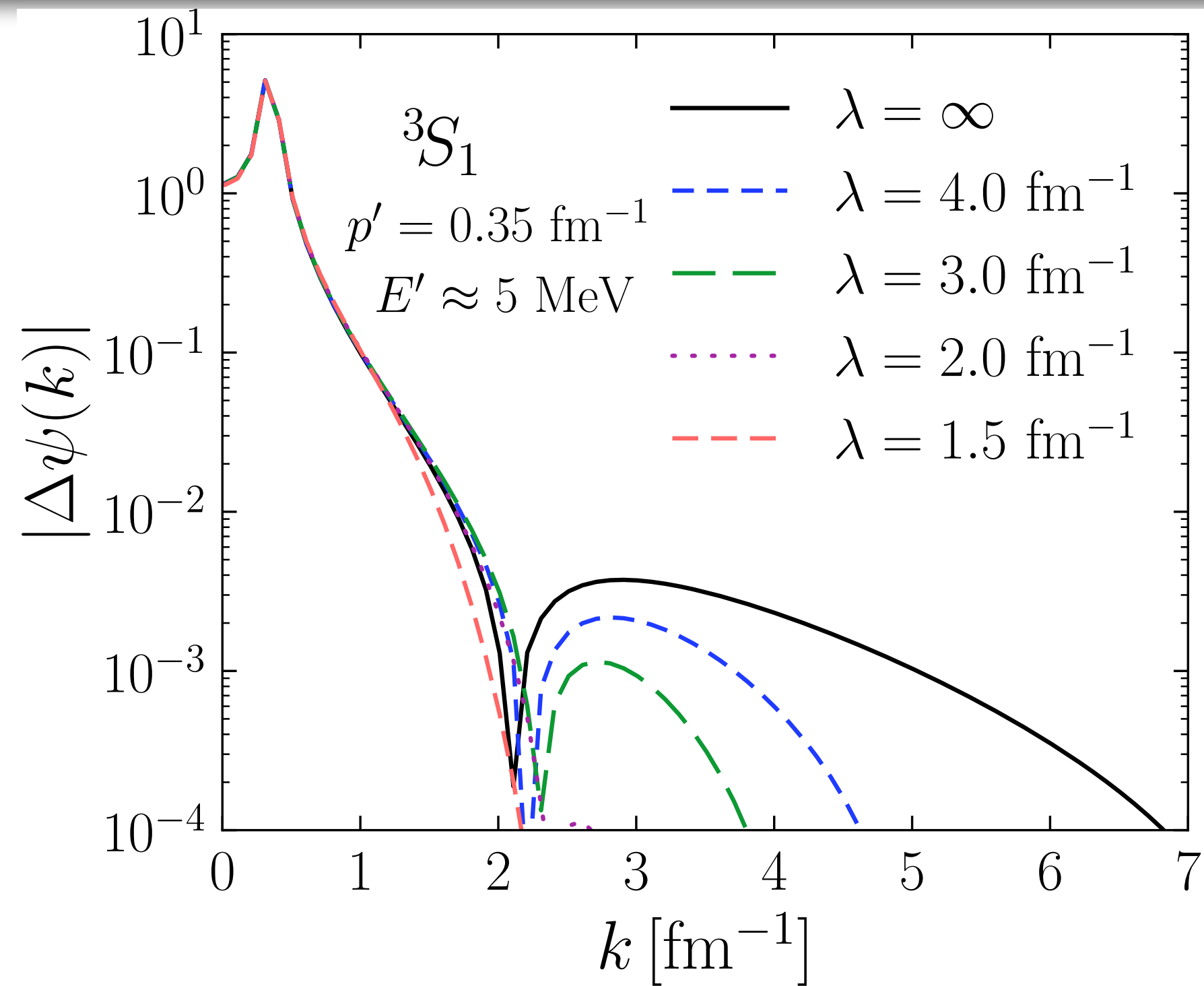


$k < \lambda$  components invariant  $\Leftrightarrow$  RG preserves long-distance physics

$k > \lambda$  components suppressed  $\Leftrightarrow$  short-range correlations blurred out

**Folklore:** Simple wave functions at low  $\lambda$   $\Leftrightarrow$  more complicated operators?  
especially for high- $q$  processes?

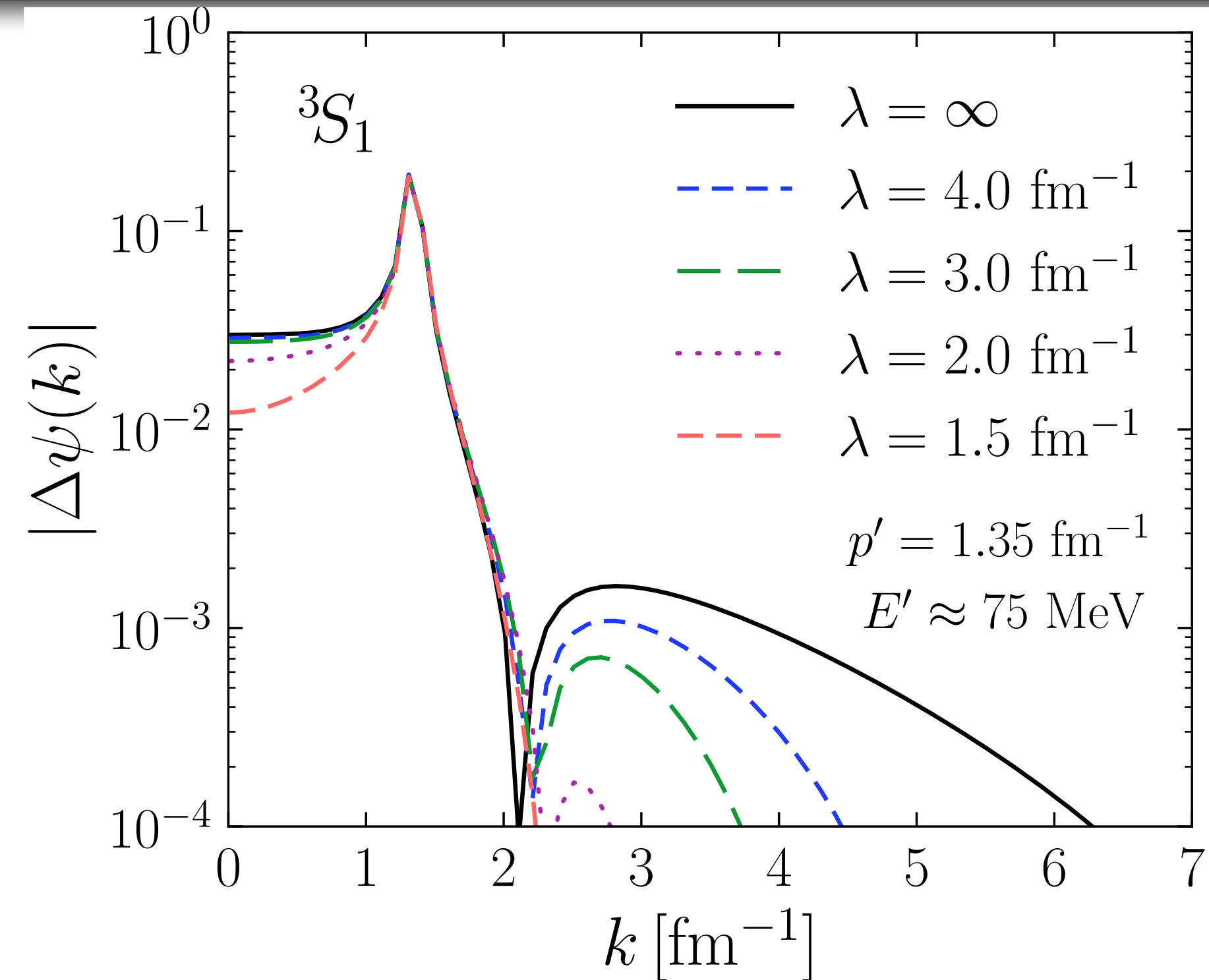
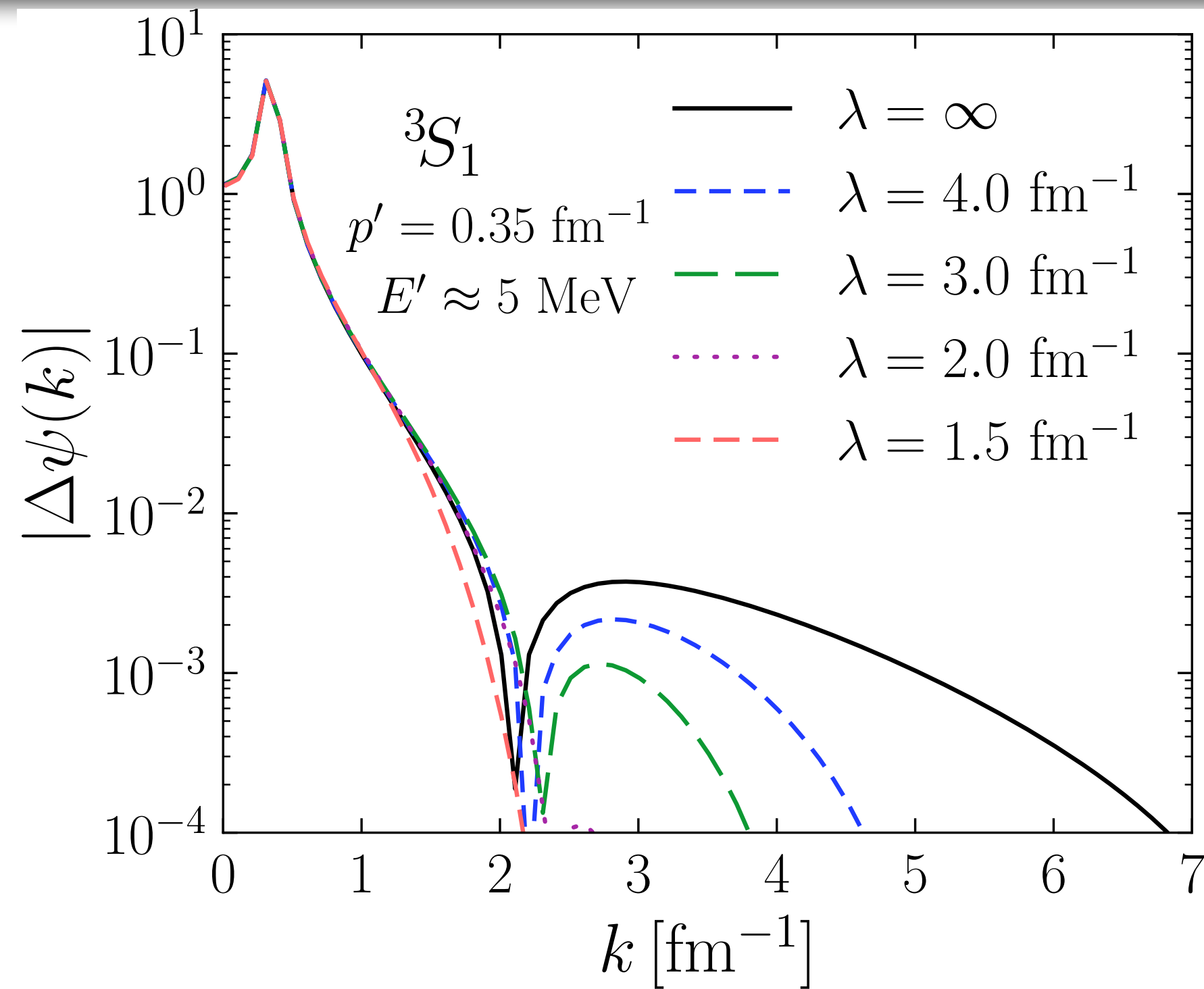
# Final-state wave function evolution



$$\psi_f^\lambda(p'; k) = \underbrace{\phi_{p'}}_{\text{IA}} + \underbrace{\Delta\psi_\lambda(p'; k)}_{\text{FSI}}$$



# Final-state wave function evolution



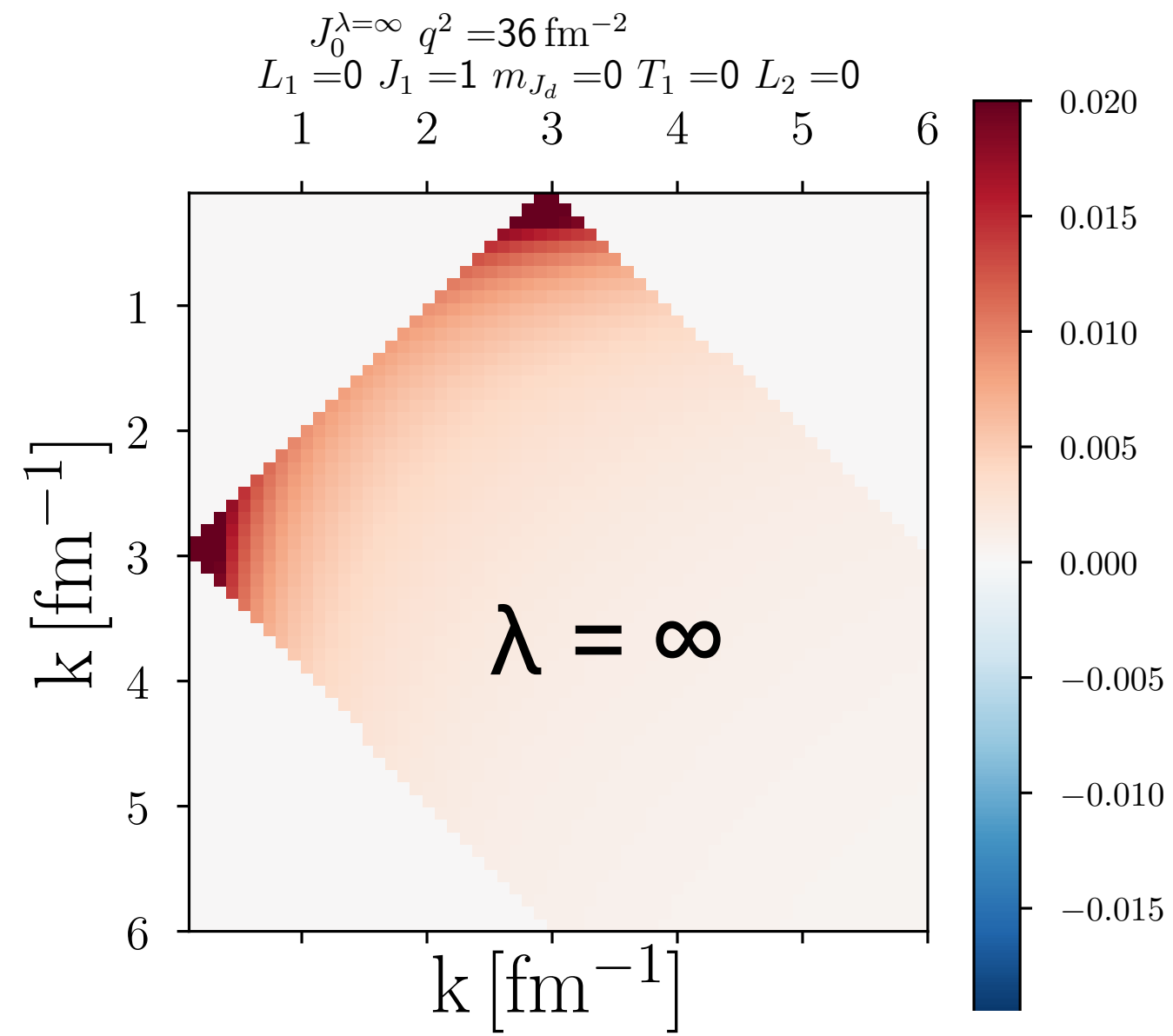
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- High-k tail suppressed with evolution

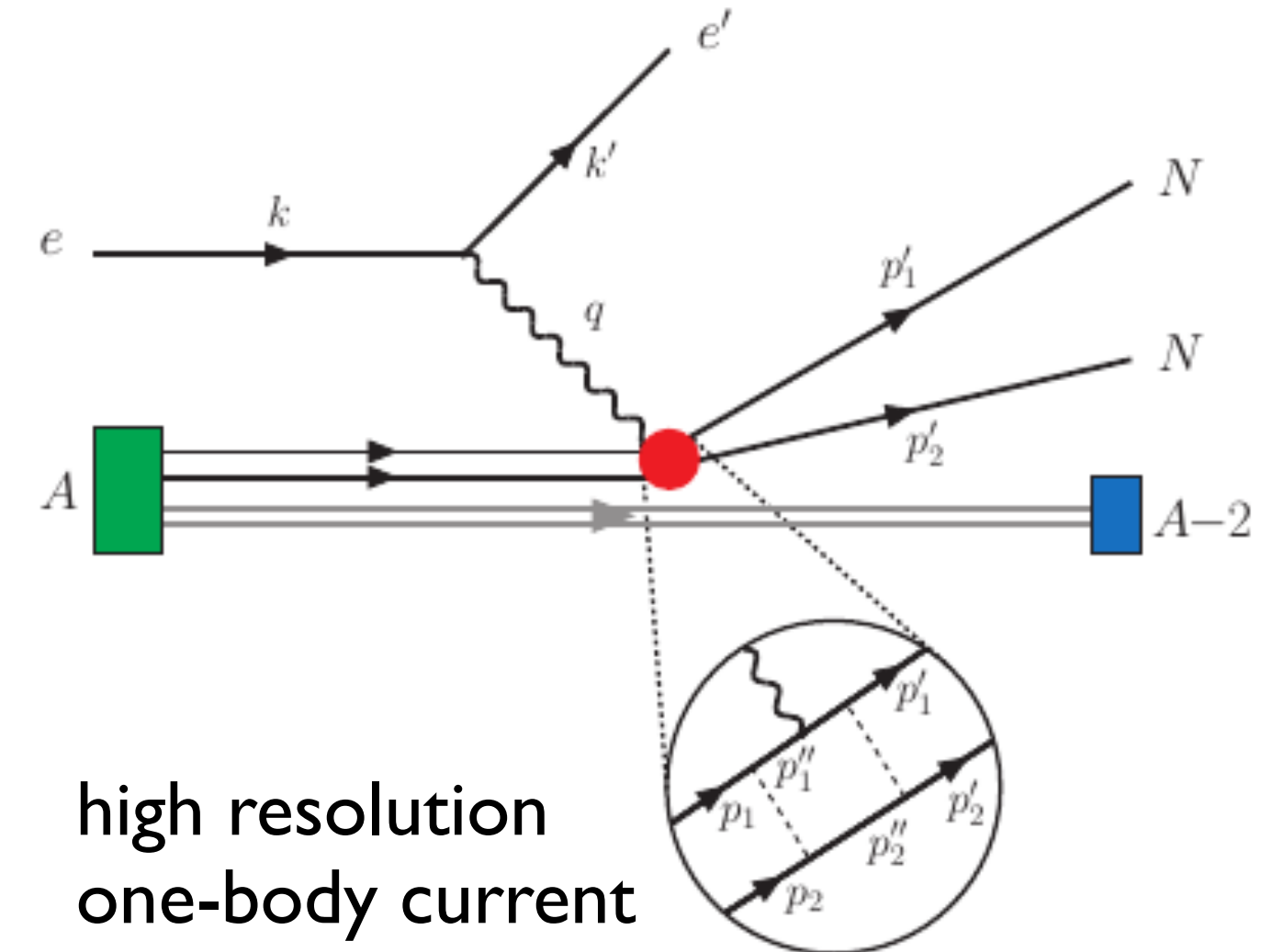
- For  $p' \gtrsim \lambda$ ,  $\Delta\psi_f^\lambda(p'; k)$  localized around outgoing  $p'$

“local decoupling” Dainton et al. PRC 89 (2014)

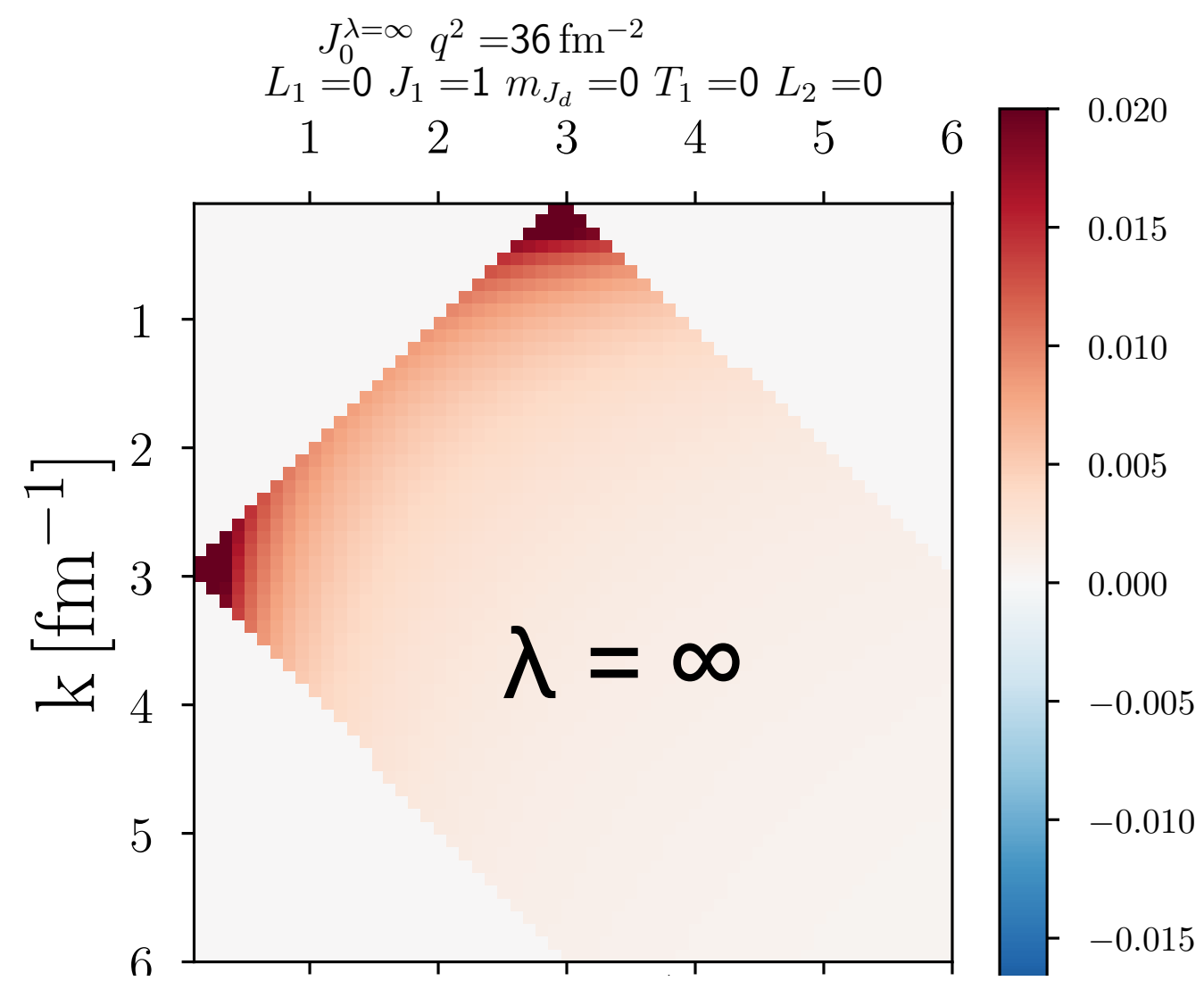
# Current operator evolution



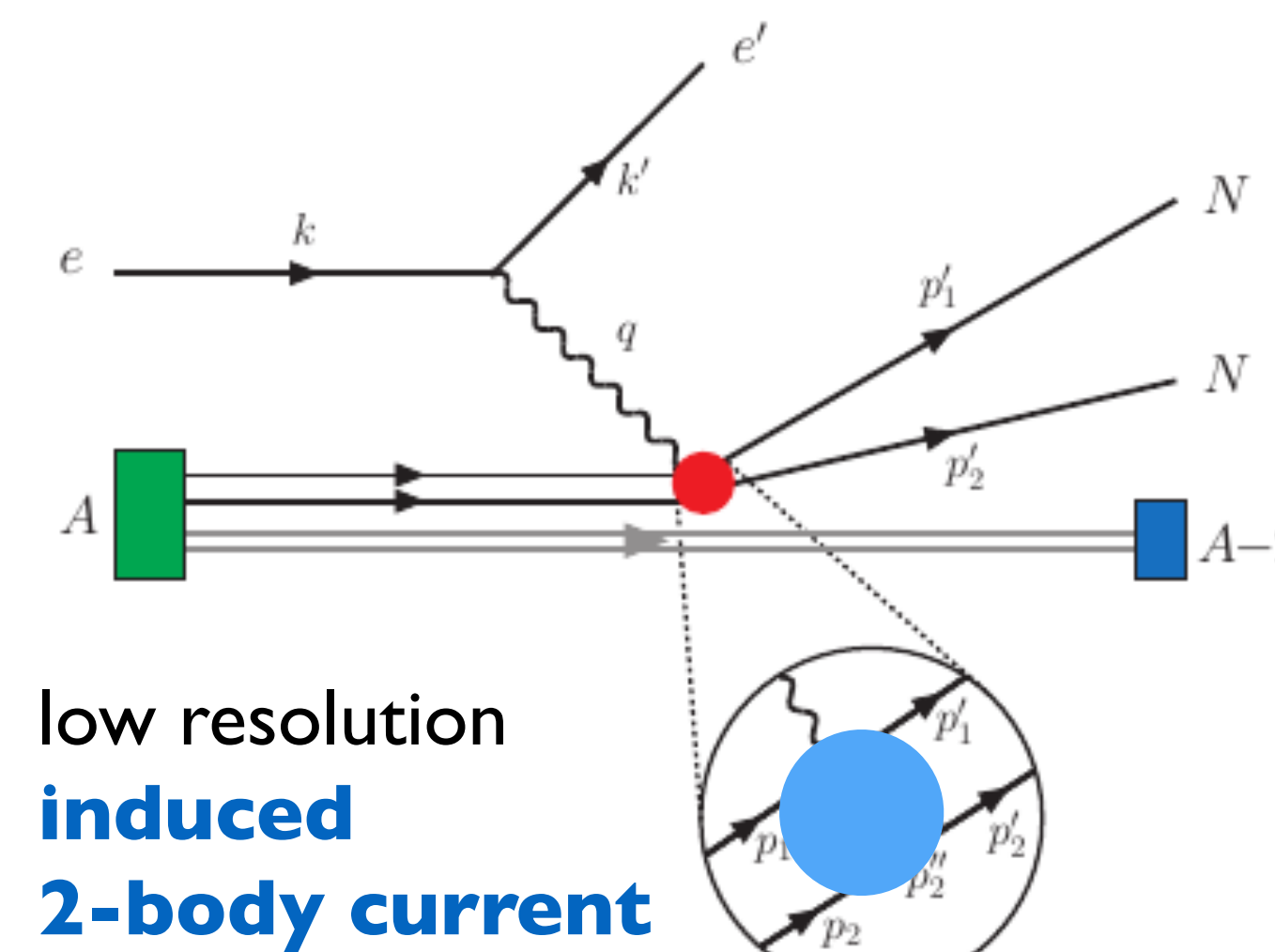
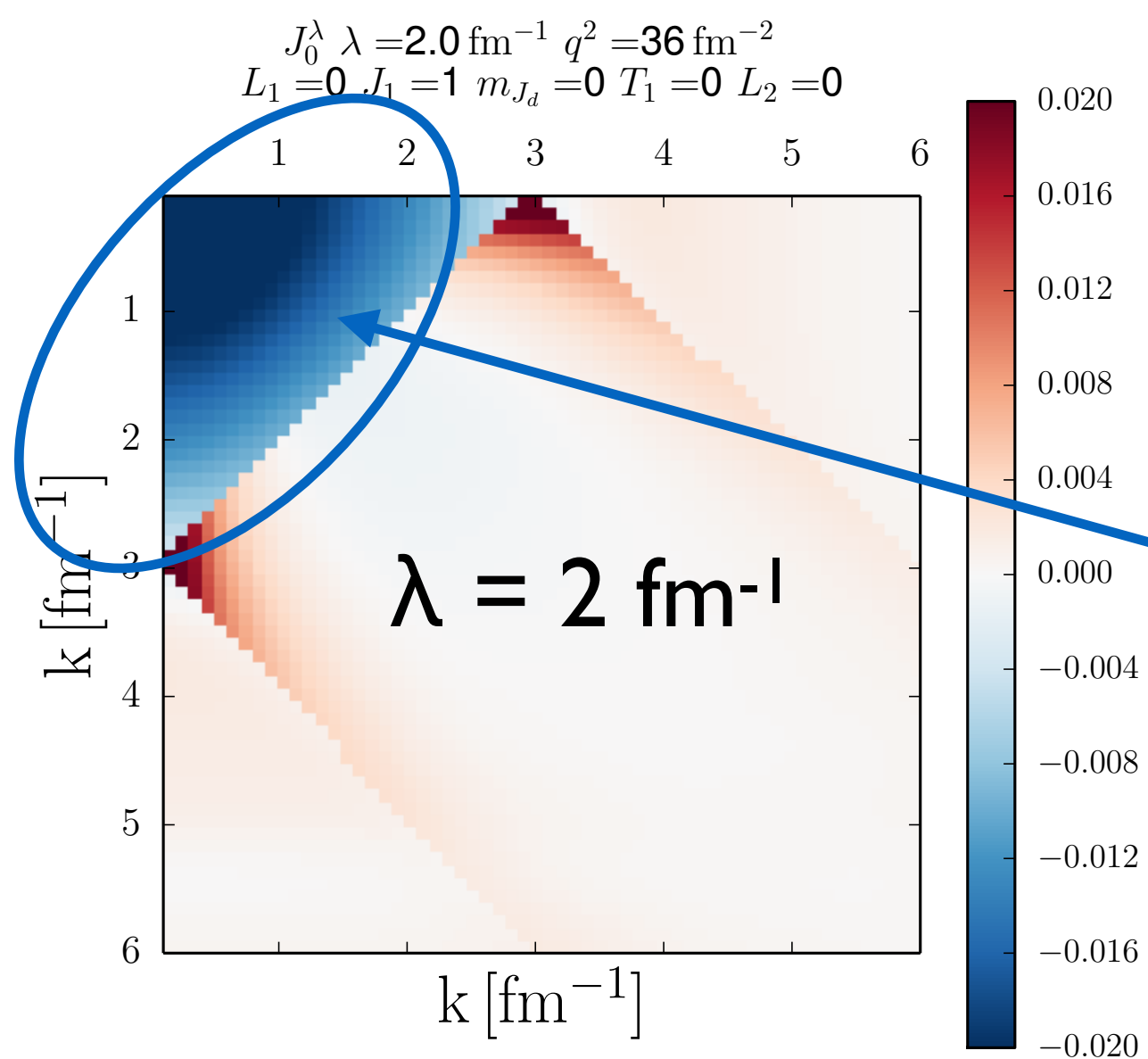
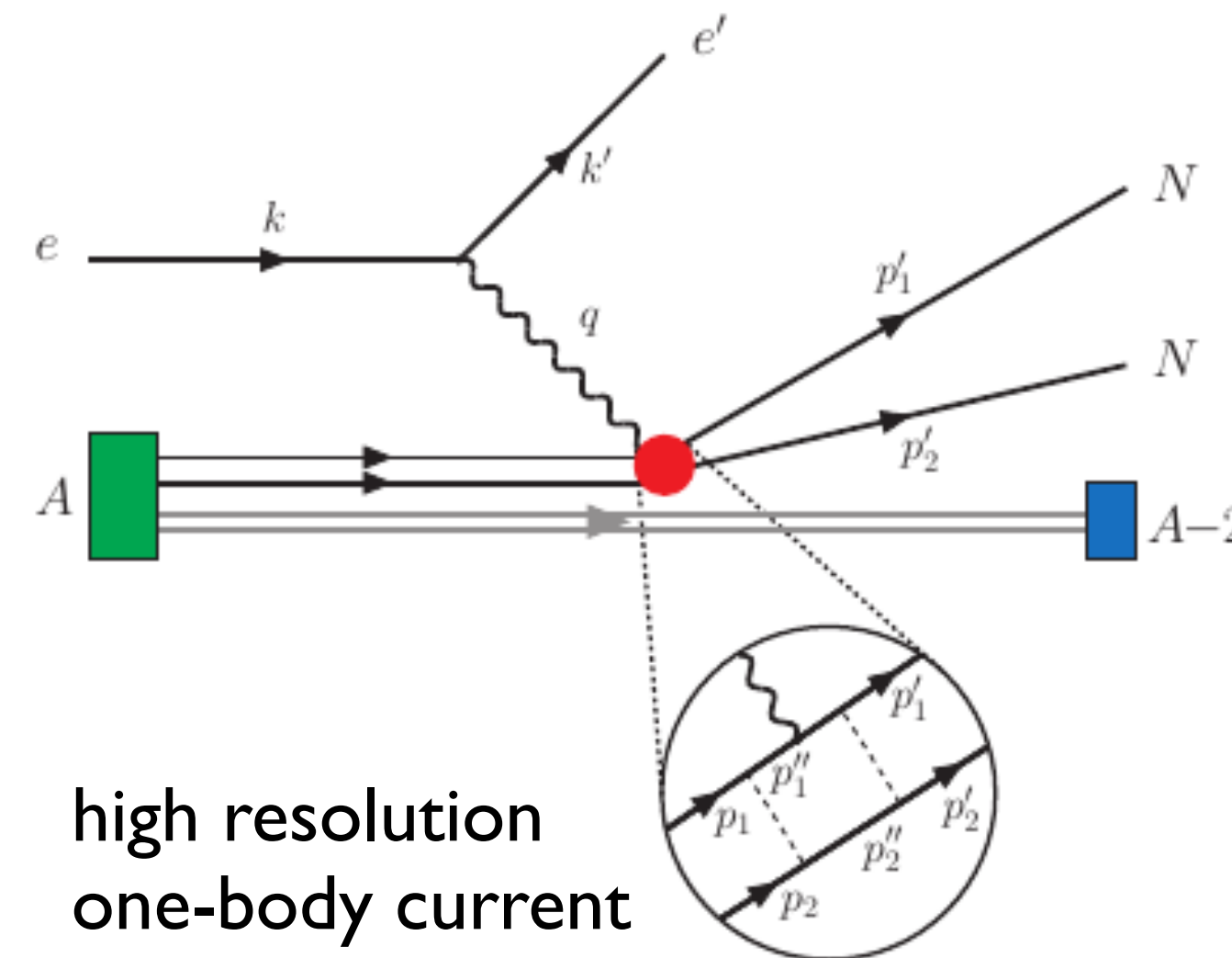
$^3S_1$  channel  
 $q^2 = 36 \text{ fm}^{-2}$



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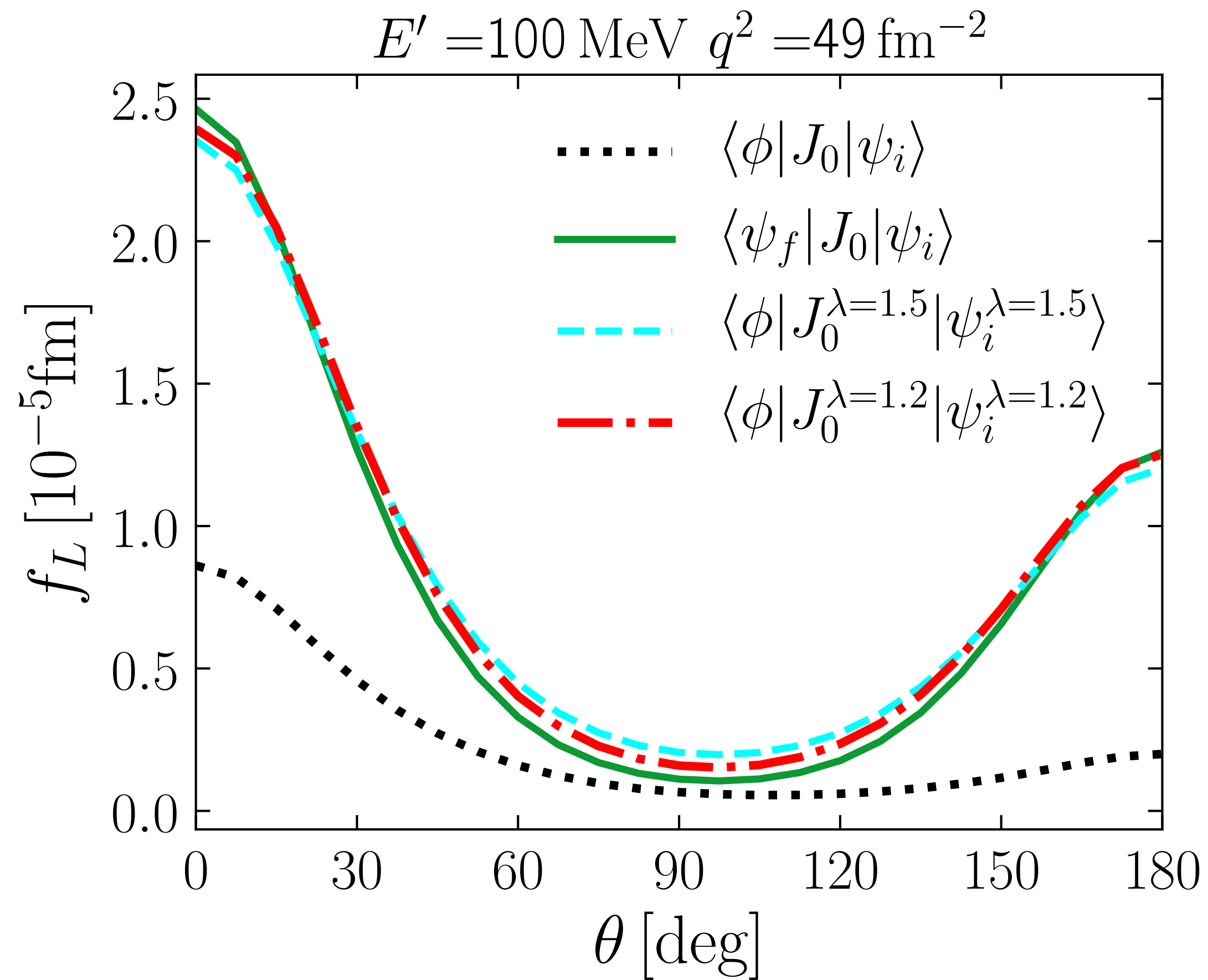


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# Scale Dependence of Final State Interactions

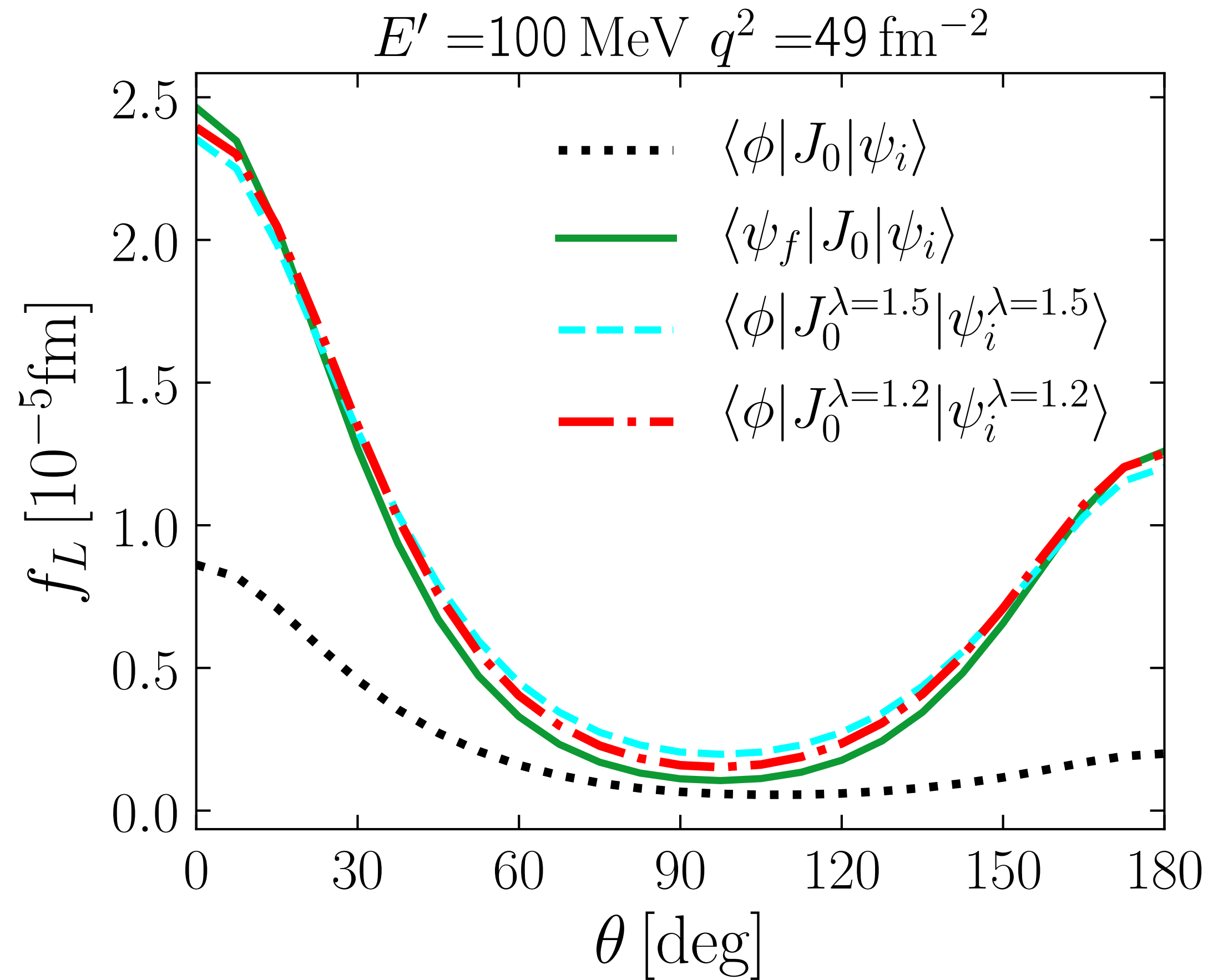
Look at kinematics relevant to SRC studies



$x_B = 1.64$ ,  $Q^2 = 1.78 \text{ GeV}^2$

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Look at kinematics relevant to SRC studies



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FSI sizable at large  $\lambda$   
but negligible at low-resolution!

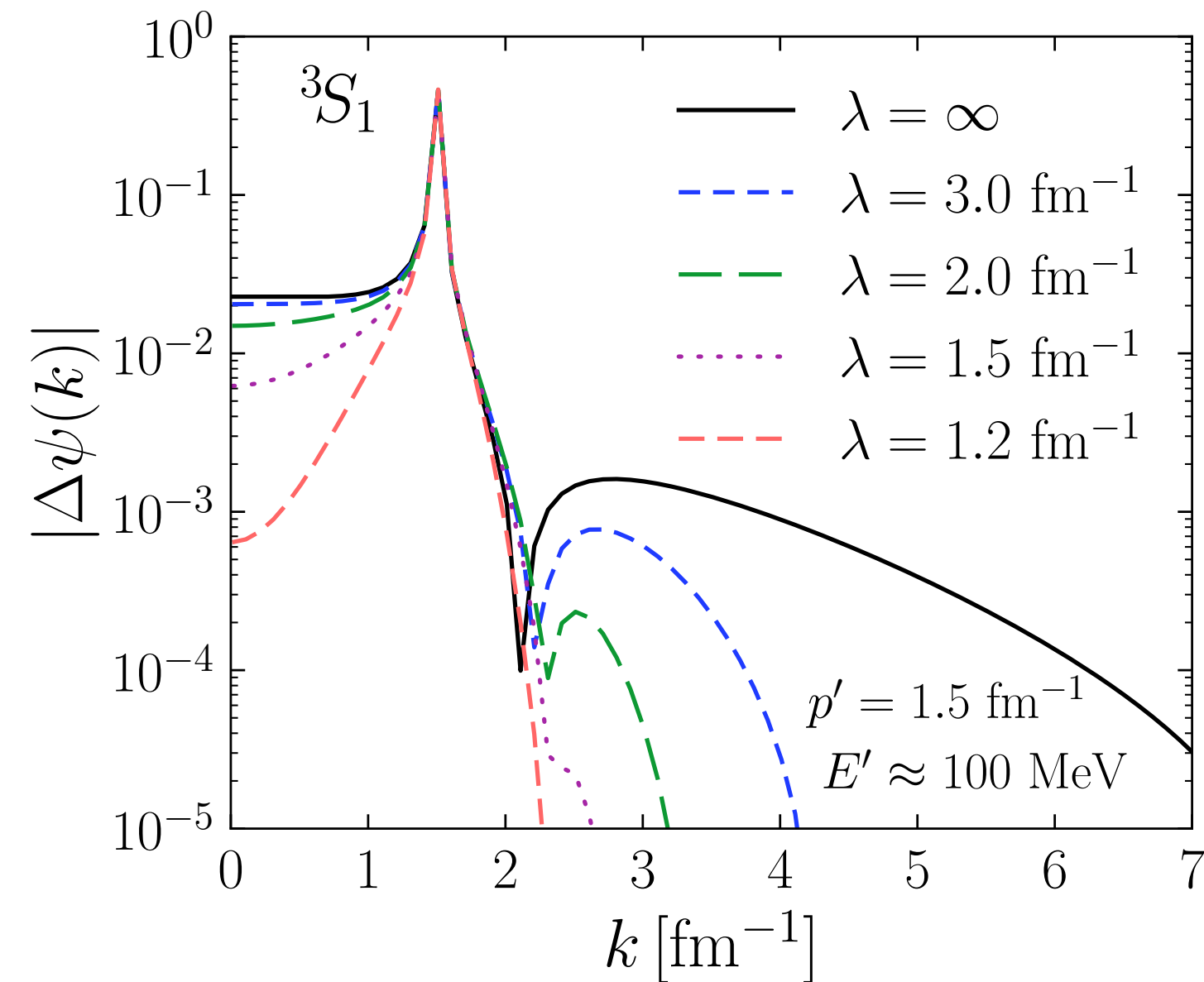
Folklore:

shouldn't hard processes  
be complicated in low resolution  
( $\lambda \ll q$ ) pictures?

**Why are FSI so small at low  $\lambda$   
in these kinematics ?**

# Scale Dependence of Final State Interactions

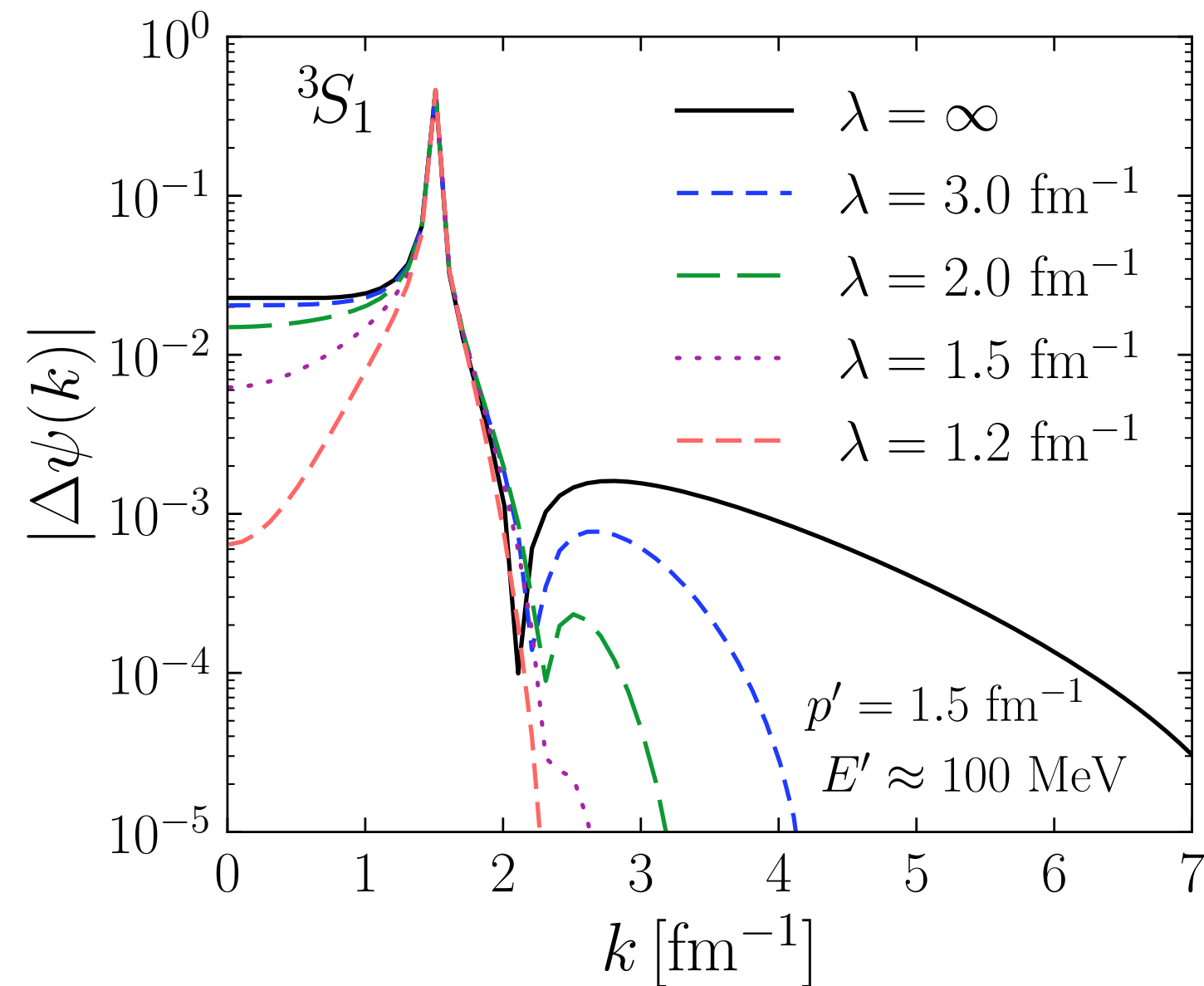
final state wf (FSI piece)



For  $p' \gtrsim \lambda$ , interacting part of final state wf localized at  $k \approx p'$

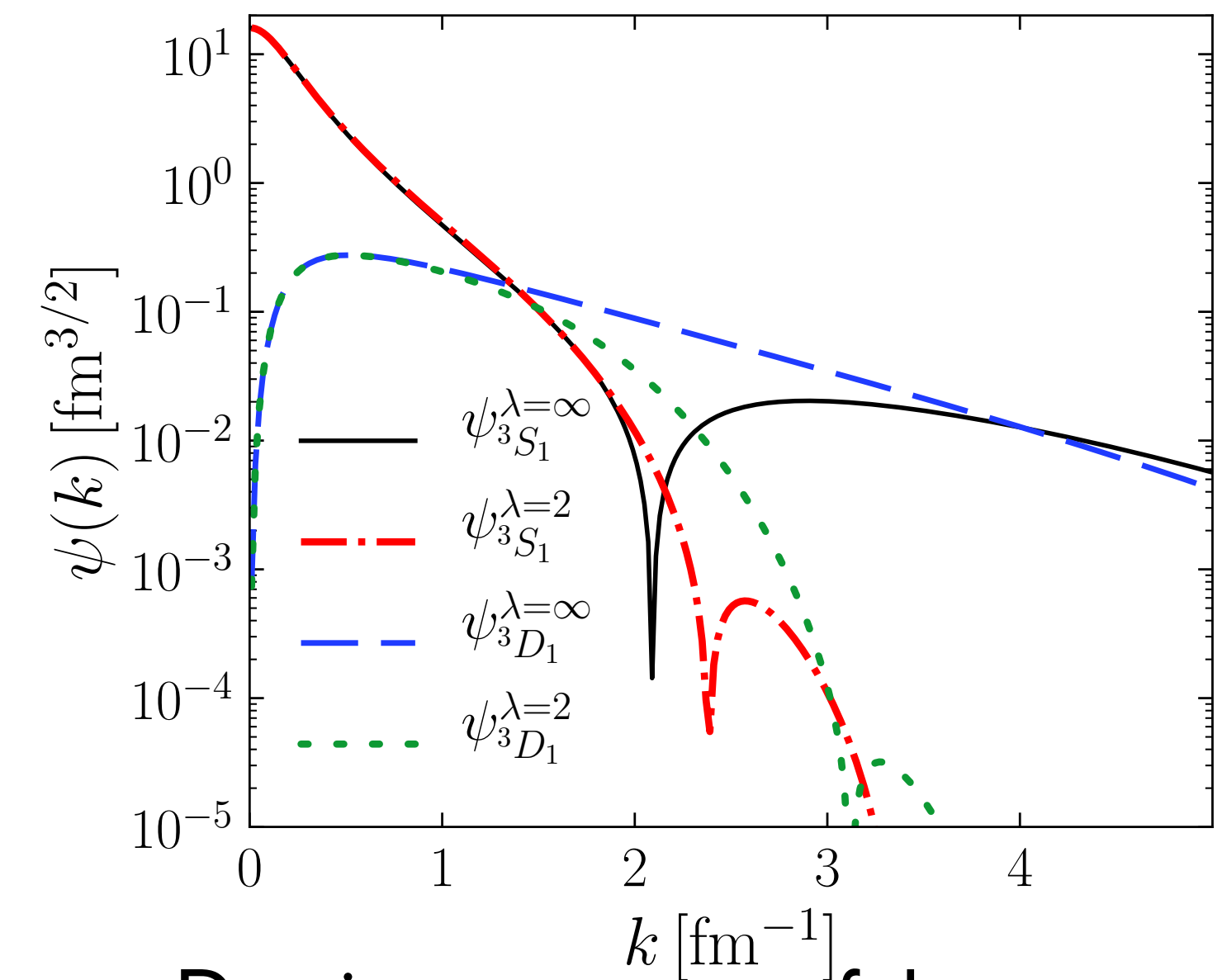
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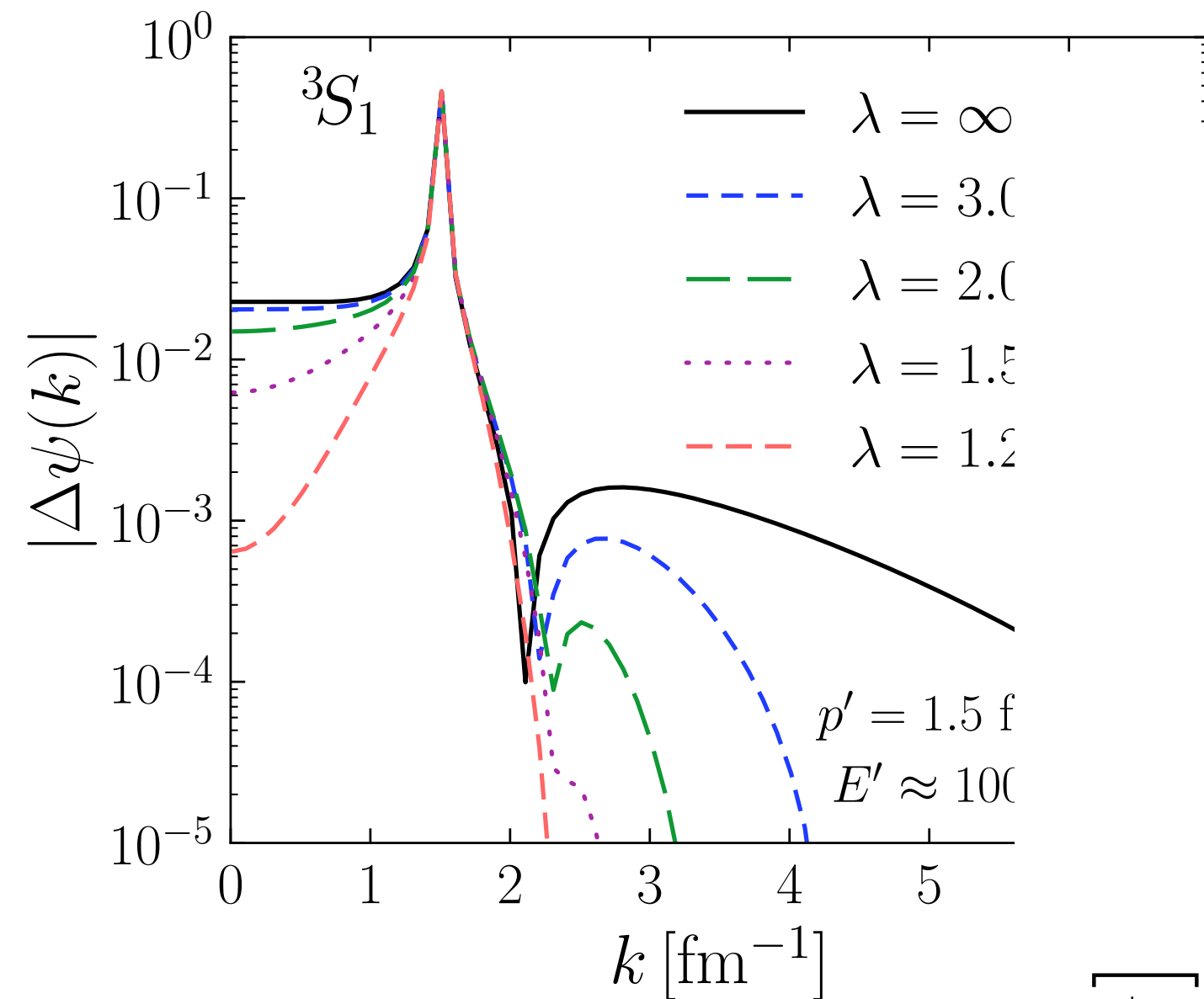
initial state (deuteron) wf



Dominant support of deuteron wf at  $k \lesssim \lambda$

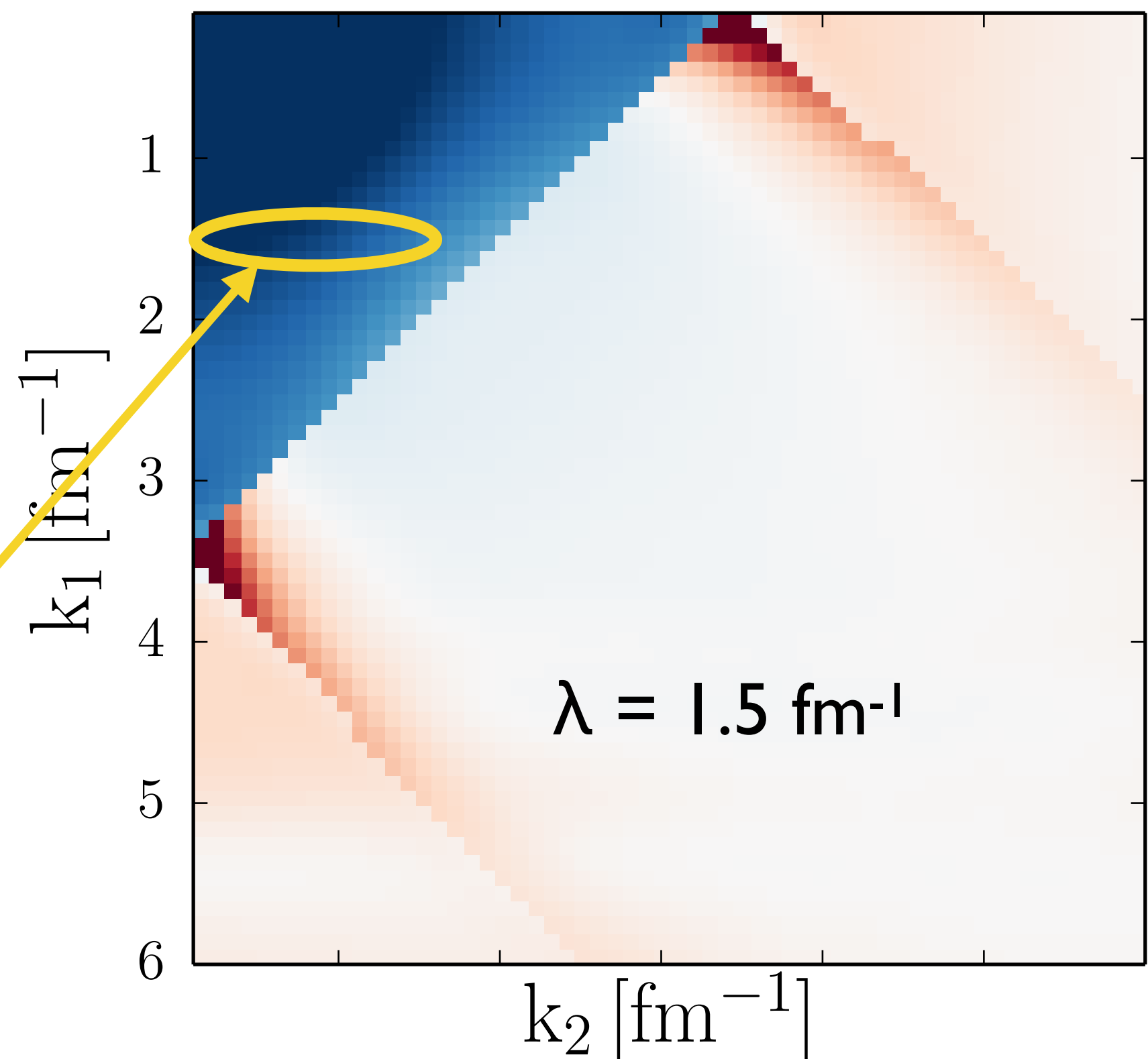
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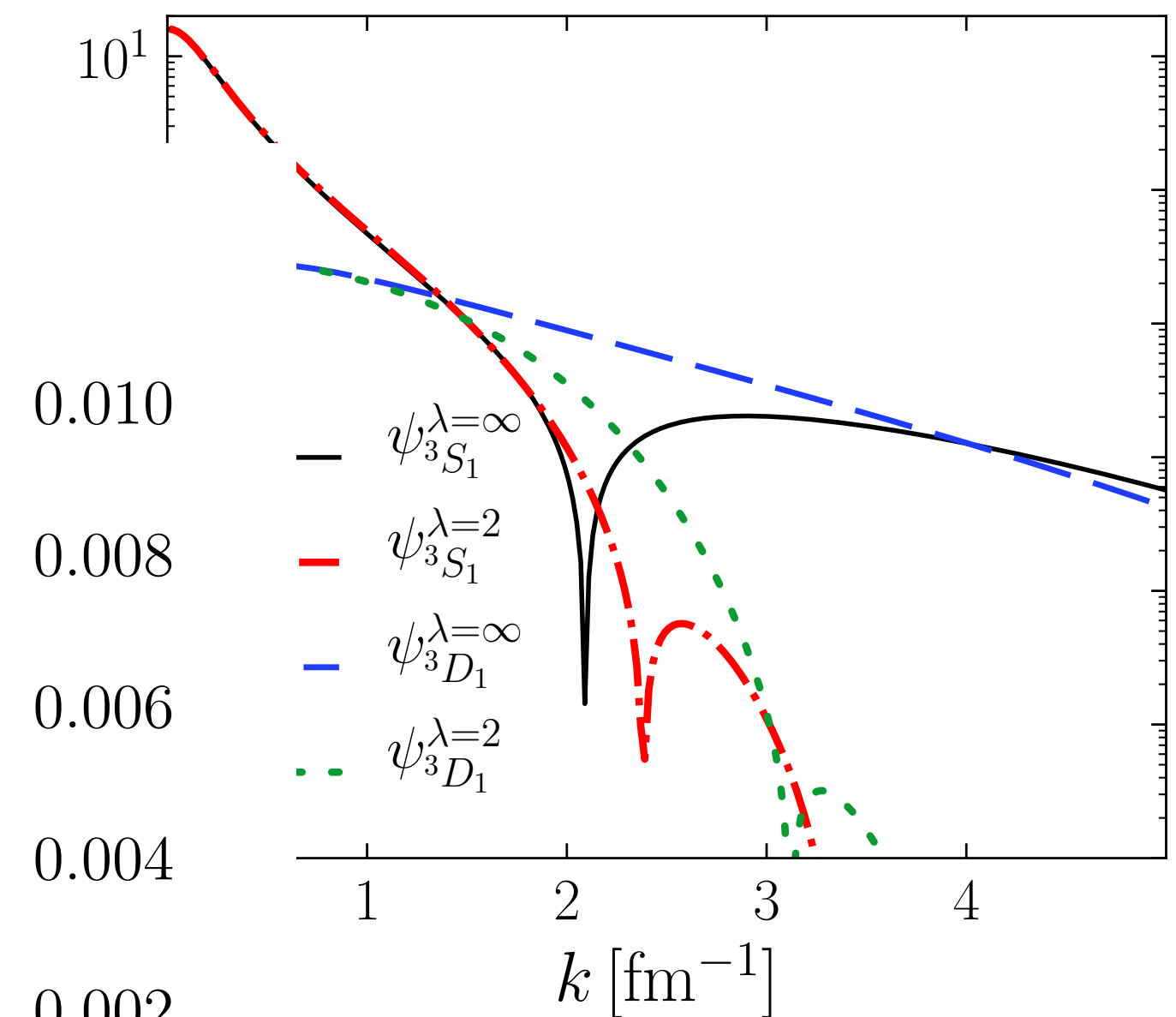
$J_q^\lambda(k', k)$   
probed by  
transition  
(smooth and  
non-singular)

$$\langle {}^3S_1; k_1 | J_0^{\lambda=1.5} | {}^3S_1; k_2 \rangle \quad q^2 = 49 \text{ fm}^{-2}$$



$\lambda = 1.5 \text{ fm}^{-1}$

initial state (deuteron) wf

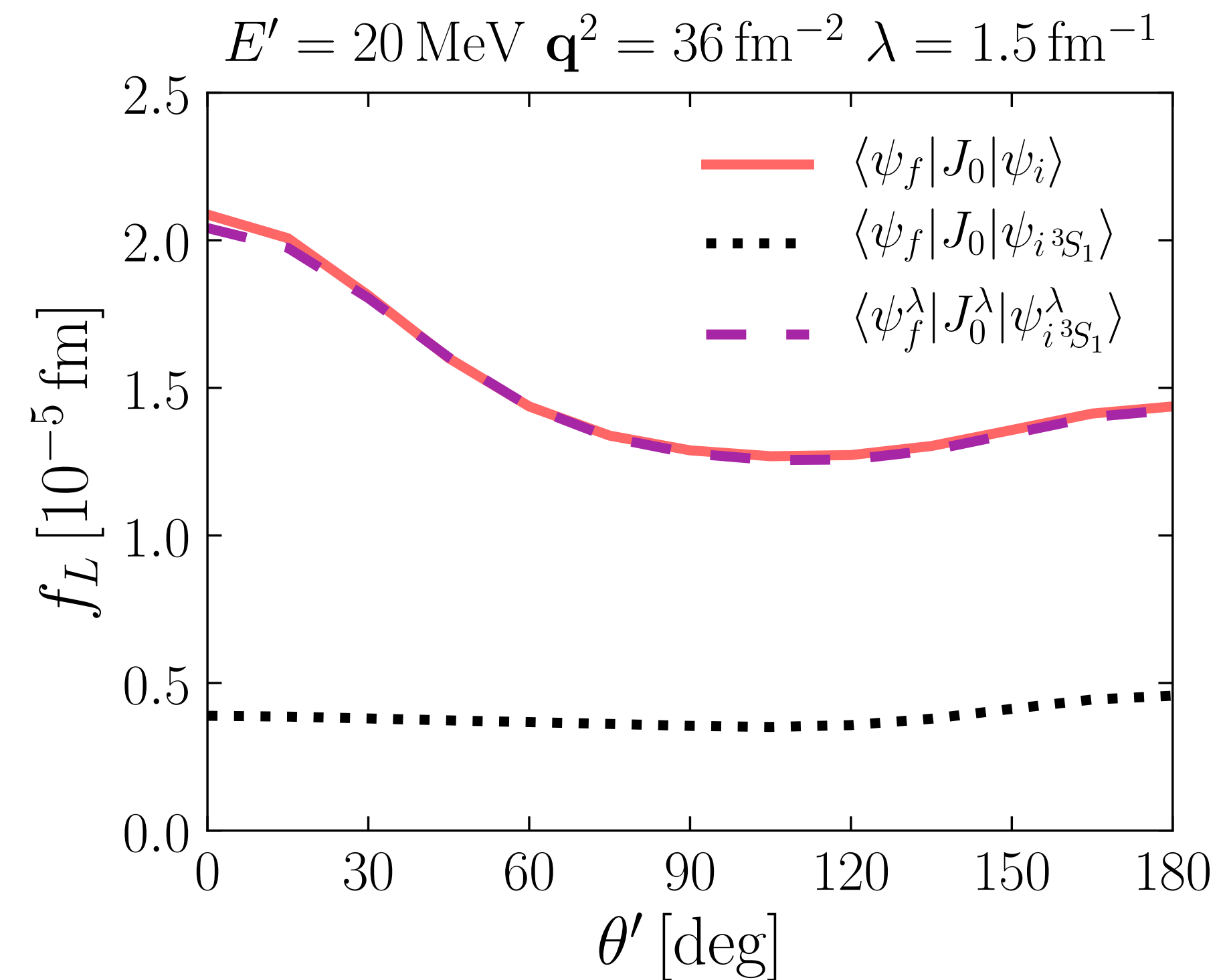
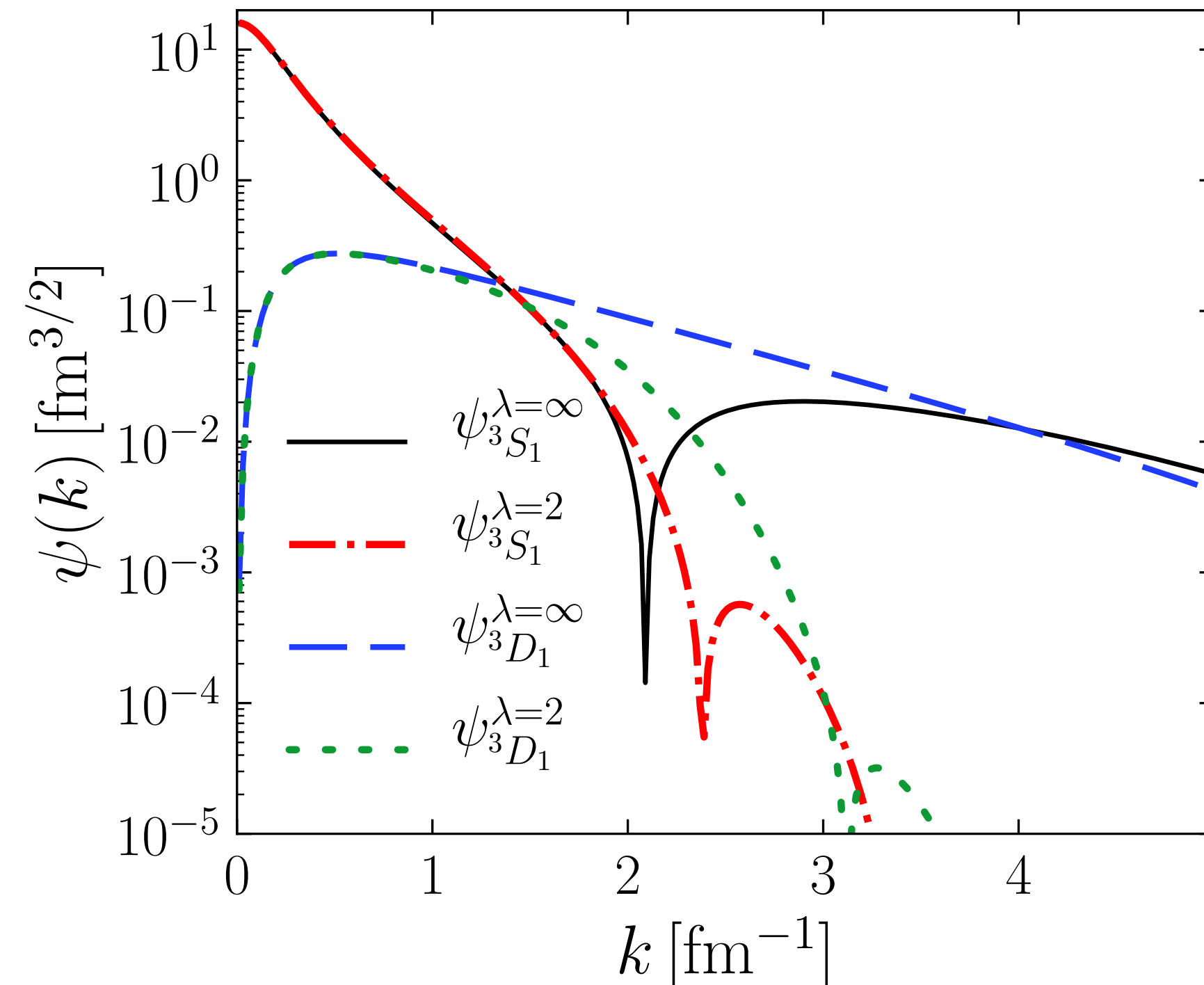


∴ FSI  $\sim T(p', p')$   
(small!)



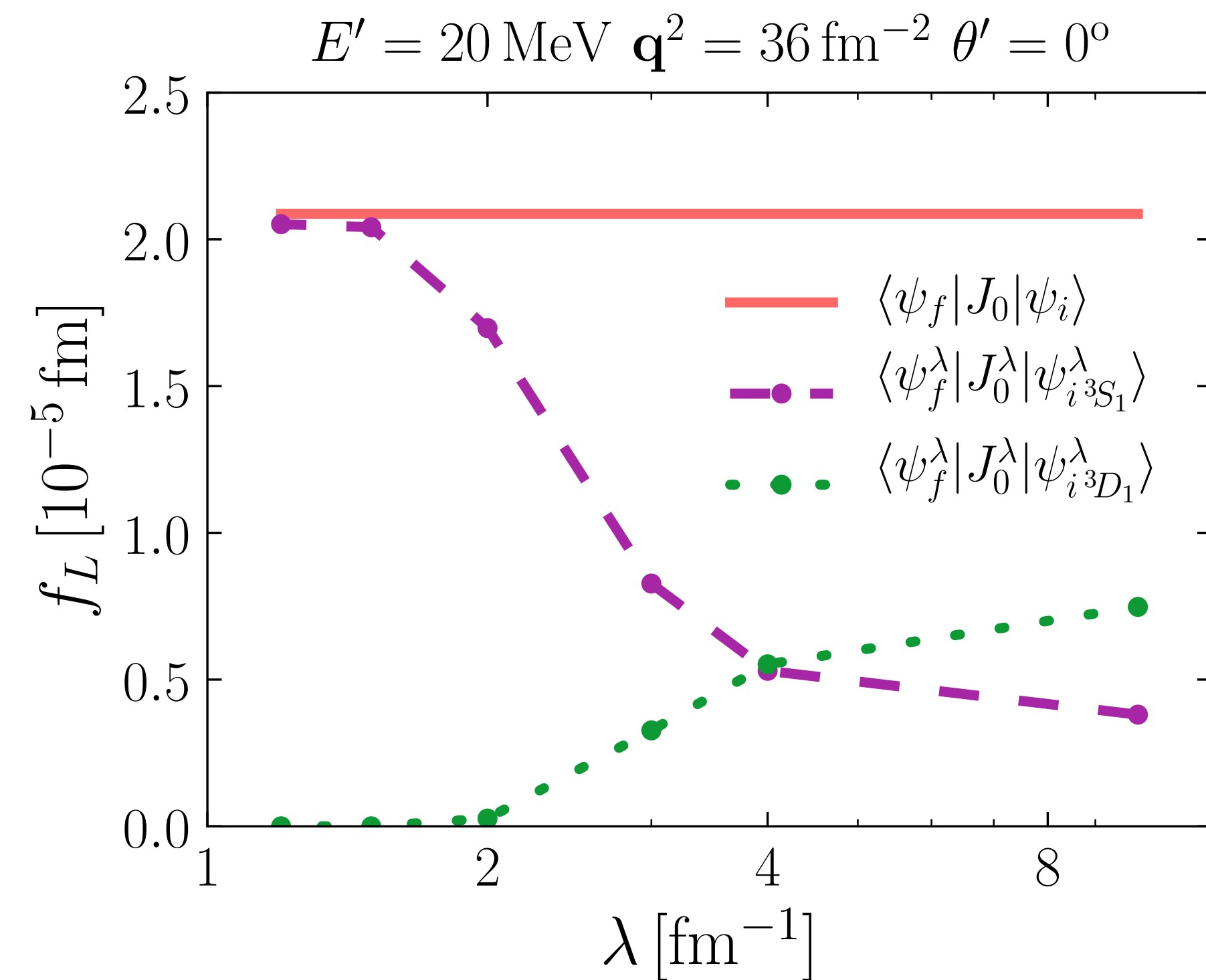
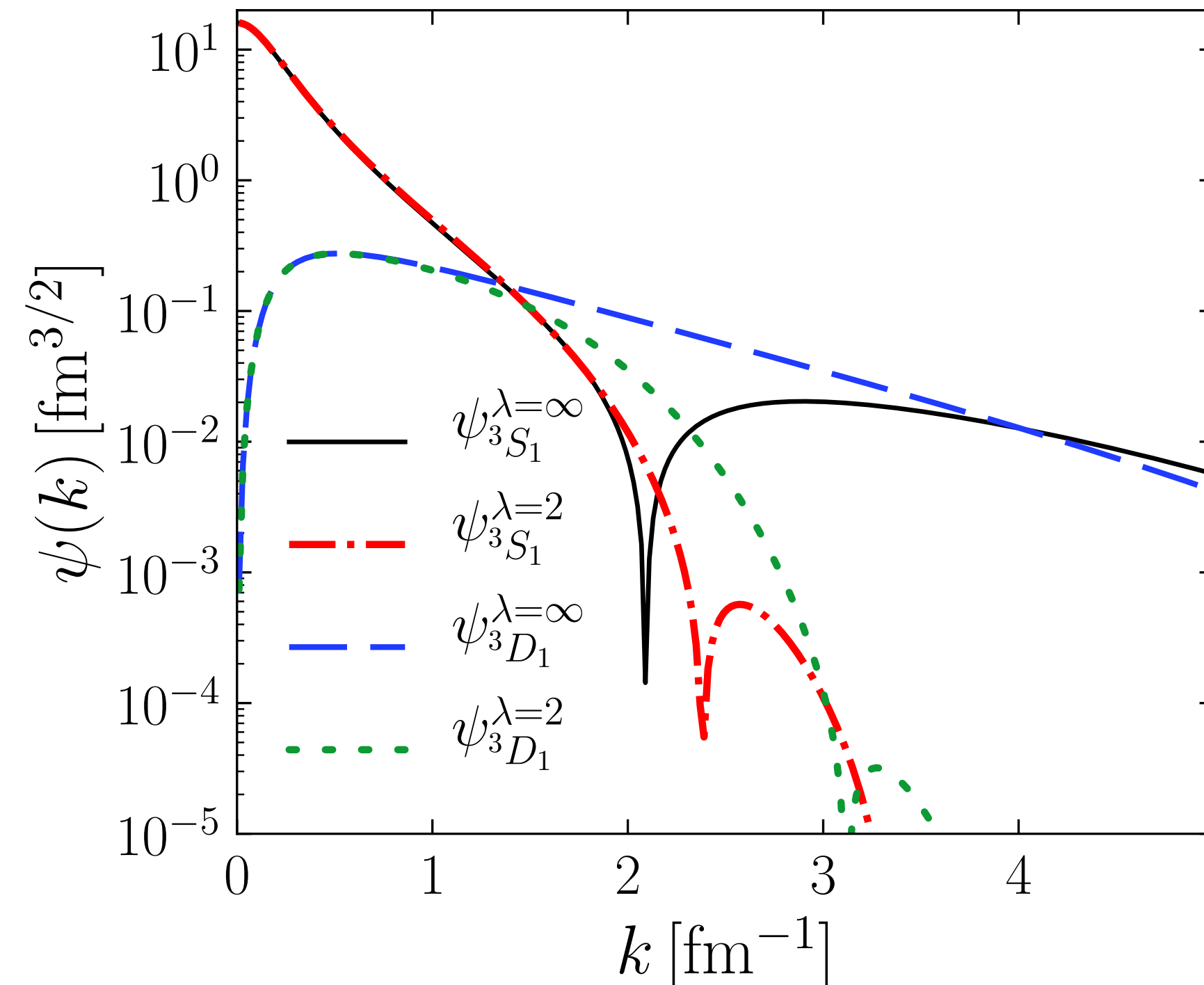
# Scale Dependence of Interpretations

- Analysis/interpretation of a reaction involves understanding which part of wave functions probed (**highly scale dependent!**)
- E.g., sensitivity to D-state w.f. in large  $\mathbf{q}^2$  processes



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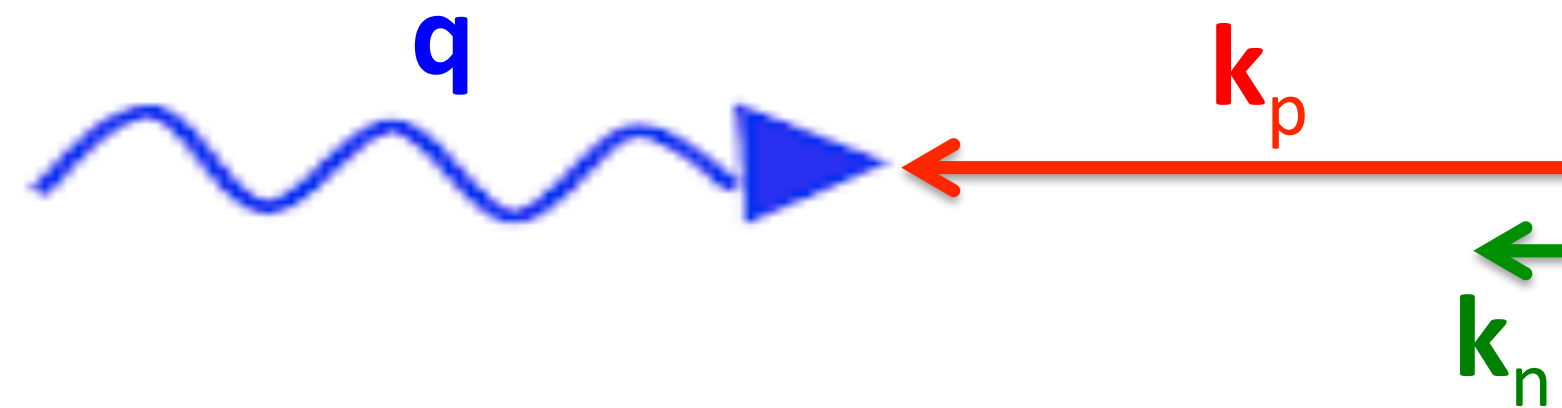
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# Scale Dependence of SRC Interpretation

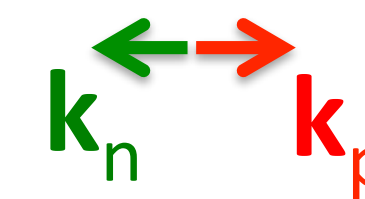
- Consider large  $q^2$  near threshold (small  $p'$ ) for  $\theta=0$  in **high-resolution** picture (COM frame of outgoing np)

Before



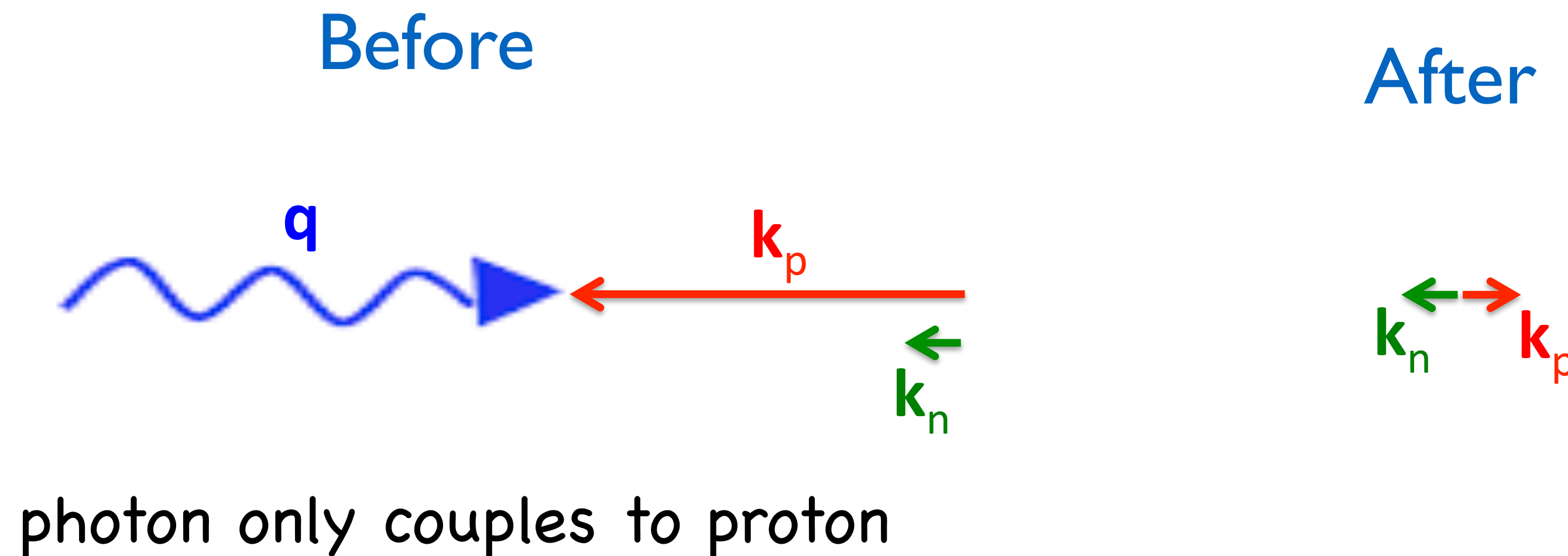
photon only couples to proton

After



# Scale Dependence of SRC Interpretation

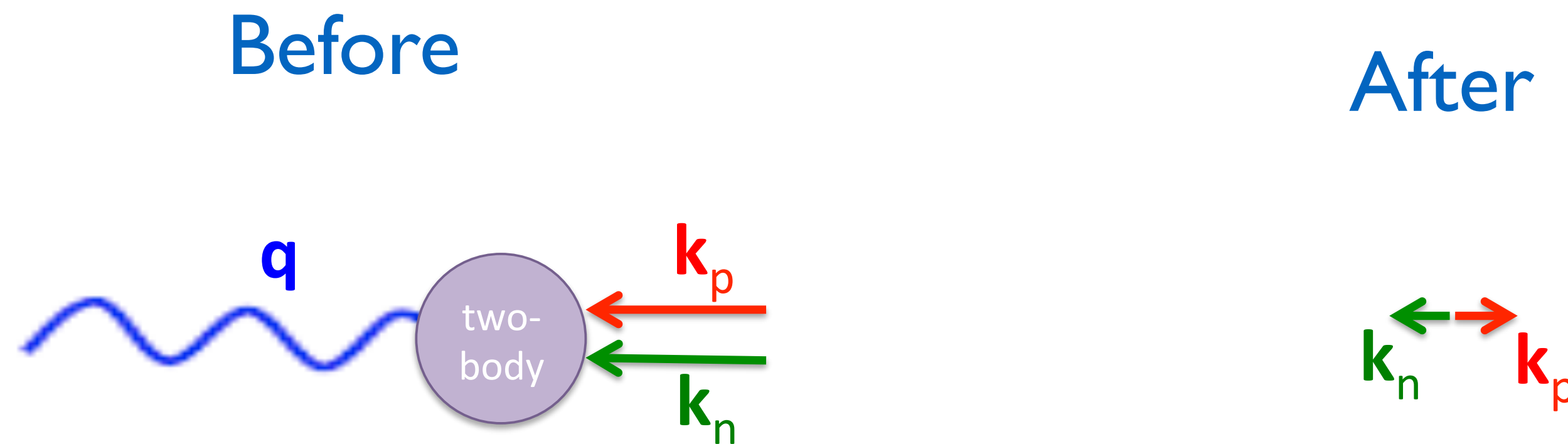
- Consider large  $q^2$  near threshold (small  $p'$ ) for  $\theta=0$  in **high-resolution** picture (COM frame of outgoing np)



- proton has large momentum  $\Rightarrow$  initial large relative momentum  
(i.e., SRC pair)

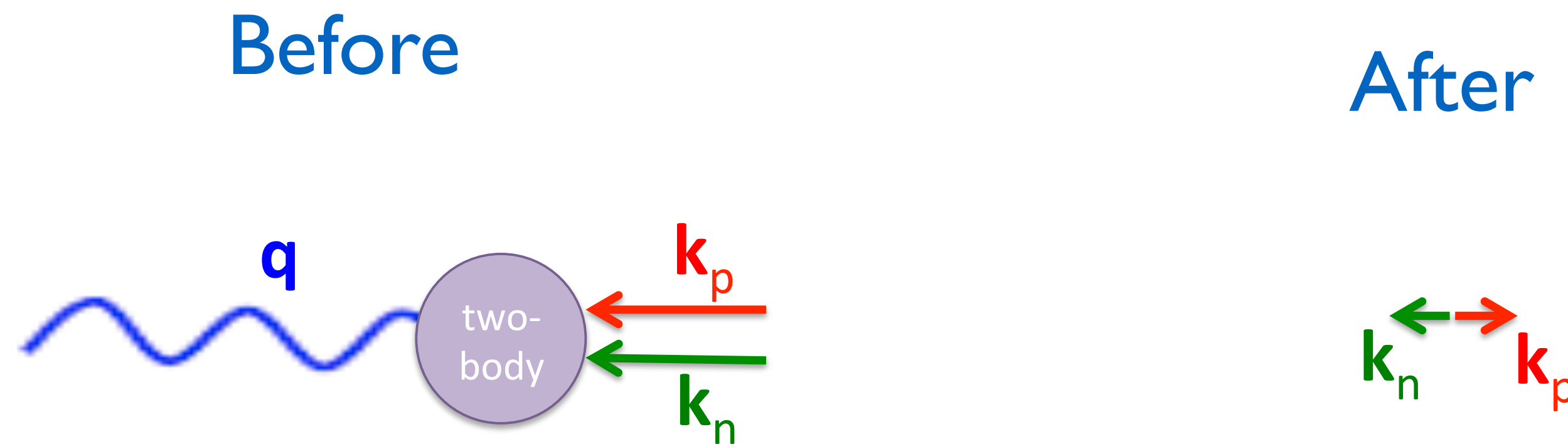
# Scale Dependence of SRC Interpretation

- Consider large  $q^2$  near threshold (small  $p'$ ) for  $\theta=0$  in **low-resolution** picture (COM frame of outgoing np)



# Scale Dependence of SRC Interpretation

- Consider large  $q^2$  near threshold (small  $p'$ ) for  $\theta=0$  in **low-resolution** picture (COM frame of outgoing np)



no large relative momentum in evolved deuteron wf

1-body current makes no contribution

∴ 2-body current mostly stops the low-relative momentum np pair



SRC

EFT

20% High-p  
Tails

Chiral  
Interactions

Hard  
Interactions

Soft  
Potentials

1B Reaction  
Currents

Transformed  
operators

## Summary/Questions

RG methods smoothly connect high- and low-resolution pictures. There is no "correct" picture (e.g., can reproduce SRC phenom. in a low resolution picture)

Interpretations vary with resolution scale (FSI, etc.), as do ease of calculations (simple wf's + complicated currents vs. complicated wf's + simple currents). Can we exploit this?

Can we use RG methods to connect SF's in low-resolution shell model picture and SRCs in high-resolution picture?

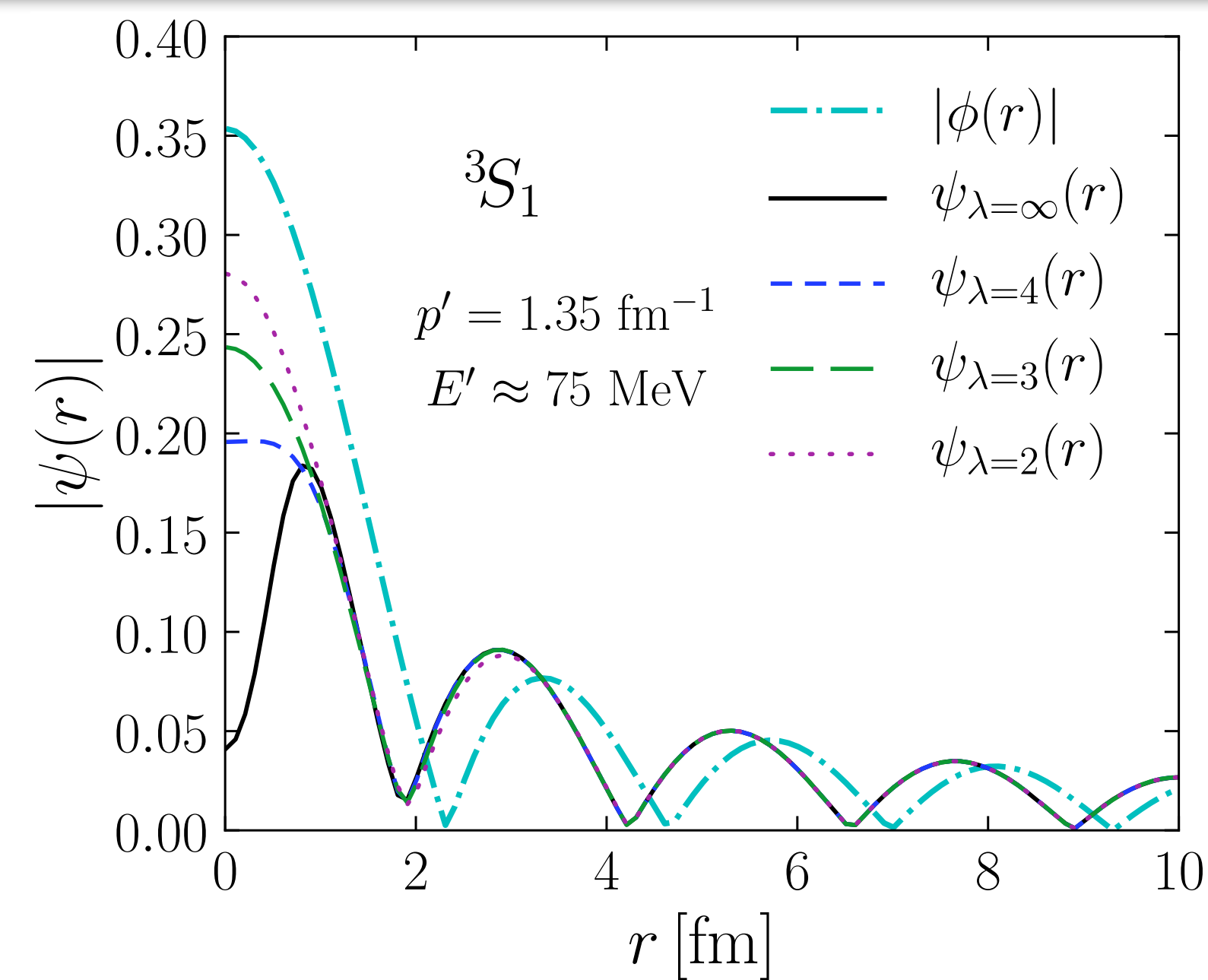
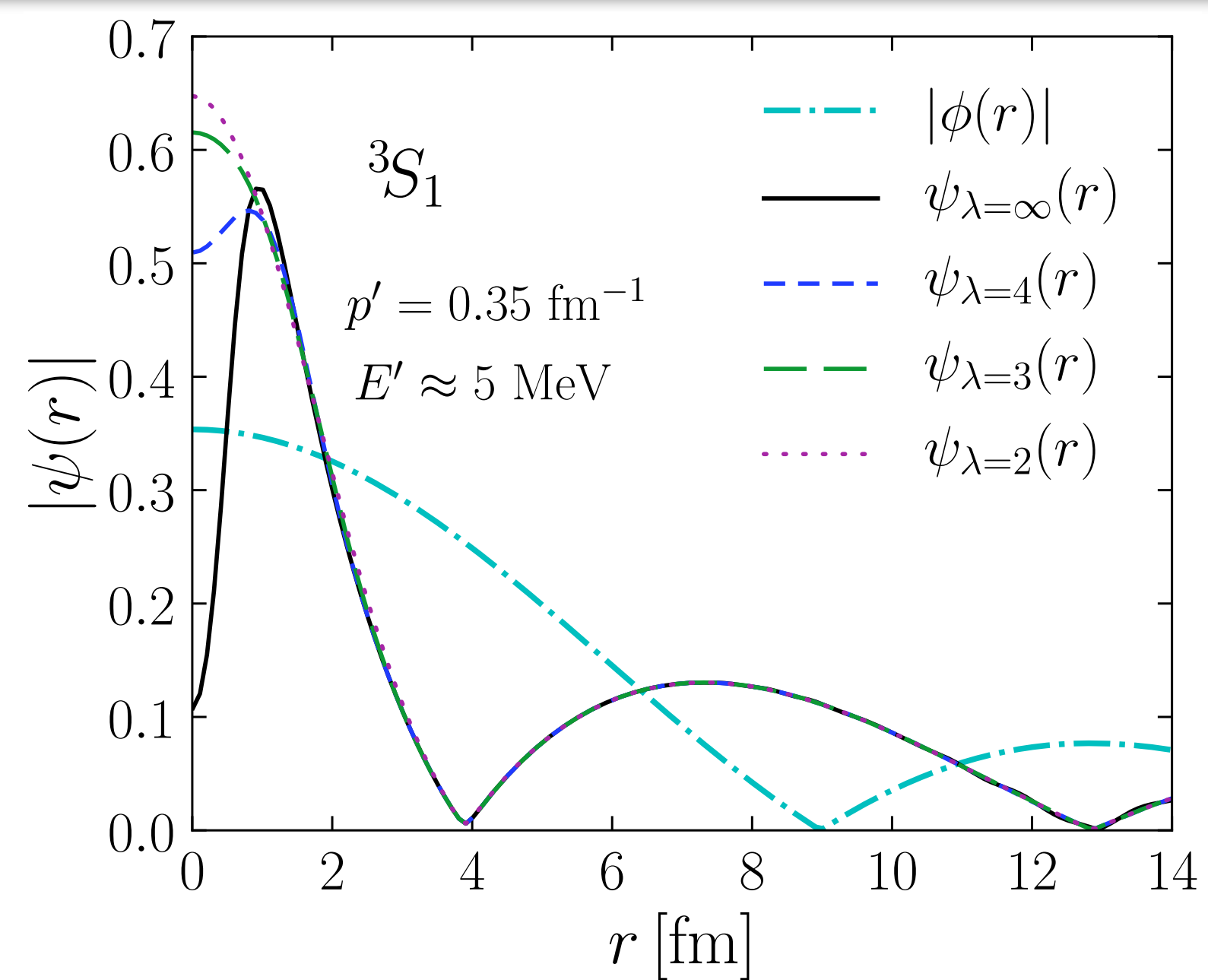
Can we use OPE + SRC/high- $q$  measurements to extend reach of low-resolution theories ?

Can we use simpler low-resolution wf's + OPE for to do high- $q$  electron scattering in medium mass nuclei?

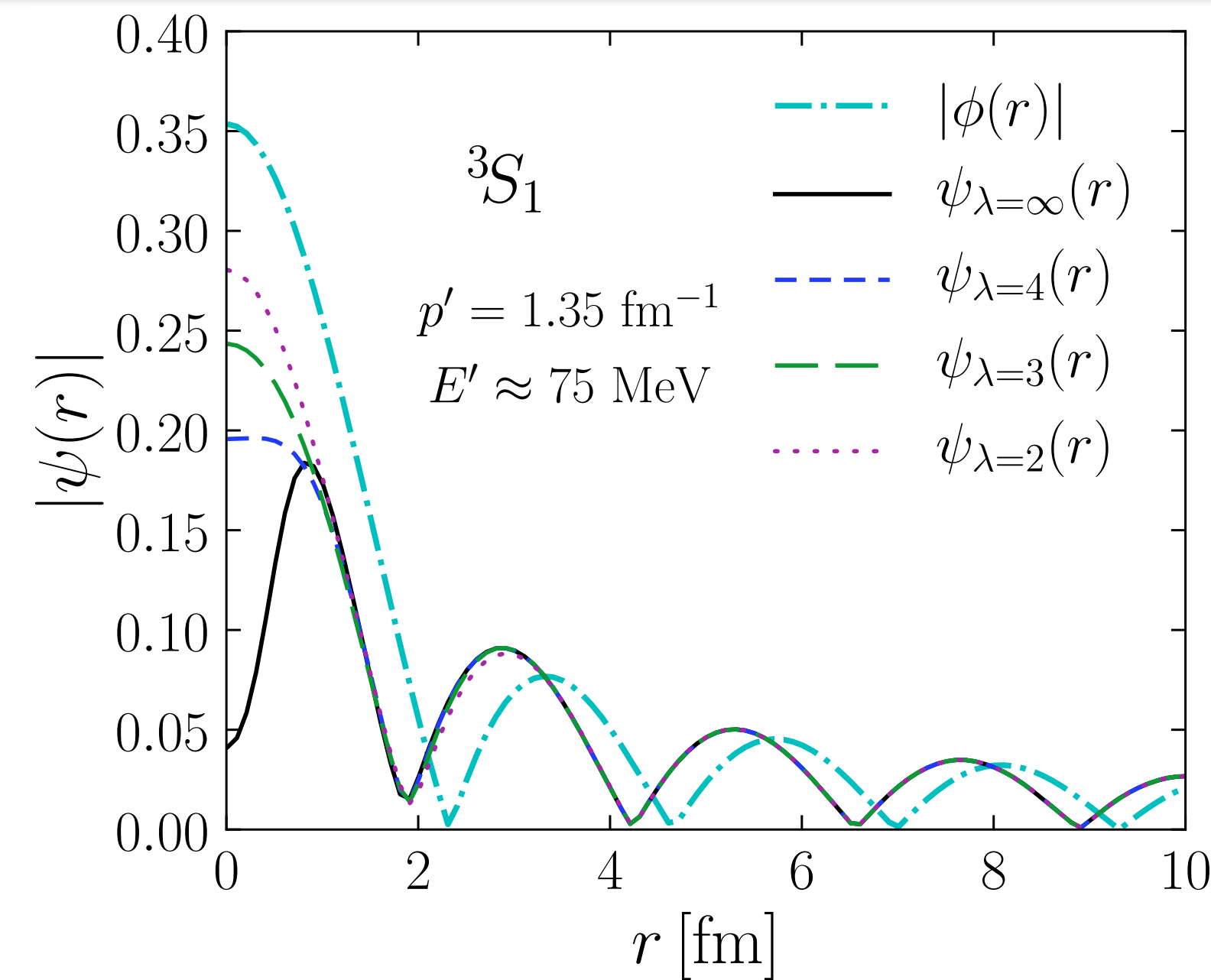
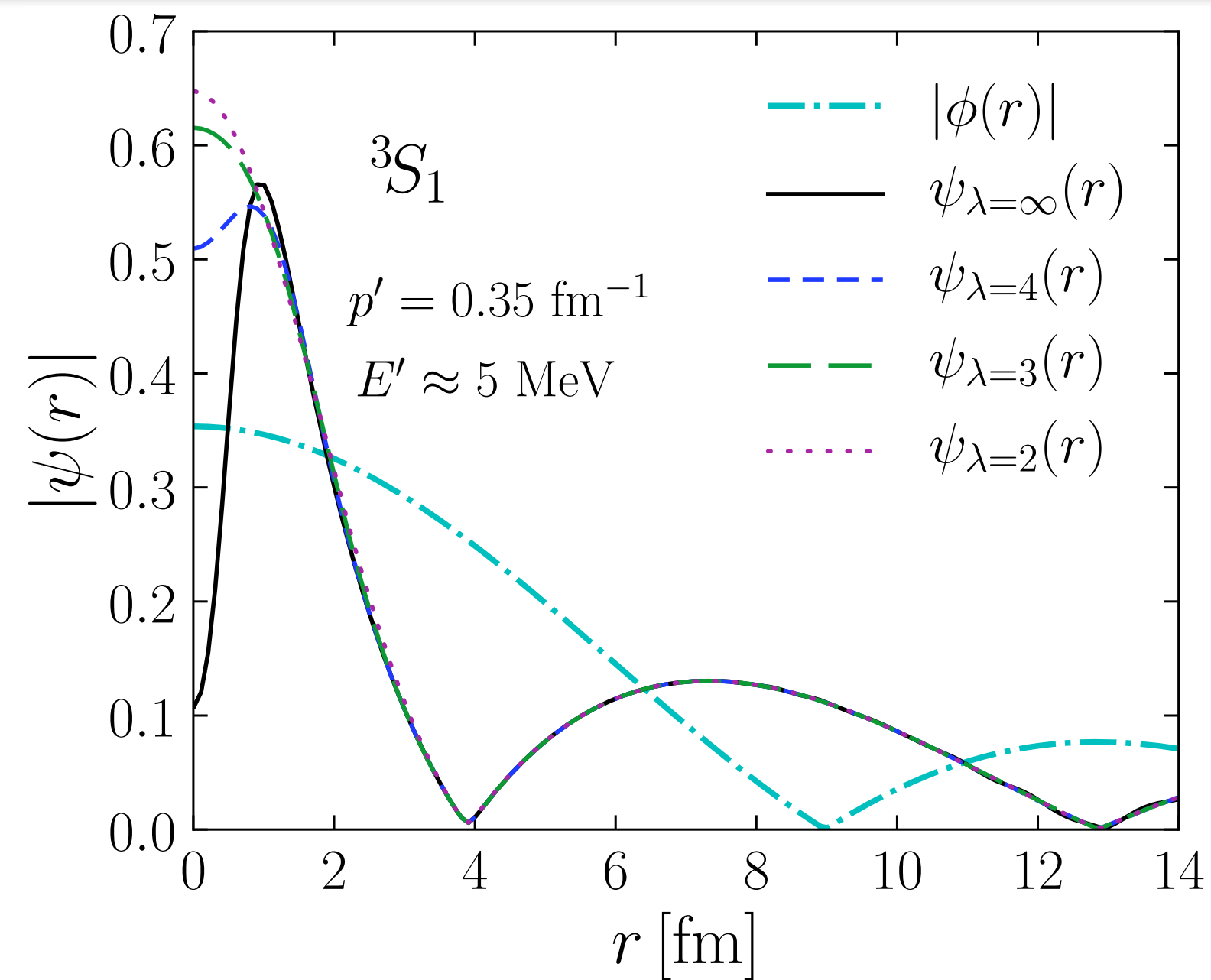


# Extras

# Final-state wave function evolution

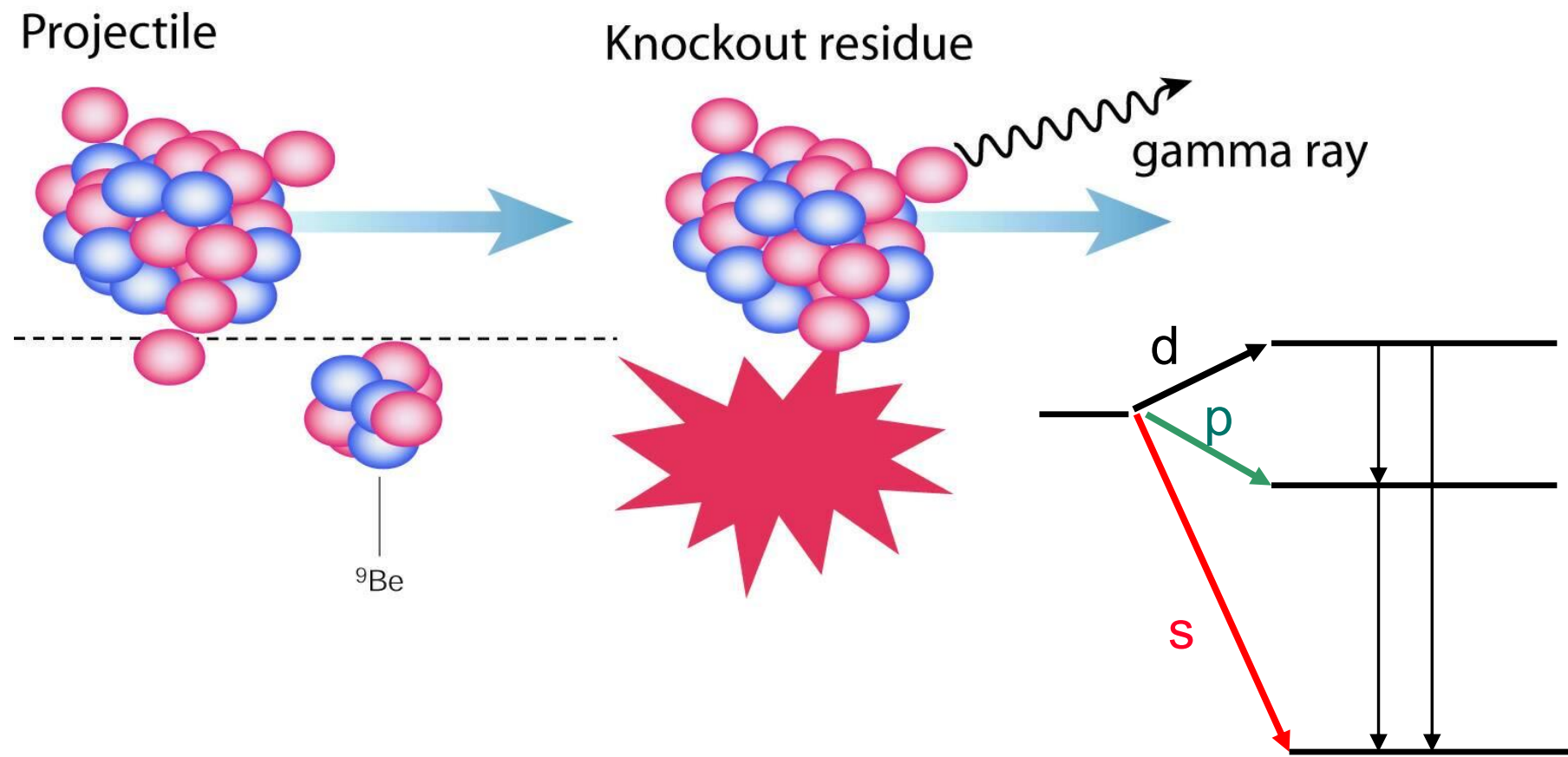


# Final-state wave function evolution



- Correlation “wound” at small  $r$  smoothed out under evolution
- Long-distance tail invariant (phase shifts preserved)

# Other exclusive knock-out reactions [pictures from A. Gade]



Exclusive reactions, theory vs. experiment

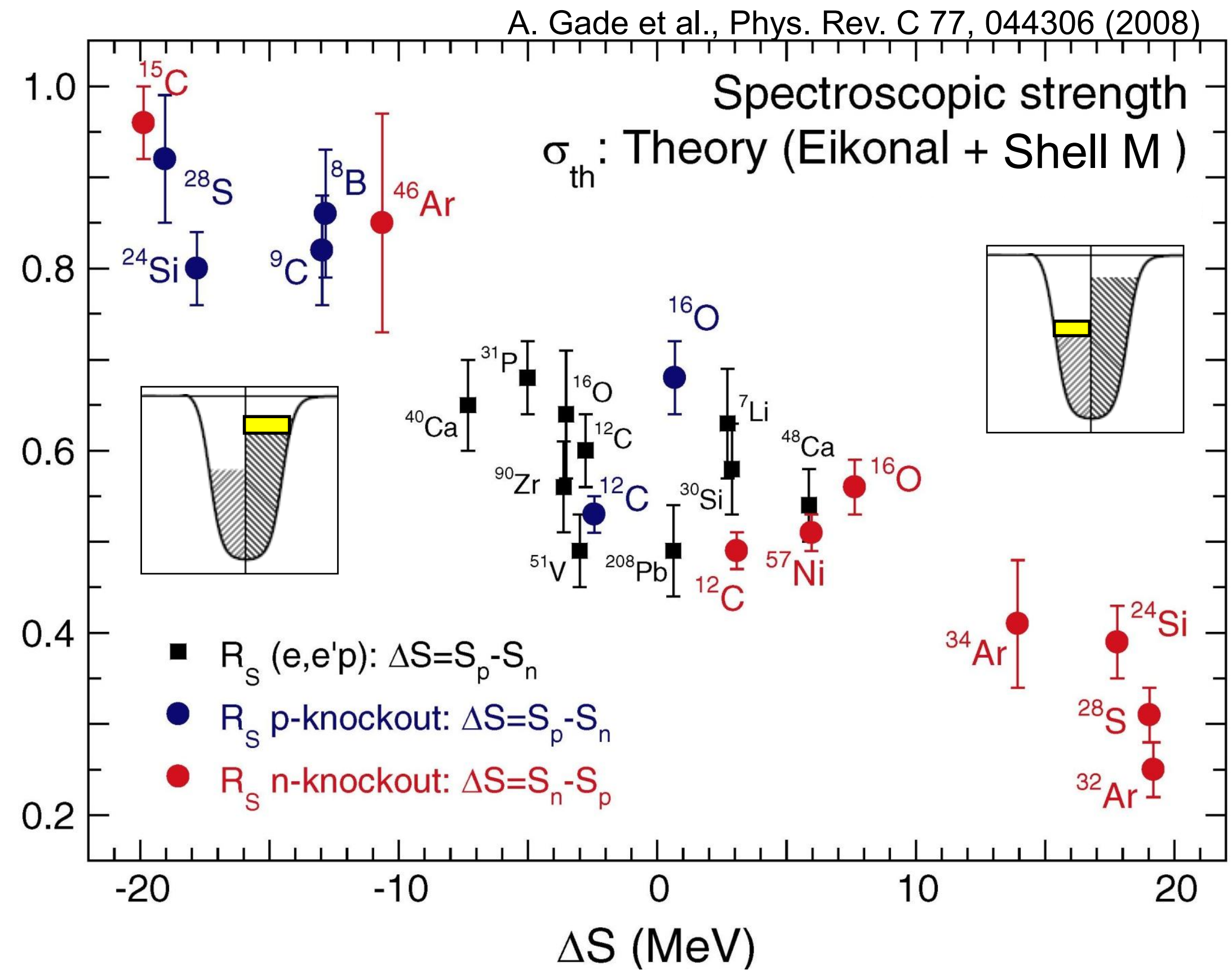
$$\sigma(j^\pi) = \left( \frac{A}{A-1} \right)^N C^2 S(j^\pi) \sigma_{sp}(j, S_N + E_x[j^\pi])$$

Structure theory

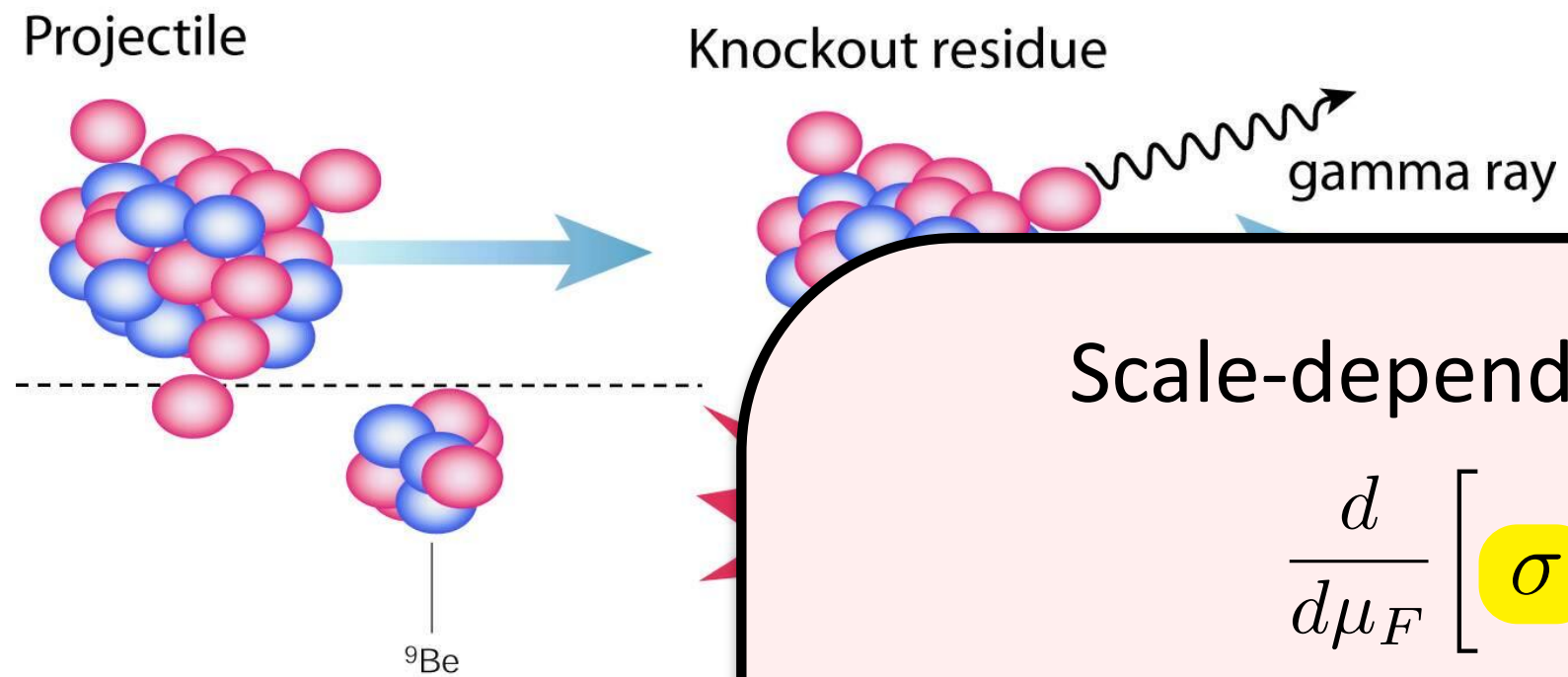
Reaction theory

Origin and systematics of  $R = \sigma_{\text{exp}} / \sigma_{\text{th}} < 1$  are not understood (includes e,e'p results)

$$R_S = \sigma_{\text{exp}} / \sigma_{\text{th}}$$



# Other exclusive knock-out reactions [pictures from A. Gade]



A. Gade et al., Phys. Rev. C 77, 044306 (2008)

Scale-dependent (RG) view of how these reactions are treated

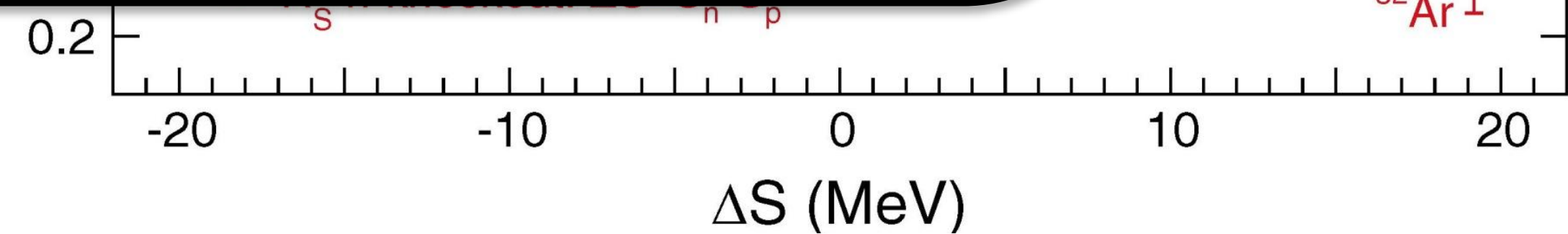
$$\frac{d}{d\mu_F} \left[ \sigma = \underbrace{\text{reaction}}_{\mu_F} \otimes \text{structure} \right] = 0$$

- Analysis mixes a high-resolution reaction mechanism (single-particle) with a low-resolution structure description.
- Theory is greater than experiment because missing induced current (e.g., 2-body for  $e^-$ ) does not exclude flux.
- **Plan: use SRG on reaction operator here and exploit factorization**

Exclusive reaction

$$\sigma(j^\pi) = \left( \frac{A}{A-1} \right)^N$$

Origin and systematics of  $R = \sigma_{\text{exp}} / \sigma_{\text{th}} < 1$  are not understood (includes  $e, e'p$  results)



# Deuteron electrodisintegration kinematics

