

Lifetime of the key ${}^{30}P(p,\gamma){}^{31}S$ resonance in novae

Barry Davids, Lijie Sun, Chris Wrede Spokespersons for S2373 TRIUMF NP-EEC Meeting, Jan 30, 2024











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Bottleneck for Nova Nucleosynthesis



Presolar grains of nova origin J. José *et al.*, Astrophys. J. 612, 414 (2004).

Nuclear thermometers for novae L. N. Downen *et al.*, Astrophys. J. 762, 105 (2013).

Nuclear mixing meters for novae K. J. Kelly *et al.*, Astrophys. J. 777, 130 (2013).



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J. José *et al.*, Nucl. Phys. A 777, 550 (2006). C. Wrede, AIP Advances 4, 041004 (2014).

Thermonuclear ${}^{30}P(p,\gamma){}^{31}S$ Reaction Rate

Resonant reaction rate

$$N_A \langle \sigma \nu \rangle_{\rm res} = N_A \left(\frac{2\pi}{\mu kT}\right)^{3/2} \hbar^2 \omega \gamma \exp\left(-\frac{E_r}{kT}\right)$$

Resonance strength

$$\omega \gamma = \frac{2J_r + 1}{(2J_p + 1)(2J_T + 1)} \frac{\Gamma_p \Gamma_\gamma}{\Gamma} \qquad \Gamma = \frac{\hbar}{\tau}$$
$$= \frac{2J_r + 1}{(2J_p + 1)(2J_T + 1)} \frac{B_p (1 - B_p)}{\tau}$$

C. Iliadis, Nuclear Physics of Stars, Wiley-VCH, 2015.





T. Budner et al., Phys. Rev. Lett. 128, 182701 (2022).

Key ³⁰P(*p*,γ)³¹S Resonance



S_p: A. Kankainen *et al.*, Phys. Rev. C 82, 052501(R) (2010). L. Canete *et al.*, Eur. Phys. J. A 52, 124 (2016).

FRIB

Facility for Rare Isotope Beams U.S. Department of Energy Office of Science | Michigan State University 640 South Shaw Lane • East Lansing, MI 48824, USA frib.msu.edu $J^{\pi} = 3/2^+$, $E_x = 6390.2(7)$ keV M. B. Bennett *et al.*, Phys. Rev. Lett. 116, 102502 (2016).

 $B_p = 2.5^{+0.4}_{-0.3} \times 10^{-4}$ T. Budner *et al.*, Phys. Rev. Lett. 128, 182701 (2022).

 $\tau = 4.4 \text{ fs} \quad (\text{USDA})$ $\tau = 26 \text{ fs} \quad (\text{USDB})$ $\tau = 3.4 \text{ fs} \quad (\text{USDC}), 1.3 \text{ fs} (\text{USDC-shifted})$ $\tau = 1.9 \text{ fs} \quad (\text{USDE})$ $\tau = 3.5 \text{ fs} \quad (\text{USDI})$ B. A. Brown. Shell Model Calculations. $\tau < 10 \text{ fs} \quad (^{31}\text{P} \ 3/2^{+} \ 6381 \text{ keV})$ E.O. de Neijs *et al.*, Nucl. Phys. A 254, 45 (1975).

 $\tau < 20 \text{ fs}$ (TRIUMF S1582) L. J. Sun *et al.*, Phys. Lett. B 839, 137801 (2023).

Doppler Shift Attenuation Method



$$E_{\gamma} = E_0 \frac{\sqrt{1-\beta^2}}{1-\beta\cos\theta}, \qquad \beta = \frac{\nu}{c}$$

D. Branford *et al.*, Nucl. Instrum. Methods 106, 437 (1973).
T. K. Alexander *et al.*, Adv. Nucl. Phys. 10, 197 (1978).
P. J. Nolan *et al.*, Rep. Prog. Phys. 42, 1 (1979).



Doppler Shift Lifetimes (DSL) @ ISAC-II





³He(³²S,αγ)³¹S Reaction with DSL1





DSL2 Upgrade

DSL Phase II (2020-)



 α detection efficiency ×11 γ detection efficiency ×1.3 ~1000 counts with 19 shifts

S2193	²³ Mg	L. E. Weghorn, PhD Thesis Project.
S2373	³¹ S	Proposed to NP-EEC 202401



Bayesian Uncertainty Quantification



S1582 L. J. Sun *et al.*, Phys. Lett. B 839, 137801 (2023).



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Ö. Sürer et al., Phys. Rev. C 106, 024607 (2022).

Prior Lineshape





Posterior Predictive Lineshape









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Resonance Strength





³⁰P(p, y)³¹S Reaction Rate



Monte Carlo reaction rate calculation <u>https://github.com/rlongland/RatesMC</u> Nova nucleosynthesis model calculation by Jordi José.



Summary

- •Our simulation has demonstrated that with 19 shifts, we can either obtain a finite value as low as 3 fs or set a strong upper limit of 2.5 fs on the lifetime of the astrophysically important $3/2^+$, 260-keV ${}^{30}P(p,\gamma){}^{31}S$ resonance.
- •This level of precision will put the ${}^{30}P(p,\gamma){}^{31}S$ reaction rate on a fully experimental footing for the first time, and will potentially eliminate the largest nuclear uncertainty associated with a series of isotopic and elemental abundance ratios in nova ejecta.
 - •The new rate will provide valuable insights into identifying the origin of several presolar grains using the ²⁹Si:²⁸Si and ³⁰Si:²⁸Si ratios.
 - The new rate will calibrate nova thermometers based on the elemental abundance ratios of **O:S, O:P, S:AI, and P:AI**.
 - The new rate will make the **Si/H** abundance ratio a more accurate constraint on the degree of mixing between the white dwarf's outer layers and the accreted envelope.



Collaboration

B. Davids	TRIUMF/SFU
L. J. Sun	MSU/FRIB
C. Wrede	MSU/FRIB
A. Adams	MSU/FRIB
C. Angus	University of York
A. Banerjee	SINP
T. Budner	ANL
M.Y.H. Chan	Northwestern U.
J. Chen	MSU/FRIB
J. Dopfer	MSU/FRIB
N. Esker	SJSU
M. Friedman	HUJI
C. Fry	LANL
A.B. Garnsworthy	TRIUMF
G. Hackman	TRIUMF
K. Hudson	SFU
J. Jose	UPC

V. Karayonchev O.S. Kirsebom L. Le R. Mahajan M. Oliver C. Ruiz R. Russell J. Surbrook Ö. Sürer V. Vedia L. Wagner L. Weghorn T. Wheeler J. Wilkinson E.J. Williams D. Yates

TRIUMF Dalhousie U. SJSU LSU/FSU SJSU TRIUMF University of Surrey LANL Miami U. TRIUMF TRIUMF MSU/FRIB MSU/FRIB LLNL TRIUMF UBC



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Known Lifetimes of ³¹S Excited States

³¹ S state		Evaluated τ (fs)	Literature
3/2+	1248	720(90)	Engmann NPA 1971 Doornenbal NIMA 2010 Tonev PLB 2021 Herlitzius Thesis 2013 Sun PLB 2023
5/2+	2234	290(80)	Engmann NPA 1971 Sun PLB 2023
1/2+	3076	<11	Sun PLB 2023
3/2+	3435	<16	Sun PLB 2023
7/2-	4450	790(250)	Pattabiraman PRC 2008 Tonev PLB 2021
(3/2)-	4971	<7	Sun PLB 2023
1/2+	5156	<15	Sun PLB 2023
9/2-	6376	245(45)	Pattabiraman PRC 2008
11/2(-)	6833	180(35)	Pattabiraman PRC 2008



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Known Decay Scheme of the Key Resonance



J. Chen et al., Nucl. Data Sheets 184, 29 (2022).



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γ-ray measurements of ³¹S states around 6.39 MeV

E_x (keV)	J^{π}_{i}	E_{γ} (keV)	E_f (keV)	J^{π}_{f}	Reaction / Decay	Reference
6202 7(E)	11/2+	1090.7(10)	5301	9/2+	$12C(20N_{0}, n)31S$	Jenkins
0393.7(3)		3042.4(4)	3351	7/2+		
6201 2(2)	11/2+	1091.2(4)	5301.7(3)	9/2+	28 5;(4U 0 p)31 5	Doherty
0394.2(2)		3042.9(1)	3351.3(2)	7/2+	31(-110,11)3	
6202	11/2+	1090.5(7)	5301	9/2+	$24M_{\odot}(12C_{\rm out})^{31}C_{\rm out}$	Testov
0393	11/2	3042.2(10)	3351	7/2+	mg(C, <i>anj</i> 3	
6392.5(2)	5/2+	5143.1(2)	1248.5(1)	3/2+	²⁸ Si(⁴ He, <i>n</i>) ³¹ S	Doherty
6392.5(2)	5/2+	5145(3)	1248	3/2+	2 H(30 P,n) 31 S	Kankainen
	2 /2+	2182.52(25)	4207.7(31)	3/2+		Bennett
		3106.28(31)	3283.76(31)	5/2+		
(200.2(7))		3313.56(33)	3076.40(31)	1/2+	31C1(P+3)31C	
0390.2(7)	5/2	4155.84(31)	2234.06(20)	5/2+	$-CI(p \gamma)^{-3}$	
		5141.3(6)	1248.43(20)	3/2+		
		6389.5(7)	0	1/2+		



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³He Implanted Target





Reaction Kinematics





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Geant4 Simulation





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Geant4 Simulation





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Solid Angle





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Geant4 Simulation





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Bayesian Model Emulation & Model Calibration





Bayesian-related talks at each APS DNP Meeting





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Advantages of Inverse Kinematics

While DSAM can in principle be done in either normal or inverse kinematics, there are several advantages to inverse kinematics, where a heavy beam impinges on a light target, including:

- •Recoils tend to be in the <u>electronic stopping power</u> regime. Electronic stopping powers are both more easily measured and have smaller angular scattering, leading to smaller uncertainties in the extracted lifetimes.
- <u>Velocity change can be larger</u>, due to a maximum in stopping powers for typical systems around recoil energies of tens of MeV for nuclei with mass number 20 < A < 40. This is particularly helpful for extracting finite lifetimes for short-lived states.
- Higher recoil velocities (β = 0.074) lead to <u>larger Doppler shifts</u>, making it easier to distinguish unshifted peaks from shifted peaks, particularly at low energies, where the absolute shift is small.

C. Fry, Ph.D. Thesis, Michigan State University, USA, 2018.



Readiness

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Isotope	Minimum Intensity (s-1)	Requested Intensity (s-1)	Target & Ion Source	Energy	Comment
³² S	0.6E+10	1.8E+10	OLIS	128 MeV	

Comments: ³²S 7+ was already delivered to DSL in Aug 2018 at 1E+11 pps, with acceptable contaminant levels. Thus this is feasible. However the new OLIS team (C. Charles) has not yet had experience of producing this beam so some practice time with OLIS should be taken into account well in advance of experiment.

Classification	for Beam Delive	ry or Developn	nent		
Proposals or P	rogress Reports:		Letters of In	tent:	
🛛 feasible	marginal	infeasible	trivial	complex	major

We possess all the equipment and targets needed for the experiment at TRIUMF.



Beam Delivery







 $\omega \gamma_{260} = 80(48) \,\mu eV$

T. Budner et al., Phys. Rev. Lett. 128, 182701 (2022).



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$$\omega \gamma_{260} = 22 - 62 \,\mu \text{eV}$$

1000 counts τ = 3 fs



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$$\omega \gamma_{260} = 18 - 57 \,\mu \text{eV}$$

500 counts τ = 3 fs



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Resonant Contribution				
Note: G1 = entrance channel, G2 = exit channel, G3 = spectator channel !! Ecm, Exf in (keV); wg, Gx in (eV) !				
Note: if Er<0, theta^2=C2S*theta_sp^2 must be entered instead of entrance channel partial width				
Ecm DEcm wg Dwg J G1 DG1 L1 G2 DG2 L2 G3 DG3 L3 Exf Int Corr/Frac				
27.05 0.38 1.1e-33 0.4e-33 3.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
124.45 0.38 9.5e-12 9.5e-12 0.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
148.45 0.27 2.1e-9 2.1e-9 1.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
196.15 0.55 5.1e-7 2.7e-7 1.5 0 0 0 0 0 0 0 0 0 0 0 0 0				
226.41 0.33 0.7e-6 0.7e-6 2.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
244.95 0.38 6.1e-8 2.1e-8 4.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
259.81 0.29 8.0e-5 4.8e-5 1.5 0 0 0 0 0 0 0 0 0 0 0 0 $\omega\gamma = 80(48) \mu\text{eV}$				
261.51 0.33 5.8e-7 2.5e-7 2.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
262.28 0.31 1.0e-30 1.0e-30 5.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
271.35 2.01 1.0e-30 1.0e-30 3.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
410.95 0.47 0.85e-4 0.85e-4 1.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
451.75 2.01 1.0e-30 1.0e-30 3.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
503.99 0.33 1.0e-30 1.0e-30 4.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
589.35 1.03 0.072 0.072 2.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0				
618.35 2.01 0.200 0.200 1.5 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0				

